

# Mid term Solution

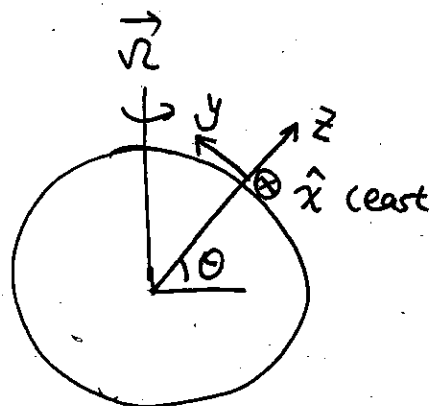
Prob 1

$$m \ddot{\vec{r}} = m \vec{g} + \vec{F}_c$$

Set the frame  $\hat{z}$  as the vertically up

$\hat{x}$  as the east

$\hat{y}$  as the north,



$\theta$  is the latitude angle. The earth spinning angular velocity

$$\vec{\omega} = (0, \omega \cos \theta, \omega \sin \theta)$$

(x2)  $\vec{r} \times \vec{\omega} = (\dot{y} \omega \sin \theta - \dot{z} \omega \cos \theta, -\dot{x} \omega \sin \theta, -\dot{x} \omega \cos \theta)$

$\Rightarrow \ddot{\vec{r}} = -g \hat{z} + 2 \vec{r} \times \vec{\omega} \Rightarrow \begin{cases} \ddot{x} = 2\omega (\dot{y} \sin \theta - \dot{z} \cos \theta) \\ \ddot{y} = -2\omega \dot{x} \sin \theta \\ \ddot{z} = -g + 2\omega \dot{x} \cos \theta \end{cases}$

1)

• zero th order solution

(x3)  $\begin{cases} \ddot{x}_0 = \ddot{y}_0 = 0 \\ \ddot{z}_0 = -g \end{cases} \Rightarrow \begin{cases} x_0(t) = y_0(t) = 0 \\ z_0(t) = v_0 t - \frac{1}{2} g t^2 \end{cases}$   
with  $v_0^2 = 2gh$

and the time for the ball returns to the ground is  $T = 2 \sqrt{2h/g}$ .

• 1st order  $\ddot{x}_1 = -2\omega \cos \theta \dot{z}_0 = -2\omega \cos \theta [v_0 - g t]$

with  $x_1(0) = 0, \dot{x}_1(0) = 0$

$\Rightarrow x_1 = -2\omega \cos \theta \left[ \frac{v_0}{2} t^2 - \frac{g}{6} t^3 \right]$

(x3)

$$\Rightarrow x_1 = -\Omega \cos\theta \left( v_0 - \frac{g}{3}t \right) t^2$$

plug in  $t = T = \sqrt{\frac{8h}{g}}$

$$v_0 = \sqrt{2gh}$$

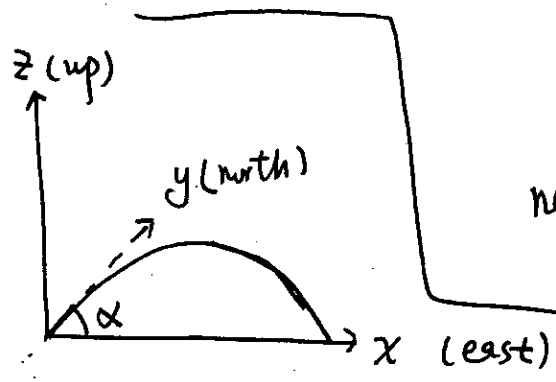
$$\Rightarrow x_1 = -\Omega \cos\theta \frac{v_0}{3} t^2 \quad (42)$$

$$= -\Omega \cos\theta \frac{1}{3} \sqrt{2gh} \frac{8h}{g}$$

$$= -\frac{8\Omega}{3} \sqrt{\frac{2h^3}{g}} \cos\theta$$

negative, deviation toward west.

