

Last time: Čech cohomology

$\mathcal{U} = \{U_\alpha\}_\alpha$ is a good cover of M .

$$\left[\begin{array}{l} \uparrow \text{each } U_{\alpha_0 \dots \alpha_p} \cong U_{\alpha_0} \cap U_{\alpha_1} \cap \dots \cap U_{\alpha_p} \\ \text{are } \cong \text{ball} \end{array} \right]$$

to each $U_{\alpha\beta}$ associate

$$a \quad \mathcal{H} = \text{span} \{ | \underline{\sigma}_{\alpha\beta} \rangle, \sigma_{\alpha\beta} \in A \}$$

↑ labeling group

eg $A = \mathbb{Z}_2$

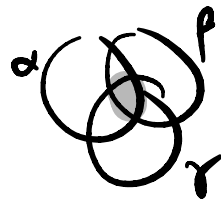
[A 1-cochain $\sigma \in C^1(\mathcal{U}, A)$
specifies a basis state

Take $H = - \sum_{U_\alpha} A_\alpha - \sum_{U_{\alpha\beta\gamma}} B_{\alpha\beta\gamma}$

$$B_{\alpha\beta\gamma} | \{ \sigma \} \rangle = (-1)^{(f\sigma)_{\alpha\beta\gamma}} | \{ \sigma \} \rangle$$

$$f\sigma : C^1 \rightarrow C^2$$

$$\sigma \mapsto (f\sigma)_{\alpha\beta\gamma} = \underline{\sigma_{\alpha\beta} + \sigma_{\beta\gamma} + \sigma_{\gamma\alpha}}$$



$$B|\{\sigma\} = |\{\sigma\}$$

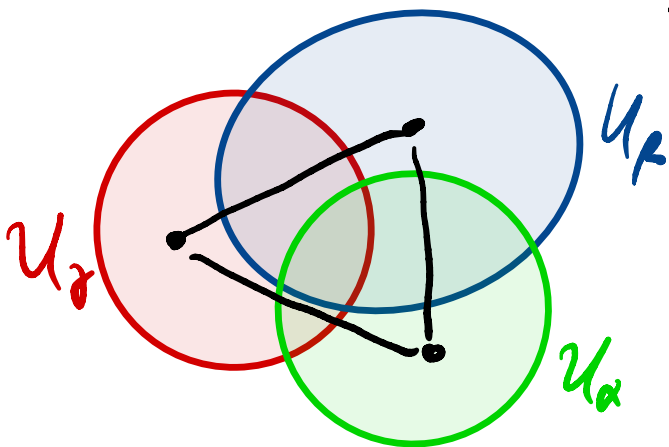
de R

says

σ

is a 1-cocycle

$$\delta\sigma = 0 \pmod{2}$$



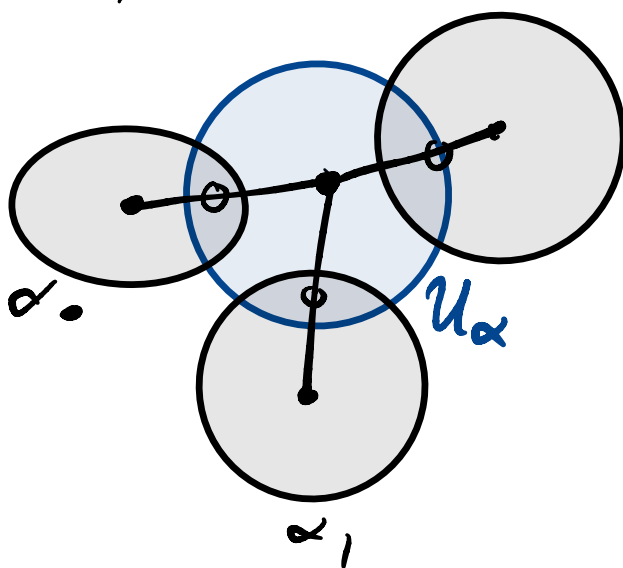
$B_{\text{de R}} \sim \text{play. up}$



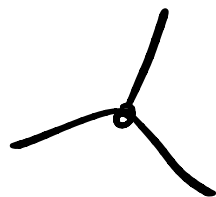
$$\delta(\lambda_\alpha)$$

$$A_\alpha |\{\sigma\} = |\{ \sigma_{\alpha\alpha_0} + \delta_{\alpha\alpha_0} - \delta_{\alpha\alpha_1} \}_{\alpha_0, \alpha_1}$$

$$\lambda_\alpha \equiv \begin{cases} 1 & \text{on } U_\alpha \\ 0 & \text{else} \end{cases}$$



star operator



$$\delta^2 = 0$$

$$\Rightarrow [A, B] = 0$$

same as T.C.

