

See p. 27 for range calculation

See also p. 34 for errors

CONTRIBUTIONS TO THE THEORY *
OF ATMOSPHERIC REFRACTION

Part II. Refraction Corrections in Satellite Geodesy

A. Photogrammetric Refraction

Depending on its definition, satellite photogrammetry deals with vertical photography of the earth's surface taken from an orbiting satellite, or with photography of an orbiting satellite taken from the surface of the earth against the stellar background. The photogrammetric refraction is essentially the same in both cases.

To determine the geometric relationship between astronomical refraction Δz_1 and photogrammetric refraction $\Delta\theta_2$, we have from Fig. 3

$$\overline{OQ} = \frac{r_1 \sin(z_1 + \Delta z_1)}{\sin z_2} \quad \overline{P_1Q} = \frac{r_1 \sin(z_1 + \Delta z_1 - z_2)}{\sin z_2}$$

and

$$d = (r_2 - \overline{OQ}) \sin z_2 = r_2 \sin z_2 - r_1 \sin(z_1 + \Delta z_1)$$

$$s = \overline{P_1Q} + d \operatorname{ctg} z_2 = r_2 \cos z_2 - r_1 \cos(z_1 + \Delta z_1)$$

The photogrammetric refraction is therefore obtained from the astronomical refraction by the formulas

$$\left. \begin{aligned} \Delta\theta_2 = \frac{d}{s} = \frac{r_2 \sin z_2 - r_1 \sin(z_1 + \Delta z_1)}{r_2 \cos z_2 - r_1 \cos(z_1 + \Delta z_1)} \end{aligned} \right\} \quad (45)$$

$$r_2 \sin z_2 = n_1 r_1 \sin z_1 \quad (2')$$

applicable to both cases of satellite photogrammetry defined above.

* -- Suite et fin de l'article publié dans les n^{os} 105 et 106 du Bulletin Géodésique.

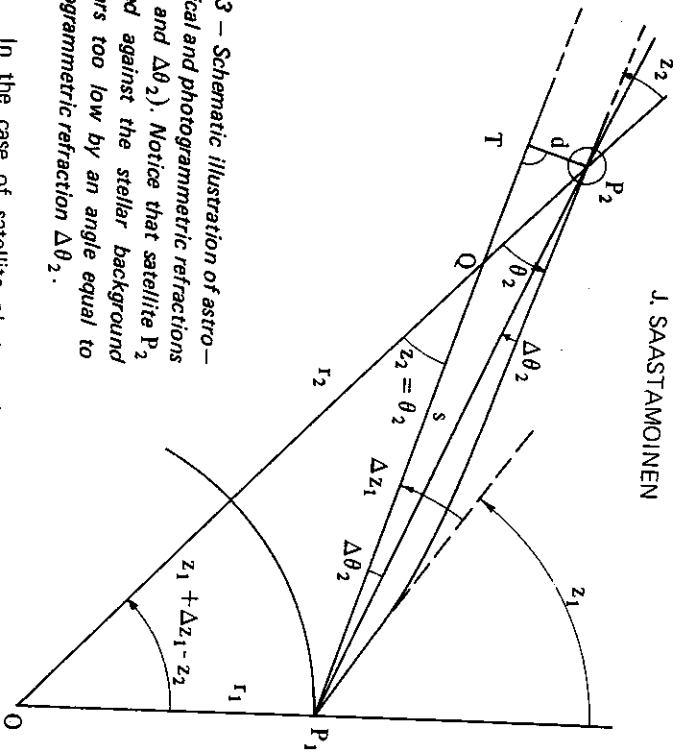


Fig. 3 — Schematic illustration of astronomical and photogrammetric refractions (Δz_1 and $\Delta \theta_2$). Notice that satellite P_2 viewed against the stellar background appears too low by an angle equal to photogrammetric refraction $\Delta \theta_2$.

In the case of satellite photography taken from a terrestrial camera station, distance s may be determined from orbital elements or computed by the formula

$$s = \sqrt{r_2^2 - n_1^2 r_1^2 \sin^2 z_1} - \sqrt{r_1^2 - (n_1 r_1 \sin z_1 - d)^2} \approx r_1 \left(\sqrt{(r_2/r_1)^2 - \sin^2 z_1} - \cos z_1 \right) \quad (46)$$

For departure d , we have from (30)

$$\begin{cases} \sin(\Delta z_1) \approx \Delta z_1 = \tan z_1 (n_1 - 1) \left[1 - \frac{RT_1}{r_1 g} \sec^2 z_1 + \frac{1}{2} \tan^2 z_1 (n_1 - 1) \right] + \delta \\ \cos(\Delta z_1) \approx 1 - \frac{1}{2} (\Delta z_1)^2 \end{cases}$$

which gives

$$d = \left[\frac{(n_0 - 1) RT_0 \tan z_1}{p_0 g \cos z_1} p_1 - r_1 \cos z_1 \delta \right] \left[1 - \frac{(n_0 - 1) T_0}{p_0} \tan^2 z_1 \left(\frac{p_1}{T_1} \right) \right] \quad (47)$$

where correction term $\delta = \delta_1 - \delta_2 - \delta_3 + \delta_4$ is expressed in radians. Numerically, assuming an effective wavelength of 0.554 microns,

$$d = 0.002317 \left[\frac{\tan z}{\cos z} \left(\frac{p}{1000} \right) - 13.35 \cos z \delta'' \right] \left[1 - 0.000079 \tan^2 z \left(\frac{p}{T} \right) \right] \quad (47a)$$

where d is the departure in kilometres, p is the total barometric pressure in millibars, T is the absolute temperature in degrees Kelvin, z is the apparent zenith distance of the satellite, and δ'' is the correction term given in Table III on page 387 (B.G. n° 106). Table VI gives the resulting values of $\Delta \theta_2$ at standard atmospheric pressure and temperature for different zenith distances and orbital heights. The satellite will appear too low against the stars; the image displacement (in microns) near the principal point of the photograph is equal to the tabular values multiplied by the focal length of the camera (in metres).

Table VI.