2)



Zeroth order

$$\begin{cases} \ddot{x}_0 = \ddot{y}_0 = 0 \\ \ddot{z}_0 = -g \end{cases}$$

with  $\dot{z}(0) = v_0 \sin\alpha$

$\dot{x}(0) = v_0 \cos\alpha$

$$\Rightarrow \begin{cases} x_0(t) = v_0 \cos\alpha t \\ y_0(t) = 0 \\ z(t) = v_0 \sin\alpha t - \frac{1}{2} g t^2 \end{cases}$$

the flying time is

$$T = \frac{2v_0 \sin\alpha}{g}$$

Consider the transverse deviation - correct to 1st order

$$\ddot{y}_1 = -2\Omega \dot{x}_0 \sin\theta = -2\Omega v_0 \cos\alpha \sin\theta$$

plug in  $y_1(0) = 0, \dot{y}_1(0) = 0$

negative, deviation towards south.

$$\Rightarrow y_1(t) = -\Omega v_0 \cos\alpha \sin\theta t^2$$

$$\Rightarrow \begin{aligned} y(T) &= -\Omega v_0 \cos\alpha \sin\theta \left( \frac{2v_0 \sin\alpha}{g} \right)^2 \\ &= -\frac{4v_0^3 \sin^2\alpha \cos\alpha \sin\theta}{g^2} \end{aligned}$$

(42)

Prob 2

$$I_{ij} = \int dV \rho(\vec{r}) (r^2 \delta_{ij} - r_i r_j)$$

after the rotation  $r'_i = T_{ij} r_j$

$$I'_{ij} = \int dV' \rho'(\vec{r}') (r'^2 \delta_{ij} - r'_i r'_j) \quad +4$$

according to  $\rho'(\vec{r}') = \rho(\vec{r})$ ,  $r'^2 = r^2$  (-2)

$$\Rightarrow I'_{ij} = \int dV \rho(\vec{r}) (r^2 \delta_{ij} - T_{i i'} r_{i'} T_{j j'} r_{j'})$$

$$= T_{i i'} \left[ \int dV \rho(\vec{r}) (r^2 \delta_{i' j'} - r_{i'} r_{j'}) \right] T_{j' j}^T$$

$$I_{i' j'}$$

we used have  $T_{i i'} \delta_{i' j'} T_{j j'}^T = \delta_{ij}$  in the above expression

$$\Rightarrow I'_{ij} = T_{i i'} I_{i' j'} T_{j j'}^T$$

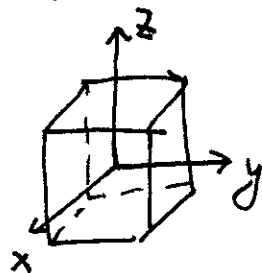
$$\text{or } \boxed{I' = T I T^{-1}}$$

(2) First we choose xyz axis in the symmetric position

we have  $I_{xy} = I_{yz} = I_{zx} = 0$

$$I_{xx} = I_{yy} = I_{zz} = \int_{-a/2}^{a/2} \int_{-a/2}^{a/2} \int_{-a/2}^{a/2} \rho (y^2 + z^2) dx dy dz$$

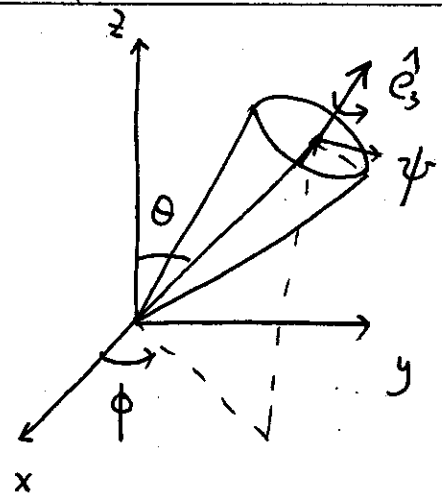
$$= \rho \cdot a^2 \cdot \left[ \frac{y^3}{3} \right]_{-a/2}^{a/2} \times 2 = 2 \rho a^2 \frac{a^3}{12} = \frac{M}{6} a^2 \Rightarrow I = \frac{M a^2}{6} \begin{bmatrix} 1 & & \\ & 1 & \\ & & 1 \end{bmatrix}$$



(2)  $\Rightarrow$  under an arbitrary frame  $I' = T I T^{-1} = I = \frac{M a^2}{6} \begin{bmatrix} 1 & & \\ & 1 & \\ & & 1 \end{bmatrix}$ .