\Rightarrow Čech cohomology \cong cellular homology.

Nice thing abt Čech complex:

we can take weird coefficients

Proof of equivalence $\check{H}_{\mathbb{R}}^p(M) \cong \check{H}^p(\mathcal{U}, \mathbb{R})$

↑
good covers
of M .

$$M \leftarrow \coprod_{\alpha} U_{\alpha} \rightrightarrows \coprod_{\alpha_0 \alpha_1} U_{\alpha_0 \alpha_1} \rightrightarrows \coprod_{\alpha_0 \alpha_1 \alpha_2} U_{\alpha_0 \alpha_1 \alpha_2} \rightrightarrows \dots$$

⇒ generalized Mayer-Vietoris seq on $\check{\Omega}_{\mathbb{R}}^i(M)$

$$0 \rightarrow \check{\Omega}^i(M) \xrightarrow{\text{res}} \bigoplus_{\alpha} \check{\Omega}^i(U_{\alpha}) \xrightarrow{\delta} \bigoplus_{\alpha \beta} \check{\Omega}^i(U_{\alpha \beta}) \xrightarrow{\delta}$$

$$i_{\alpha} : U_{\alpha \beta \gamma} \rightarrow U_{\beta \gamma}$$

$$\bigoplus_{\alpha \beta \gamma} \check{\Omega}^i(U_{\alpha \beta \gamma}) \rightarrow \dots$$

$$\leadsto \delta_{\alpha} : \check{\Omega}^i(U_{\beta \gamma}) \rightarrow \bigoplus_{\alpha} \check{\Omega}^i(U_{\alpha \beta \gamma})$$

$$\delta \equiv \sum_i (-1)^i \delta_{\alpha_i}$$

$$(\delta \omega)_{\alpha_0 \dots \alpha_{p+1}} = \sum_{i=0}^{p+1} (-1)^i \omega_{\alpha_0 \dots \hat{\alpha}_i \dots \alpha_{p+1}} \quad \text{a form on } U_{\alpha_0 \dots \alpha_{p+1}}$$

(Prev. $M=V$ seq. was the special case $M=U \cup V$.)

Facts: ① $\delta^2 = 0$

② seq is exact.

why: make a homotopy operator

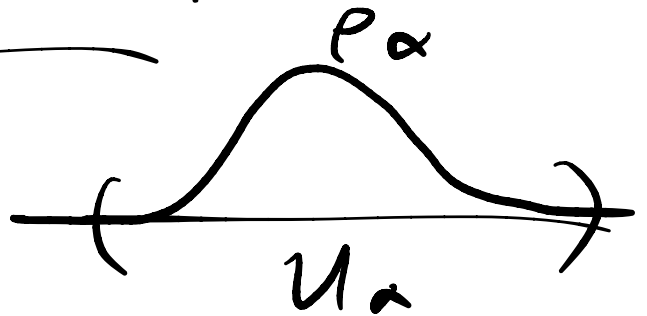
$$K : \bigoplus_{\alpha_0 \dots \alpha_p} \Omega^p(U_{\alpha_0 \dots \alpha_p}) \rightarrow \bigoplus_{\alpha_0 \dots \alpha_{p-1}} \Omega^{p-1}(U_{\alpha_0 \dots \alpha_{p-1}})$$

$$(K\omega)_{\alpha_0 \dots \alpha_{p-1}} = \sum_{\alpha} \rho_{\alpha} \omega_{\alpha \alpha_0 \dots \alpha_{p-1}}$$

where $\sum_{\alpha} \rho_{\alpha} = 1$ partition of unity.