Differential Refraction in Microradians Between Satellite and Stellar Background of White Stars

Apparent Zenith Distance	Orbital Height					
	250 km	500 km	750 km	1000 km	1500 km	2000 km
10°	1.7	0.8	0.6	0.4	0.3	0.2
20°	3.4	1.7	1.1	0.9	0.6	0.4
30°	5.5	2.7	1.8	1.4	0.9	0.7
40°	8.0	4.0	2.7	2.1	1.4	1.1
50°	11.5	5.9	4.0	3.0	2.1	1.6
60°	16.8	8.8	6.1	4.7	3.3	2.6
70°	28.4	15.4	11.0	8.7	6.3	5.0
80°	69.1	41.2	30.9	25.3	19.2	15.8

Sunlight reflected from the satellite is nearly white in colour, and the differential refractions in Table VI therefore apply to white stars only. Differences in colour as well as in zenith distance of individual stars can be taken into account by corrections to their astronomical refractions.

In the case of vertical photography of the earth's surface taken from an orbiting satellite, the photogrammetric refraction is obtained similarly by computing first zenith distance z_1 from nadir distance θ_2 measured from the photograph:

$$\sin z_1 = (r_2 / r_1) \left[1 - \frac{(n_0 - 1) T_0}{p_0} \left(\frac{p_1}{T_1} \right) \right] \sin \theta_2 \quad (48)$$

or numerically, for $\lambda = 0.554$ microns:

$$\sin z = (r_2 / r_1) [1 - 0.000079 (p / T)] \sin \theta \quad (48a)$$

Table VII gives the resulting standard values of $\Delta\theta_2$ for super-wide angle photography from different camera heights. Due to these refraction angles, the photographic image points will appear to be shifted radially away from the principal point of the photograph; the linear displacements (in microns) are equal to the tabular values multiplied by $f \sec^2 \theta$, f denoting the focal length of the camera (in metres).

Table VII.

Photogrammetric Refraction in Microradians in Vertical Photography from Satellites

Apparent Nadir Distance	Orbital Height					
	250 km	500 km	750 km	1000 km	1500 km	2000 km
10°	1.7	0.9	0.6	0.5	0.3	0.3
20°	3.6	1.9	1.3	1.0	0.7	0.6
30°	5.8	3.1	2.2	1.7	1.3	1.1
40°	8.6	4.7	3.4	2.8	2.3	2.3
50°	12.7	7.3	5.8	5.4	6.9	-
59°	19.3	13.1	14.3	38.6	-	-

In photogrammetric mapping from low-altitude satellites, corrections for refraction will be required if horizontal control is extended by photogrammetric means. They can be obtained more conveniently using the approximate formulas

$$\left\{ \begin{aligned} A^2 &= \left(\frac{r}{r+h} \right)^2 - \sin^2 \theta \\ \Delta\theta &= \frac{2.32 p r \sin \theta}{(r+h)^2 A^2 (\cos \theta - A)} \end{aligned} \right. \quad (49)$$

where $\Delta\theta$ is the photogrammetric refraction in microradians at apparent nadir distance θ , p is the barometric pressure in millibars at the ground level, h is the height of the camera in kilometres above the ground level ($h > 50$ km), and r is the radius of the earth in kilometres.

B. — The Atmospheric Corrections for Troposphere and Stratosphere in Electromagnetic Ranging of Satellites

Due to the retarding effect of the atmosphere on the propagation of electromagnetic waves, a range correction

$$\Delta S = \int_S (n - 1) dS$$

where n is the refractive index referred to the group velocity of propagation, must be subtracted from observed electromagnetic distance S to obtain the true measured length of the effective ray path. As far as the electrically nonconducting lower atmosphere, the troposphere and the stratosphere, is concerned, this correction can be derived on the basis of the refraction theory developed previously in Part I for the determination of astronomical refraction.

Derivation of General Formula for Range Correction

In a spherically layered atmosphere, the basic mathematical expression of the range correction becomes

$$(0 \leq z_1 \leq 90^\circ); \quad \Delta S = \int_{r_1}^{r_1'} (n - 1) \sec z \, dr \quad (1^*)$$

which corresponds to equation (1) of astronomical refraction. Setting $n r / (n_1 r_1) = y$ for brevity, we now have from (2)

$$\sin^2 z = (\sec^2 z_1 - 1) / (y^2 \sec^2 z_1)$$

$$\cos^2 z = (y^2 \sec^2 z_1 - \sec^2 z_1 + 1) / (y^2 \sec^2 z_1)$$

and

$$\sec z = y \sec z_1 \left[1 + \sec^2 z_1 (y^2 - 1) \right]^{-\frac{1}{2}} = y \sec z_1 - \frac{1}{2} y (y^2 - 1) \sec^3 z_1 +$$

$$\frac{3}{8} y (y^2 - 1)^2 \sec^5 z_1 - \frac{5}{16} y (y^2 - 1)^3 \sec^7 z_1 + \frac{35}{128} y (y^2 - 1)^4 \sec^9 z_1 - \dots$$

Neglecting the subsequent terms in the binomial expansion, the first four may be written identically

$$\begin{aligned} \sec z &= \sec z_1 + \left(\frac{r}{r_1} - y \right) \left(\sec^3 z_1 - \sec z_1 \right) - \left(\frac{r-r_1}{r_1} \right) \left(\sec^3 z_1 - \sec z_1 \right) - \\ &- \left(1 + \frac{1}{2} y \right) (y-1)^2 \sec^3 z_1 + \frac{3}{8} y (y+1)^2 (y-1)^2 \sec^5 z_1 - \\ &- \frac{5}{16} y (y+1)^3 (y-1)^3 \sec^7 z_1 \end{aligned}$$

into which we substitute the approximations

$$\begin{aligned} \frac{r}{r_1} - y &= \left(\frac{r}{r_1} \right) \left(\frac{n_1 - n}{n_1} \right) = n_1 - n \\ \left(1 + \frac{1}{2} y \right) (y-1)^2 &= \frac{3}{2} (y-1)^2 = \frac{3}{2 r_1^2} (r-r_1)^2 \\ \frac{3}{8} y (y+1)^2 (y-1)^2 &= \frac{3}{2} (y-1)^2 = \frac{3}{2} \left[\left(\frac{r}{r_1} - 1 \right) - (n_1 - n) \right]^2 = \\ &= \frac{3}{2 r_1^2} (r-r_1)^2 - \frac{3}{r_1} (n_1 - n) (r-r_1) \\ \frac{5}{16} y (y+1)^3 (y-1)^3 &= \frac{5}{2 r_1^3} (r-r_1)^3 \end{aligned}$$

and obtain

$$\begin{aligned} \sec z &= \sec z_1 - A_1 (r-r_1) + A_1' (n_1 - n) + A_2 (r-r_1)^2 - \\ &- A_2' (n_1 - n) (r-r_1) - A_3 (r-r_1)^3 \end{aligned} \quad (3^*)$$

where the coefficients are :

$$\begin{aligned} A_1 &= (\sec^3 z_1 - \sec z_1) / r_1 \\ A_1' &= \sec^3 z_1 - \sec z_1 \\ A_2 &= 3 (\sec^5 z_1 - \sec^3 z_1) / (2 r_1^2) \\ A_2' &= 3 \sec^5 z_1 / r_1 \\ A_3 &= 5 \sec^7 z_1 / (2 r_1^3) \end{aligned} \quad (4^*)$$