Prob 3

1) at  $t=0$ , we have  $\vec{L} = \lambda_3 \omega \hat{e}_3$



(x3)

$$\Rightarrow L_3 = \lambda_3 \omega$$

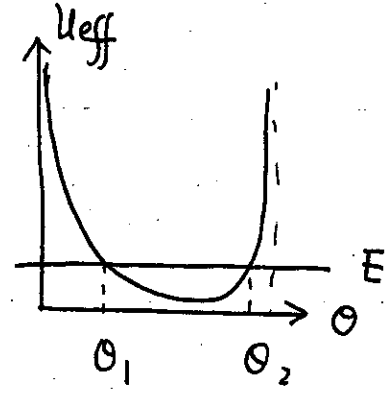
$$L_z = \vec{L} \cdot \hat{z} = \lambda_3 \omega \cos \theta_1$$

$$\Rightarrow E = \frac{1}{2} \lambda \dot{\theta}^2 + \frac{(\lambda_3 \omega)^2 (\cos \theta_1 - \cos \theta)^2}{2 \lambda \sin^2 \theta} + \frac{L_3^2}{2 \lambda_3} + M g R \omega \sin \theta$$

because at  $t=0$ ,  $\theta(t=0) = \theta_1$  and  $\dot{\theta}(t=0) = 0$

$\Rightarrow \theta_1$  is the turning point with

$$E = U_{eff}(\theta_1) = \frac{L_3^2}{2 \lambda_3} + M g R \omega \sin \theta_1$$



2) at another turning point  $\theta_2$ , we have  $\dot{\theta} = 0$

(x4)

$$\Rightarrow U_{eff}(\theta_2) = U_{eff}(\theta_1)$$

$$\Rightarrow \frac{(\lambda_3 \omega)^2 (\cos \theta_1 - \cos \theta_2)^2}{2 \lambda \sin^2 \theta_2} + M g R \omega \sin \theta_2 = M g R \omega \sin \theta_1$$

$$\Rightarrow \frac{(\lambda_3 \omega)^2 (\cos \theta_1 - \cos \theta_2)^2}{2 M g R \lambda} = \sin^2 \theta_2 = 1 - \cos^2 \theta_2$$

$$\Rightarrow 1 - \cos^2 \theta_2 - p (\cos \theta_1 - \cos \theta_2) = 0$$

x3

define  $x = \omega s \theta_1 - \omega s \theta_2 \Rightarrow \omega s \theta_2 = \omega s \theta_1 - x$   
 $\omega^2 s^2 \theta_2 = \omega^2 s^2 \theta_1 + x^2 - 2\omega s \theta_1 x$

$\Rightarrow 1 - \omega^2 s^2 \theta_1 - x^2 + 2\omega s \theta_1 x - p x = 0$

$\rightarrow x^2 + (p - 2\omega s \theta_1) x - \sin^2 \theta_1 = 0$  if

$\Rightarrow x = \frac{1}{2} \left[ -p + 2\omega s \theta_1 \pm \sqrt{(p - 2\omega s \theta_1)^2 + 4\sin^2 \theta_1} \right]$  We take the value  $x > 0$

i.e.  $\omega s \theta_1 > \omega s \theta_2$

(x1)

$\approx \frac{1}{2} \left[ -(p - 2\omega s \theta_1) + (p - 2\omega s \theta_1) \left( 1 + \frac{4\sin^2 \theta_1}{(p - 2\omega s \theta_1)^2} \right)^{\frac{1}{2}} \right]$

$\approx \frac{1}{2} \cdot \frac{4\sin^2 \theta_1}{p - 2\omega s \theta_1} \cdot \frac{1}{2} \approx \frac{\sin^2 \theta_1}{p - 2\omega s \theta_1} \approx \frac{\sin^2 \theta_1}{p}$

3)  $\Rightarrow \omega s \theta_1 - \omega s \theta_2 \approx \sin^2 \theta_1 \frac{2\lambda M g R}{\lambda_3^2 \omega^2}$

(x2)

as time evolves, due to resistance,  $\omega$  decreases thus nutation becomes more and more important.

4)  $\dot{\phi} = \frac{\lambda_3 \omega (\omega s \theta_1 - \omega s \theta_2)}{\lambda \sin^2 \theta}$  We replace  $\omega s \theta_1 - \omega s \theta_2$

as  $\frac{1}{2} (\omega s \theta_1 - \omega s \theta_2) = \frac{x}{2}$

at  $p \gg 1$ .

(x3)

we also replace  $\overline{\sin^2 \theta} = \sin^2 \theta_1$  at  $p \gg 1$

Same as we get in the limit  $p \gg 1$ .

$\Rightarrow \overline{\dot{\phi}} = \frac{\lambda_3 \omega x}{2\lambda \sin^2 \theta_1} = \frac{\lambda_3 \omega}{2p\lambda} = \frac{\lambda_3 \omega}{(\lambda_3 \omega)^2 / M g l} = \frac{M g l}{\lambda_3 \omega}$