claim

$$\underline{K\delta + \delta K = 1.}$$



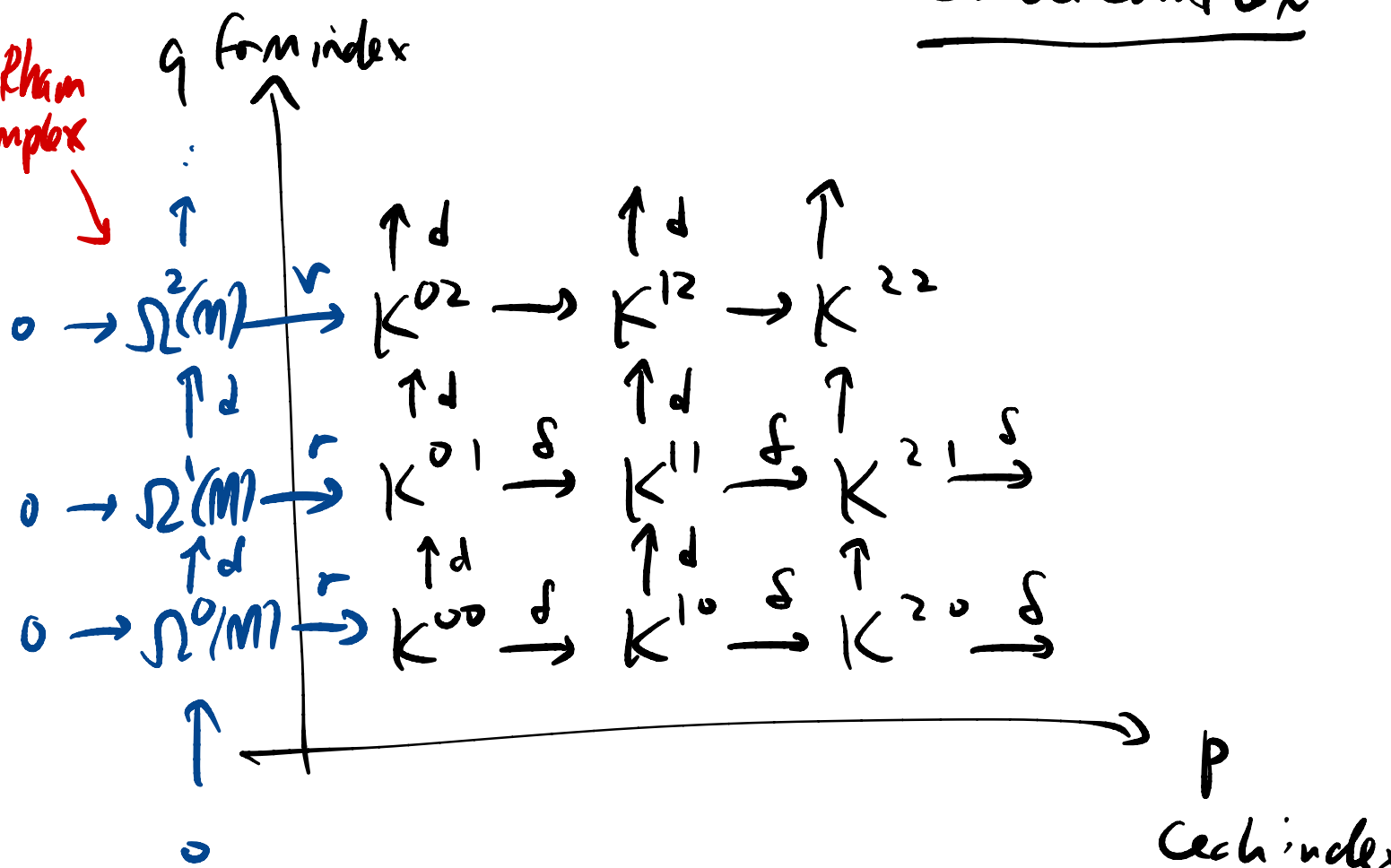
(Note: can't do this on C^0 !)

$$K^{p,q} \equiv C^p(\mathcal{U}, \Omega^q) \equiv \bigoplus_{\alpha_0 \dots \alpha_p} \Omega^q(\mathcal{U}_{\alpha_0 \dots \alpha_p})$$