By the substitution of (3*), integral (1*) breaks down into a linear function of six atmospheric integrals :

$$\begin{aligned} \Delta S &= \sec z_1 \int_{r_1}^{r'} (n-1) dr - A_1 \int_{r_1}^{r'} (n-1)(r-r_1) dr + \\ &+ A_1' \int_{r_1}^{r'} (n-1)(n_1 - n) dr + A_2 \int_{r_1}^{r'} (n-1)(r-r_1)^2 dr - \\ &- A_2' \int_{r_1}^{r'} (n-1)(n_1 - n)(r-r_1) dr - A_3 \int_{r_1}^{r'} (n-1)(r-r_1)^3 dr \end{aligned} \quad (5^*)$$

The first term in equation (5*) is given by refractivity integral (8) discussed previously. The five remaining integrals will now be determined, as follows.

$$\text{Integrals } \int_{r_1}^{r'} \frac{(n-1)(r-r_1)^q dr}{r_1}$$

Using the substitutions

$$\begin{aligned} u &= n - 1 & v &= (r - r_1)^q + 1 \\ du &= dn & dv &= (q + 1)(r - r_1)^{q-1} dr \end{aligned}$$

and integrating by parts, we obtain for these integrals

$$(q = 1, 2, 3); \quad \int_{r_1}^{r'} (n-1)(r-r_1)^q dr = \left(\frac{1}{q+1} \right) \int_{r_1}^{n_1} (r-r_1)^{q+1} dn \quad (50)$$

The values of the integrals on the right side have been determined previously in Part I, equations (20), (28) and (29).

$$\frac{\text{Integral}}{\int_{r_1}^{r'} \frac{(n-1)(n_1-n) dr}{g}}$$

Since identically

$$(n-1)(n_1-n) = (n_1-1)(n-1) - (n-1)^2$$

we obtain immediately, applying integrals (9), (22) and (23)

$$\int_{r_1}^{r'} (n-1)(n_1-n) dr = \frac{R}{g} (n_1-1)^2 T_1 - \left(\frac{R}{2g+R\beta} \right) [(n_1-1)^2 T_1 + \frac{1}{2} (R\beta/g)(n^0-1)^2 T^0] \quad (51)$$

$$\frac{\text{Integral}}{\int_{r_1}^{r'} \frac{(n-1)(n_1-n)(r-r_1) dr}{\beta}}$$

Using the same identity as in the previous case, we have first

$$\int_{r_1}^{r'} (n-1)(n_1-n)(r-r_1) dr = (n_1-1) \int_{r_1}^{r'} (n-1)(r-r_1) dr - \int_{r_1}^{r'} (n-1)^2 (r-r_1) dr$$

For the stratospheric component of the second integral above, we have from (11)

$$\begin{aligned} \int (n-1)^2 (r-r^0) dr &= (n^0-1)^2 \int (r-r^0) e^{2m} (r-r^0) dr = \\ &= (n^0-1)^2 \frac{e^{2m}(r-r^0)}{4m^2} [2m(r-r^0)-1] + C = \\ &= (n-1)^2 \left[\frac{1}{2m} (r-r^0) - \frac{1}{4m^2} \right] + C \end{aligned}$$

and

$$\int_{r^0}^{r'} (n-1)^2 (r-r^0) dr = \frac{(n^0-1)^2}{4m^2} = \frac{R^2}{4g^2} (n^0-1)^2 T^0$$

Consequently, using the identity $r-r_1 = r^0-r_1+r-r^0$ and applying integral (22), the total stratospheric component is obtained as

$$\int_{r^0}^{r'} (n-1)^2 (r-r_1) dr = \frac{R^2}{4g^2} (n^0-1)^2 T^0 + \frac{R}{2g} (r^0-r_1)(n^0-1)^2 T^0$$

For the tropospheric component of the same integral we have from (15) and (17)

$$\begin{aligned} \int (n-1)^2 (r-r_1) dr &= \frac{(n_1-1)^2 T_1}{\beta^2} \int \left(\frac{T}{T_1} \right)^{2m'} \left(\frac{T}{T_1} - 1 \right) dT = \\ &= \frac{(n_1-1)^2 T_1^2}{\beta^2} \left[\left(\frac{1}{2m'+2} \right) \left(\frac{T}{T_1} \right)^{2m'+2} - \left(\frac{1}{2m'+1} \right) \left(\frac{T}{T_1} \right)^{2m'+1} \right] + C = \\ &= \frac{(n-1)^2 T_1 T}{\beta^2} \left[\left(\frac{1}{2m'+2} \right) \frac{T}{T_1} - \frac{1}{2m'+1} \right] + C = \\ &= \frac{(n-1)^2 T_1 T}{\beta^2 (2m'+1)} \left[\left(\frac{2m'+1}{2m'+2} \right) \frac{T}{T_1} - 1 \right] + C = \frac{(n-1)^2 T_1 T}{\beta^2 (2m'+1)} \left[\frac{T}{T_1} - 1 - \left(\frac{1}{2m'+2} \right) \frac{T}{T_1} \right] + C = \\ &= \frac{(n-1)^2 (r-r_1) T}{\beta (2m'+1)} - \frac{(n-1)^2 T^2}{\beta^2 (2m'+1)(2m'+2)} + C = \\ &= - \left(\frac{R}{2g+R\beta} \right) (r-r_1)(n-1)^2 T - \frac{R^2}{2g(2g+R\beta)} (n-1)^2 T^2 + C \end{aligned}$$

The tropospheric component is accordingly

$$\begin{aligned} \int_{r_1}^{r'} (n-1)^2 (r-r_1) dr &= \frac{R^2}{2g(2g+R\beta)} [(n_1-1)^2 T_1^2 - (n^0-1)^2 T^0] - \\ &- \frac{R}{2g+R\beta} (r^0-r_1)(n^0-1)^2 T^0 \end{aligned}$$

which added to the stratospheric component gives the total value of the integral

$$\int_{r_1}^{r_1'} (n-1)^2 (r-r_1) dr = \frac{R^2}{2g^2} \left[\frac{(n_1-1)^2 T_1^2 - (n^0-1)^2 T_1 T^0}{2 + R\beta/g} + \frac{1}{2} (n^0-1)^2 T^0 \right] \quad (52)$$

and finally, in view of equations (50) and (20)

$$\int_{r_1}^{r_1'} (n-1) (n_1-n) (r-r_1) dr = \frac{R^2}{g} (n_1-1) \left[\frac{(n_1-1) T_1^2 - (n^0-1) T^0}{1 - R\beta/g} + (n^0-1) T^0 \right] - \frac{R^2}{2g^2} \left[\frac{(n_1-1)^2 T_1^2 - (n^0-1)^2 T_1 T^0}{2 + R\beta/g} + \frac{1}{2} (n^0-1)^2 T^0 \right] \quad (53)$$

Combining the results from the preceding developments, we may now write on the basis of equation (5*) the following expression for the range correction :

$$\Delta S = \frac{R}{g} \sec z_1 (n_1-1) T_1 - \frac{R^2}{r_1 g^2} (\sec^3 z_1 - \sec z_1) \left[\frac{(n_1-1) T_1^2 - (n^0-1) T^0}{1 - R\beta/g} + (n^0-1) T^0 \right] + \delta_5 + \delta_6 - \delta_7 - \delta_8 \quad (30^*)$$

where

$$\delta_5 = \frac{R}{g} (\sec^3 z_1 - \sec z_1) \left[\frac{(n_1-1)^2 T_1^2 - (n_1-1)^2 T_1^2 + \frac{1}{2} (R\beta/g)(n^0-1)^2 T^0}{2 + R\beta/g} \right]$$

$$\delta_6 = \frac{3R^3}{r_1^2 g^3} (\sec^5 z_1 - \sec^3 z_1) \left[\frac{(n_1-1) T_1^3 - (n^0-1) T^0}{(1 - R\beta/g)(1 - 2R\beta/g)} + (n^0-1) T^0 \right] + \frac{3R^2}{r_1^2 g^2} (\sec^5 z_1 - \sec^3 z_1) \left(1 - \frac{1}{1 - R\beta/g} \right) (r^0 - r_1) (n^0-1) T^0$$

$$\delta_7 = \frac{3R^2}{r_1 g^2} \sec^5 z_1 (n_1-1) \left[\frac{(n_1-1) T_1^2 - (n^0-1) T^0}{1 - R\beta/g} + (n^0-1) T^0 \right] - \frac{3R^2}{2r_1 g^2} \sec^5 z_1 \left[\frac{(n_1-1)^2 T_1^2 - (n^0-1)^2 T_1 T^0}{2 + R\beta/g} + \frac{1}{2} (n^0-1)^2 T^0 \right]$$

$$\delta_8 = \frac{15R^4}{r_1^3 g^4} \sec^7 z_1 \left[\frac{(n_1-1) T_1^4 - (n^0-1) T^0}{(1 - R\beta/g)(1 - 2R\beta/g)(1 - 3R\beta/g)} + (n^0-1) T^0 \right] + \frac{15R^3}{r_1^3 g^3} \sec^7 z_1 \left[\frac{1}{(1 - R\beta/g)(1 - 2R\beta/g)} \right] (r^0 - r_1) (n^0-1) T^0 + \frac{15R^2}{2r_1^3 g^2} \sec^7 z_1 \left(1 - \frac{1}{1 - R\beta/g} \right) (r^0 - r_1)^2 (n^0-1) T^0$$

represent minor correction terms, and the primed quantities refer to values corrected for ground inversion. The formal accuracy of formula (30*) has been tested in Tables IXa - IXc on the three atmospheric models introduced earlier. The refractive indices were computed using the Essen and Froome formula for radio microwaves.

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Atmospheric Model No. 2
Tropical Zone)

$$\Delta S = \int_{z_1}^{z_2} (n-1) \sec z \, dz, \text{ for } z = 60^\circ, 70^\circ, 80^\circ$$

Range Correction, $\Delta S = \int_{z_1}^{z_2} (n-1) \sec z \, dz, \text{ for } z = 60^\circ, 70^\circ, 80^\circ$

z_1 (km)	z_2 (km)	$z_1=60^\circ$	$z_1=70^\circ$	$z_1=80^\circ$
6380+				
0	276.37	2.00000	2.92380	5.75877
0.5	262.20	1.99961	2.92237	5.74674
1	250.51	1.99922	2.92092	5.73469
1.5	238.30	1.99882	2.91947	5.72264
2	226.55	1.99842	2.91800	5.71058
2.5	215.24	1.99802	2.91654	5.69851
3	204.38	1.99762	2.91506	5.68644
3.5	193.94	1.99721	2.91358	5.67439
4	183.92	1.99680	2.91209	5.66234
4.8	168.71	1.99615	2.90969	5.64930
5.6	154.49	1.99549	2.90728	5.63624
6.4	141.20	1.99482	2.90486	5.62318
7.2	128.81	1.99415	2.90242	5.61012
8	117.27	1.99347	2.89998	5.59706
8.8	106.54	1.99280	2.89752	5.58400
9.6	96.58	1.99211	2.89505	5.57094
10.4	87.35	1.99143	2.89258	5.55788
11.2	78.96	1.99074	2.89010	5.54482
12.1	71.39	1.99005	2.88762	5.53176
13.8	51.29	1.98851	2.88208	5.50923
15.5	39.30	1.98702	2.87655	5.48670
17.2	30.11	1.98552	2.87102	5.46417

Atmospheric Model No. 1
(Tropical Zone)

$$\Delta S = \int_{z_1}^{z_2} (n-1) \sec z \, dz, \text{ for } z = 60^\circ, 70^\circ, 80^\circ$$

Range Correction, $\Delta S = \int_{z_1}^{z_2} (n-1) \sec z \, dz, \text{ for } z = 60^\circ, 70^\circ, 80^\circ$

z_1 (km)	z_2 (km)	$z_1=60^\circ$	$z_1=70^\circ$	$z_1=80^\circ$
6360+				
0	261.46	2.00000	2.92380	5.75877
0.8	242.49	1.99936	2.92145	5.73910
1.6	224.61	1.99871	2.91908	5.71943
2.4	207.78	1.99806	2.91669	5.69978
3.2	191.96	1.99741	2.91429	5.68016
4	177.09	1.99674	2.91187	5.66057
4.8	163.15	1.99608	2.90944	5.64104
5.6	150.08	1.99541	2.90700	5.62156
6.4	137.84	1.99473	2.90454	5.60215
7.2	126.47	1.99405	2.90209	5.58277
8	115.05	1.99337	2.89964	5.56339
8.8	104.58	1.99269	2.89719	5.54400
9.6	94.05	1.99201	2.89474	5.52461
10.4	84.47	1.99133	2.89229	5.50522
11.2	74.84	1.99065	2.88984	5.48583
12.1	66.16	1.99000	2.88739	5.46644
13.8	48.32	1.98831	2.88294	5.42705
15.5	32.49	1.98662	2.87849	5.38766
17.2	20.65	1.98493	2.87404	5.34827

Integrals:

0 - 4 km: 1.82061 2.55859 6.7472 13.2895

4 - 10.4 km: 1.67324 2.43594 4.69389 5.20519

10.4 - 24 km: 0.97682 1.41470 2.55858 3.1087

24 - 40 km: 0.12024 0.17257 0.31087 0.4991

40 - 72 km: 0.01053 0.01495 0.02578 0.04191

Formula (30°):

1st term: 4.6154 6.7472 13.2895

2nd term: -0.0158 -0.0580 -0.4870

δ_5 : 0.0017 0.0064 0.0534

δ_6 : 0.0002 0.0016 0.0511

$-\delta_7$: -0.0000 -0.0003 -0.0095

$-\delta_8$: -0.0000 -0.0001 -0.0091

See footnote to Table 1A.

Integrals:

0 - 6.4 km: 2.48697 3.62931 7.08185

6.4 - 16.8 km: 1.65816 2.40806 4.58328

16.8 - 24 km: 0.31461 0.45419 0.84005

24 - 40 km: 0.11964 0.17171 0.30937

40 - 72 km: 0.00805 0.01143 0.01971

Formula (30°):

1st term: 4.6020 6.7277 13.2511

2nd term: -0.0164 -0.0603 -0.5064

δ_5 : 0.0016 0.0060 0.0504

δ_6 : 0.0002 0.0016 0.0536

$-\delta_7$: -0.0000 -0.0003 -0.0093

$-\delta_8$: -0.0000 -0.0001 -0.0091

See footnote to Table 1A.

Table IXc.
Atmospheric Model No. 3
(Arctic Zone)

$$\text{Range Correction, } \Delta S = \int_{z_1}^{z_2} (n-1) \sec z \, dz, \text{ for } z = 60^\circ, 70^\circ \text{ \& } 80^\circ$$

r, km 6400+	(n-1) 10 ⁶	sec z			(n-1) 10 ³ sec z			r, km 6400+	(n-1) 10 ⁶	sec z			(n-1) 10 ³ sec z		
		z ₁ =60°	z ₁ =70°	z ₁ =80°	z ₁ =60°	z ₁ =70°	z ₁ =80°			z ₁ =60°	z ₁ =70°	z ₁ =80°	z ₁ =60°	z ₁ =70°	z ₁ =80°
0	313.57	2.00000	2.92380	5.75877	0.62714	0.91682	1.80578	18.3	24.60	1.98482	2.86889	5.33617	0.04883	0.07058	0.13129
0.2	302.60	1.99988	2.92336	5.75503	0.60517	0.88462	1.74149	20.2	18.38	1.98313	2.86290	5.29420	0.03645	0.05263	0.09732
0.4	292.11	1.99975	2.92290	5.75119	0.58415	0.85381	1.67998	22.1	13.73	1.98144	2.85692	5.25303	0.02721	0.03924	0.07215
0.6	282.06	1.99963	2.92243	5.74728	0.56402	0.82432	1.62111	24	10.26	1.97974	2.85096	4.21267	0.02032	0.02926	0.05349
0.8	272.44	1.99950	2.92195	5.74333	0.54475	0.79607	1.56474	26	7.55	1.97796	2.84471	5.17109	0.01493	0.02148	0.03904
1	263.23	1.99936	2.92147	5.73926	0.52629	0.76902	1.51074	28	5.56	1.97619	2.83850	5.13043	0.01098	0.01577	0.02850
1.2	254.40	1.99923	2.92098	5.73516	0.50860	0.74309	1.45901	30	4.09	1.97441	2.83232	5.09069	0.00807	0.01158	0.02081
1.4	245.93	1.99909	2.92048	5.73100	0.49164	0.71823	1.40943	32	3.01	1.97264	2.82618	5.05184	0.00593	0.00850	0.01519
1.6	237.81	1.99896	2.91997	5.72678	0.47537	0.69440	1.36189	34	2.21	1.97087	2.82007	5.01387	0.00436	0.00624	0.01109
1.6	237.81	1.99896	2.91997	5.72678	0.47537	0.69440	1.36189	36	1.63	1.96911	2.81401	4.97675	0.00321	0.00458	0.00810
2.5	216.62	1.99824	2.91735	5.70517	0.43287	0.63197	1.23588	38	1.20	1.96736	2.80799	4.94046	0.00236	0.00336	0.00592
3.4	196.91	1.99752	2.91470	5.68354	0.39333	0.57394	1.11915	40	0.88	1.96561	2.80202	4.90498	0.00173	0.00247	0.00432
4.3	178.60	1.99679	2.91204	5.66192	0.35663	0.52009	1.01122	44	0.48	1.96213	2.79019	4.83631	0.00094	0.00133	0.00231
5.2	161.62	1.99605	2.90935	5.64032	0.32260	0.47021	0.91158	48	0.26	1.95867	2.77852	4.77054	0.00051	0.00072	0.00123
6.1	145.90	1.99531	2.90664	5.61876	0.29112	0.42408	0.81978	52	0.14	1.95524	2.76703	4.70750	0.00027	0.00039	0.00066
7	131.38	1.99456	2.90392	5.59725	0.26204	0.38151	0.73536	56	0.08	1.95183	2.75569	4.64699	0.00015	0.00021	0.00035
7.9	117.99	1.99381	2.90118	5.57581	0.23525	0.34231	0.65788	60	0.04	1.94844	2.74451	4.58886	0.00008	0.00011	0.00019
8.8	105.67	1.99305	2.89843	5.55444	0.21060	0.30627	0.58692	64	0.02	1.94508	2.73349	4.53296	0.00004	0.00006	0.00010
8.8	105.67	1.99305	2.89843	5.55444	0.21060	0.30627	0.58692	66	0.01	1.94174	2.72261	4.47916	0.00002	0.00003	0.00005
10.7	78.95	1.99145	2.89266	5.51026	0.15722	0.22837	0.43503	72	0.007	1.93843	2.71189	4.42732	0.00001	0.00002	0.00003
12.6	58.99	1.98982	2.88679	5.46610	0.11737	0.17028	0.32243								
14.5	44.07	1.98817	2.88086	5.42223	0.08762	0.12697	0.23897								
16.4	32.93	1.98650	2.87488	5.37888	0.06541	0.09467	0.17712								

Integrals *			
0 - 1.6 km :	0.87507	1.27879	2.51375
1.6 - 8.8 km :	2.37161	3.45773	6.71402
8.8 - 24 km :	1.23684	1.79310	3.38433
24 - 40 km :	0.12078	0.17337	0.31256
40 - 72 km :	0.01118	0.01588	0.02738
Range Correction, metres :	4.615	6.719	12.952
Formula (30*) :			
1st term :	4.6287	6.7667	13.3277
2nd term :	-0.0155	-0.0568	-0.4770
δ _g :	0.0017	0.0064	0.0537
δ ₆ :	0.0002	0.0015	0.0501
-δ ₇ :	-0.0000	-0.0003	-0.0095
-δ ₈ :	-0.0000	-0.0001	-0.0088
	4.615	6.717	12.936

* See footnote to Table IIa.

THEORY OF ATMOSPHERIC REFRACTION

Introduction to Practical Computation of Range Correction
Group Index of Refraction.

Radiant energy, such as an electromagnetic pulse or signal, travels with a speed referred to as the group velocity of propagation. The group velocity is different from the wave velocity in the case when the latter depends on wavelength λ, the group index of refraction being obtained from refractive index n by the formula

$$n - \lambda \frac{dn}{d\lambda}$$

The group index must be used in electromagnetic distance measurement in the computation of the velocity of propagation. It does not apply to the bending of a ray, which phenomenon depends on the wave velocity alone.

For visible light, dn/dλ can be determined with sufficient accuracy by differentiating equation (33), which gives the following formula for the group refractivity of standard air :

$$n_g - 1 = \left(\frac{173.3 + 1/\lambda^2}{173.3 - 1/\lambda^2} \right) (n_0 - 1) \quad (54)$$

where wavelength λ is expressed in microns. In laser ranging, therefore, range correction (30*) must be multiplied by factor (173.3 + 1/λ²) / (173.3 - 1/λ²) due to the effect of the group velocity of waves.

At the frequency range of radio microwaves, the atmosphere is considered nondispersive ; i.e. any waveform will be propagated without change of shape. In this case the group index is equal to the refractive index. Of the various formulas for the computation of the latter, we shall here quote the Essen and Froome formula adopted by the International Association of Geodesy [1963], converted into metric units :

$$(n-1) 10^6 = 77.624 (p/T) - 12.92 (e/T) + 371900 (e/T^2) \quad (55)$$

where p is the total pressure and e is the partial pressure of water vapour, both expressed in millibars, and T is the absolute temperature in degrees Kelvin.

Geometric Correction.

Since the effective ray path is curved due to the refraction, the difference in length between the measured arc and the corresponding chord

$$\delta_g = S - s$$

must further be subtracted from the measured distance. The total range correction then becomes

$$\Delta s = \Delta S + \delta_g \tag{56}$$

where δ_g is a geometric correction.

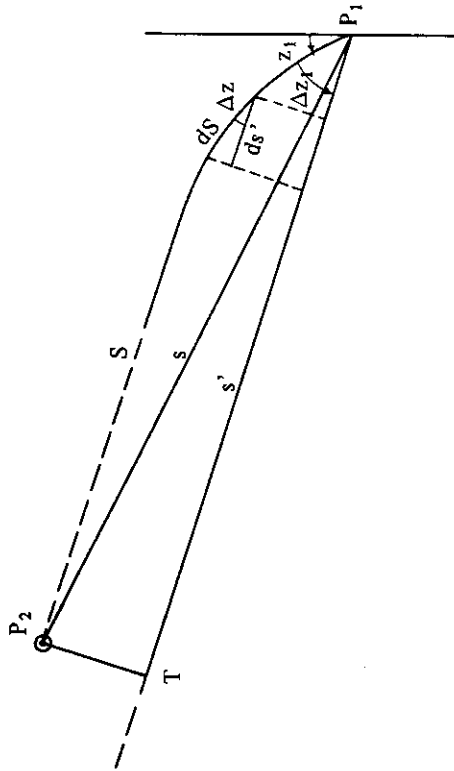


Fig. 4 — Geometric Correction $\delta_g = S - s \approx S - s'$.

Referring to Figure 4, the geometric correction is approximately given by the integral

$$\delta_g = \int_S [1 - \cos(\Delta z)] ds = \frac{1}{2} \int_{r_1}^{r'} (\Delta z)^2 \sec z \, dr$$

where Δz is the astronomical refraction at points along the ray path. Now from (30)

$$\Delta z = \tan z (n-1) - \frac{R}{rg} \tan z \sec^2 z (n-1) T \tag{57}$$

$$\frac{1}{2} (\Delta z)^2 \sec z = \frac{1}{2} \tan^2 z \sec z (n-1)^2 - \frac{R}{rg} \tan^2 z \sec^3 z (n-1)^2 T$$

and from (3) and (3*)

$$\tan^2 z \sec z = \tan^2 z_1 \sec z_1 - \frac{3}{r_1} \sec^5 z_1 (r-r_1)$$

$$\tan^2 z \sec^3 z = \sec^5 z_1$$

approximately, which gives

$$\delta_g = \frac{1}{2} \tan^2 z_1 \sec z_1 \int_{r_1}^{r'} (n-1)^2 \, dr - \frac{3}{2r_1} \sec^5 z_1 \int_{r_1}^{r'} (n-1)^2 (r-r_1) \, dr - \frac{R}{r_1 g} \sec^5 z_1 \int_{r_1}^{r'} (n-1)^2 T \, dr \tag{58}$$

The values of the first two integrals in equation (58) have been determined previously in equations (22), (23) and (52).

Equation (22) also gives the stratospheric component of the third integral in (58):

$$\int_{r_0}^{r'} (n-1)^2 T \, dr = T^0 \int_{r_0}^{r'} (n-1)^2 \, dr = \frac{R}{2g} (n^0 - 1)^2 T^0$$

In the troposphere, we have from (17)

$$\int (n-1)^2 T \, dr = \frac{(n_1-1)^2 T_1}{\beta} \int \left(\frac{T}{T_1}\right)^{2m'+1} dT = \frac{(n_1-1)^2 T_1^2}{\beta (2m'+2)} \left(\frac{T}{T_1}\right)^{2m'+2} + C = -\frac{R}{2g} (n-1)^2 T^2 + C$$

and

$$\int_{r_1}^{r_0} (n-1)^2 T \, dr = \frac{R}{2g} [(n_1-1)^2 T_1^2 - (n^0-1)^2 T^0]$$

The total value of the third integral is consequently

$$\int_{z_1}^{z_1'} (n-1)^2 T dr = \frac{R}{2g} (n_1-1)^2 T_1^2 \quad (59)$$

Substituting the values of the atmospheric integrals into equation (58), we obtain the following formula for the geometric correction :

$$\begin{aligned} \delta_g = \frac{R}{2g} (\sec^3 z_1 - \sec z_1) & \left[\frac{(n_1'-1)^2 T_1' + \frac{1}{2} (R\beta/g)(n^0-1)^2 T^0}{2 + R\beta/g} \right] - \\ & - \frac{3R^2}{4r_1 g^2} \sec^5 z_1 \left[\frac{(n_1'-1)^2 T_1'^2 - (n^0-1)^2 T_1' T^0}{2 + R\beta/g} + \frac{1}{2} (n^0-1)^2 T^0 \right] - \\ & - \frac{R^2}{2r_1 g^2} \sec^5 z_1 (n_1'-1)^2 T_1'^2 \end{aligned} \quad (60)$$

where the primed quantities, as before, refer to values corrected for ground inversion. At a maximum zenith distance $z_1 = 80^\circ$, the geometric correction amounts to about 0.03 metres. It can be readily included in the range correction by modifying appropriate terms in equations (31*).

Corrections for Vapour Pressure.

The amount and distribution of water vapour in the atmosphere varies greatly according to the prevailing conditions of evaporation and condensation. Assuming that the vapour pressure decreases with height in a similar manner as total pressure (16), we can write

$$e = e_1 \left(\frac{T}{T_1} \right)^{-\nu g/(R\beta)} \quad (61)$$

where ν is a numerical coefficient determined from local observations. Equation (61) provides a convenient means for the evaluation of various humidity integrals for the purpose of refraction.

To determine the contribution of humidity to refractivity integral (8), we have from (61)

$$\begin{aligned} \int \left(\frac{e}{T} \right) dr &= \frac{e_1}{\beta T_1} \int \left(\frac{T}{T_1} \right)^{-\nu g/(R\beta)-1} dT = -\frac{R}{\nu g} e + C \\ \int_{z_1}^{z_1'} \left(\frac{e}{T} \right) dr &= \frac{R}{\nu g} e_1 \end{aligned} \quad (62)$$

and

$$\begin{aligned} \int \left(\frac{e}{T^2} \right) dr &= \frac{e_1}{\beta T_1^2} \int \left(\frac{T}{T_1} \right)^{-\nu g/(R\beta)-2} dT = \\ &= -\frac{e_1}{\beta T_1} \left[\frac{1}{\nu g/(R\beta) + 1} \left(\frac{T}{T_1} \right)^{-\nu g/(R\beta)-1} \right] + C = -\left(\frac{R}{\nu g + R\beta} \right) \left(\frac{e}{T} \right) + C \\ & \int_{z_1}^{z_1'} \left(\frac{e}{T^2} \right) dr = \left(\frac{R}{\nu g + R\beta} \right) \left(\frac{e_1}{T_1} \right) \end{aligned} \quad (63)$$

This gives a correction for vapour pressure

$$\delta_w = \left[\frac{(n_0-1)RT_0}{\nu p_0 g} \left(1 - \frac{R}{R_w} \right) - \frac{Rc_w}{\nu g} + \left(\frac{c_w}{\nu g/R + \beta} \right) \left(\frac{1}{T_1} \right) \right] \sec z_1 e_1 \quad (64)$$

which must be added to range correction (5*).

The contribution of humidity to the second term in (5*)

$$\begin{aligned} & - \frac{Rc_w}{\nu T_1 g (\nu g/R + \beta)} \tan^2 z_1 \sec z_1 e_1 \\ & \text{and to geometric correction (58)} \\ & + \frac{Rc_w (n_1-1)}{g (\nu + 1 + 2R\beta/g) T_1} \tan^2 z_1 \sec z_1 e_1 \end{aligned}$$

are further corrections for vapour pressure, both derived on the basis of equations (61) and (7), that might be considered in the range correction. They amount to less than one percent of correction (64) and, in view of the uncertainty in the determination of the latter, may be omitted in practical applications.

At a maximum zenith distance $z_1 = 80^\circ$, and assuming an average value of $\nu = 4$, vapour pressure correction (64) may under extreme conditions exceed 2.5 metres in the radio range, and 0.03 metres in the laser range.

Formulas and Tables for the Computation of Range Correction

1. Standard Formulas.

The atmospheric correction for troposphere and stratosphere in electromagnetic ranging of satellites is given by the standard formula (a) for laser ranging :

$$\Delta s_0 = 0.002357 \sec z (p + 0.06 e - B \tan^2 z) + \delta_L \quad (56a)$$

or (b) for radio ranging :

$$\Delta s_0 = 0.002277 \sec z \left[p + \left(\frac{1255}{T} + 0.05 \right) e - B \tan^2 z \right] + \delta_R \quad (56b)$$

where Δs_0 is the range correction in metres, z is the apparent (radio) zenith distance of the satellite, p is the total barometric pressure in millibars, e is the partial pressure of water vapour in millibars, T is the absolute temperature in degrees Kelvin, and B and δ are correction quantities obtained from Tables X and XI, respectively. In radio ranging, apparent zenith distance z can be determined from true zenith distance Z of the satellite by the formula $z = Z - \Delta z$, where

$$\Delta z'' = \frac{16''.0 \tan Z}{T} \left(p + \frac{4800 e}{T} \right) - 0''.07 (\tan^3 Z + \tan Z) \left(\frac{p}{1000} \right) \quad (57a)$$

is the angle of refraction.

2. Correction for the Effective Wavelength.

Formula (56a) employs a standard wavelength of 0.6943 microns for a ruby laser. For other laser systems, the numerical coefficient of the first term in the formula is obtained from the expression

$$\frac{0.39406 (173.3 + 1/\lambda^2)}{(173.3 - 1/\lambda^2)^2}$$

where λ is the effective wavelength of the system expressed in microns.

3. Correction for Local Latitude and Station Height.

The numerical coefficient of the first term in formulas (56a) and (56b) is to some extent dependant on local latitude and station height. A locally corrected value may be obtained by applying a correction factor

$$1 + 0.0026 \cos 2 \varphi + 0.00028 H$$

Table X.

Standard Values of Correction Factor B

$$B = \frac{R}{r_1 g} \left[\frac{p_1' T_1' - (R\beta/g) p^0 T_0}{1 - R\beta/g} \right]$$

Station Height Above Sea Level	B, mb	Station Height Above Sea Level	B, mb
0 km	1.156	2 km	0.874
0.5 km	1.079	2.5 km	0.813
1 km	1.006	3 km	0.757
1.5 km	0.938	4 km	0.654
2 km	0.874	5 km	0.563

Table XI.

Correction Term δ_L in Metres

$$(\delta_R \approx \delta_L)$$

Apparent Zenith Distance	Station Height Above Sea Level									
	0 km	0.5 km	1 km	1.5 km	2 km	3 km	4 km	5 km		
60° 00'	+0.003	+0.003	+0.002	+0.002	+0.002	+0.002	+0.001	+0.001		
66° 00'	+0.006	+0.006	+0.005	+0.005	+0.004	+0.003	+0.003	+0.002		
70° 00'	+0.012	+0.011	+0.010	+0.009	+0.008	+0.006	+0.005	+0.004		
73° 00'	+0.020	+0.018	+0.017	+0.015	+0.013	+0.011	+0.009	+0.007		
75° 00'	+0.031	+0.028	+0.025	+0.023	+0.021	+0.017	+0.014	+0.011		
76° 00'	+0.039	+0.035	+0.032	+0.029	+0.026	+0.021	+0.017	+0.014		
77° 00'	+0.050	+0.045	+0.041	+0.037	+0.033	+0.027	+0.022	+0.018		
78° 00'	+0.065	+0.059	+0.054	+0.049	+0.044	+0.036	+0.030	+0.024		
78° 30'	+0.075	+0.068	+0.062	+0.056	+0.051	+0.042	+0.034	+0.028		
79° 00'	+0.087	+0.079	+0.072	+0.065	+0.059	+0.049	+0.040	+0.033		
79° 30'	+0.102	+0.093	+0.085	+0.077	+0.070	+0.058	+0.047	+0.039		
79° 45'	+0.111	+0.101	+0.092	+0.083	+0.076	+0.063	+0.052	+0.043		
80° 00'	+0.121	+0.110	+0.100	+0.091	+0.083	+0.068	+0.056	+0.047		

where φ is the latitude, and H is the station height in kilometres above the sea level. The correction factor is equal to unity for sea-level stations in the middle latitudes.

Accuracy of the Determination of Range Correction

The bending of a ray, as observed in the astronomical refraction, is caused by the same physical phenomenon which alters the speed of propagation. The errors which affect the determination of the range correction are therefore closely related to those involved in the determination of astronomical refraction.

In laser ranging, the maximum errors in range correction at $z_1 = 80^\circ$ may be considered to consist of the following :

Departure from hydrostatic equilibrium	1.5 cm	* * 10° et.
Tilt of isopycnic surfaces	2.0 cm	*
Error in formula (56a)	1.0 cm	
Local error in correction factor B	2.0 cm	
Local error in correction term δ_L	0.5 cm	
Total maximum error	3.4 cm	

→ *Common straight line*

From these round figures it can be roughly estimated that for zenith distances not exceeding 80 degrees, the standard error of range correction (56a) will amount to about 1 to 2 centimetres.

The standard error of range correction (56b) for radio ranging is probably about ten times greater. In this case, the most important single source of error is due to local variations in the vertical distribution of humidity which may cause an error of up to one-fifth of the standard correction for vapour pressure.

Reference in Part II

International Association of Geodesy, Resolution No. 1 of the 13th General Assembly, Bulletin Géo-désique, No. 70, p. 390, 1963.

* See No. 106 p 390 For discussion of these

TRANSMISSION OF PLUMB-LINE DEFLECTIONS BY GROUND OBSERVATIONS IN POLARIZED LIGHT

Abstract

The problem of transferring astrogeodetic plumb-line deflections along the net by measuring angle components between the plumb-line, produced by the direction of electric vector of a linearly polarized light wave is discussed. A mathematical solution of the problem is given and a block-diagram of a photo-electric device for measuring angle-components between plumb-lines is considered.

In establishing geodetic nets in mountainous regions the problems involved in the determination of plumb-line deflection at every station of the net take on an especially great importance. The difficulties encountered in obtaining astronomic coordinates at every point by purely astronomic means and, as is often the case, in the absence of the necessary gravimetric data, call for some other ways to be investigated for the solution of the problem [1].

In what follows we describe one of practicable methods to transfer astronomic coordinates or astro-geodetic plumb-line deflections along a geodetic net, point to point, where projections of the angle between plumb-lines into a plane perpendicular to the line of sight are used as measured quantities.

Since under general conditions the atmosphere is not optically active, and reflection angles are small, the polarization plane of plane-polarized light does not show any appreciable bending.

As a result, the above mentioned projected angles can be measured as follows. Let a source of plane-polarized light be set at one of the stations. The plane in which the electric vector oscillations take place must be oriented parallel to the plumb-line at that station, and the beam pointed at station 2. If the axis of the receiving system analyser at station 2 is oriented parallel to the plumb-lines at that point, then the value of the light-beam passing through the analyser will be the measure of the angle between the plumb-lines projected into the plane which is normal to the line joining the two stations.

To obtain the necessary formulae let us denote the astronomic and the geodetic coordinates of 3 stations by φ_1, λ_1 and B_1 and L_1 , respectively, where