Cech cochains
w/ values in Ω^q

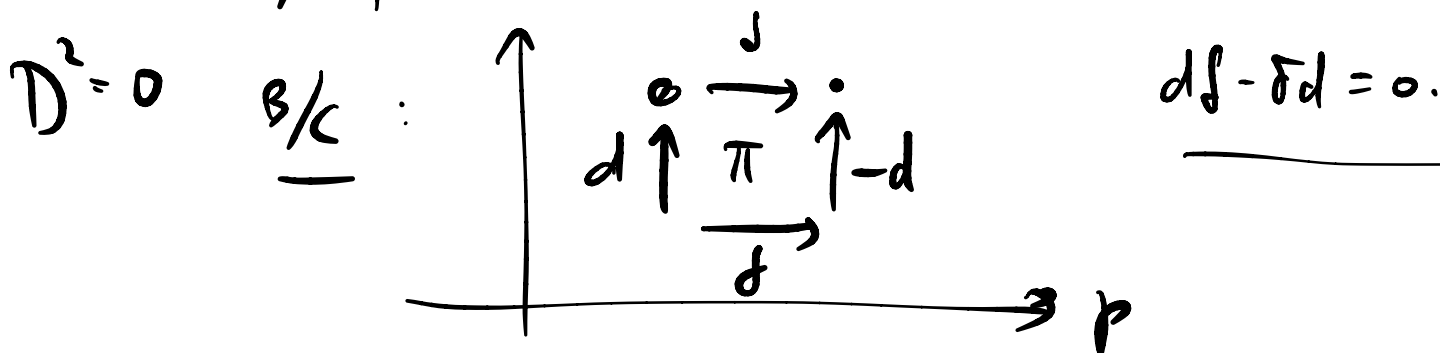
Cech-de Rham
double complex

de Rham
Complex



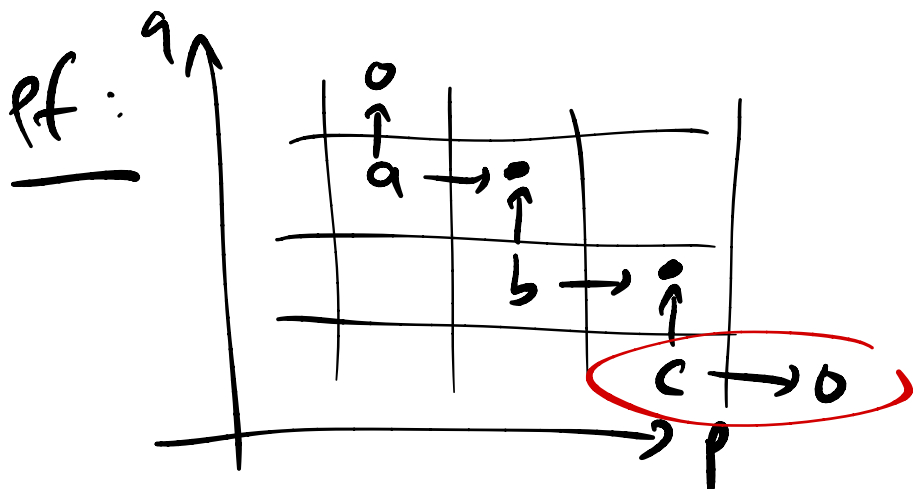
Rows are M.V exact seqs.

New coboundary operator $D \equiv \delta + (-1)^p d$



Lemma: If rows are exact then

$$H^0_D = f + (-1)^p d \cong H^0_{\text{left column}} \cong H^0_{\mathbb{R}}$$

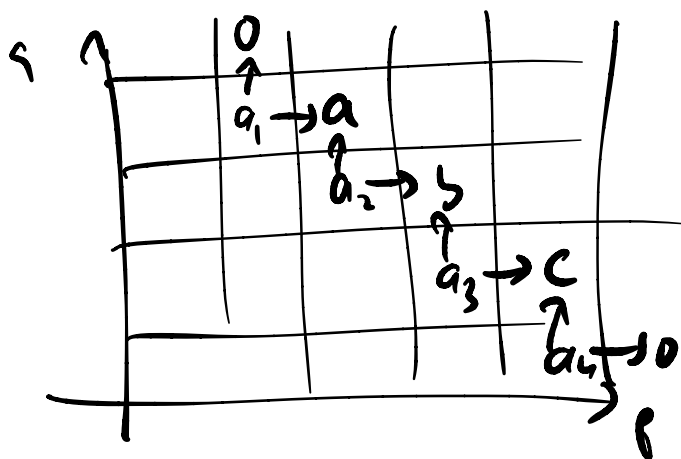


A \mathbb{D} -cocycle:

$$\phi = a + b + c$$

$$\begin{cases} da = 0 \\ \delta a = (-1)^{p+1} db \\ \delta b = (-1)^p dc \\ \delta c = 0 \end{cases}$$

A \mathbb{D} coboundary $\phi = a + b + c$



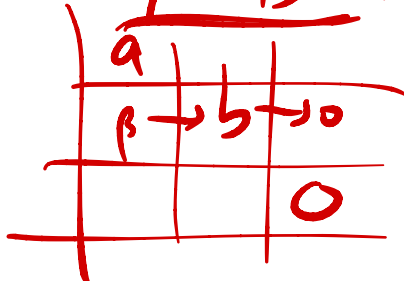
$$a = \delta a_1 + da_2$$

$$b = \delta a_2 + da_3$$

Given ϕ , \mathbb{D} -cocycle

$$\delta c = 0 \Rightarrow c = \delta \alpha$$

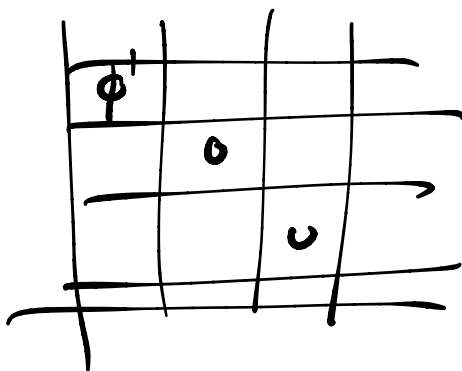
$\phi - \mathbb{D}\alpha$:



$$\delta b = 0$$

$$\Rightarrow b = \delta \beta$$

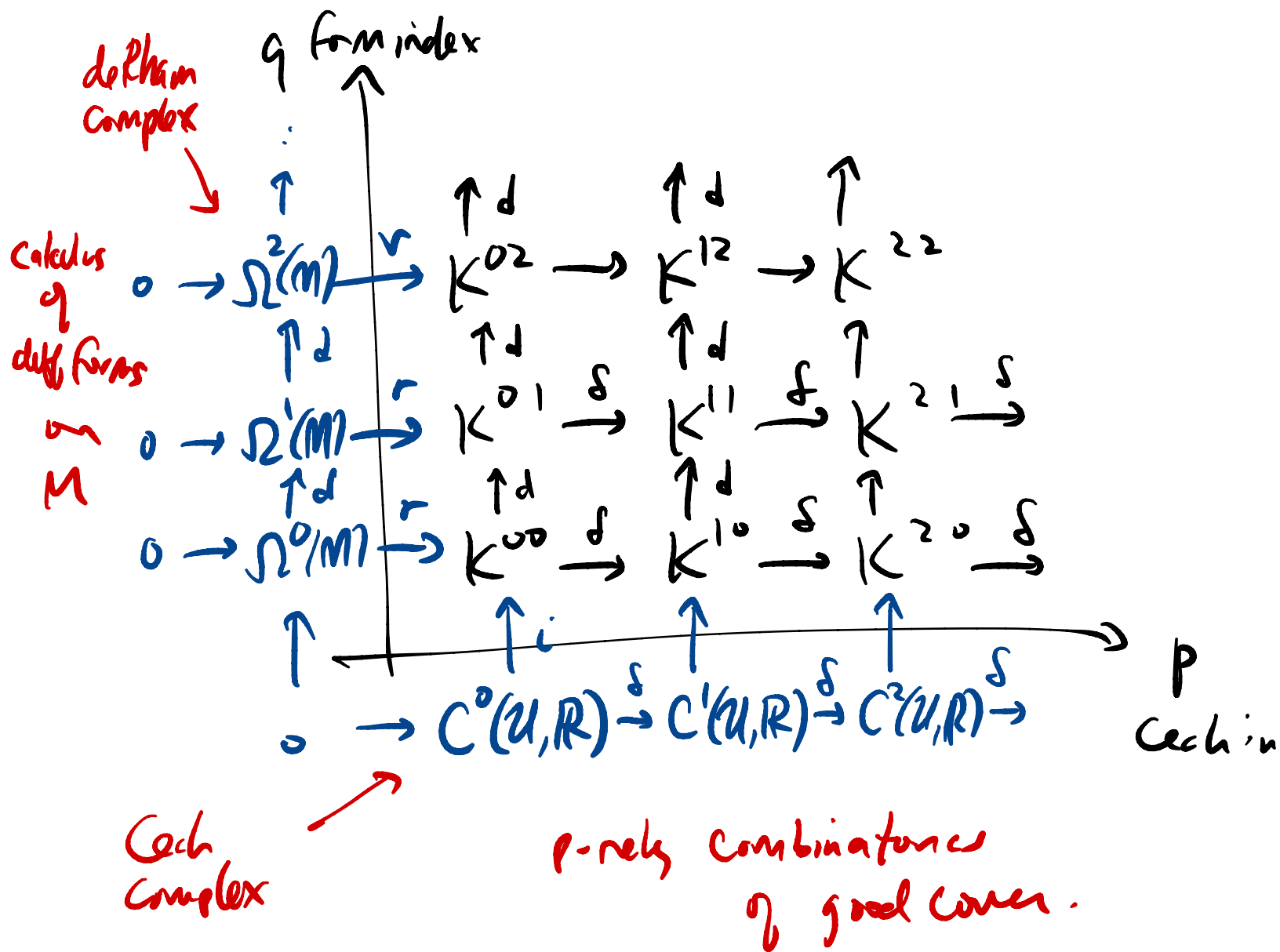
\Rightarrow a representative $\underline{\underline{[\phi]}} = [\underbrace{\phi - D\alpha - D\beta \dots}_{= \phi'}]$



$$d\phi' = 0$$

$$\delta\phi' = 0$$

$\Rightarrow \underline{\underline{\phi'}}$ is a global form on M .

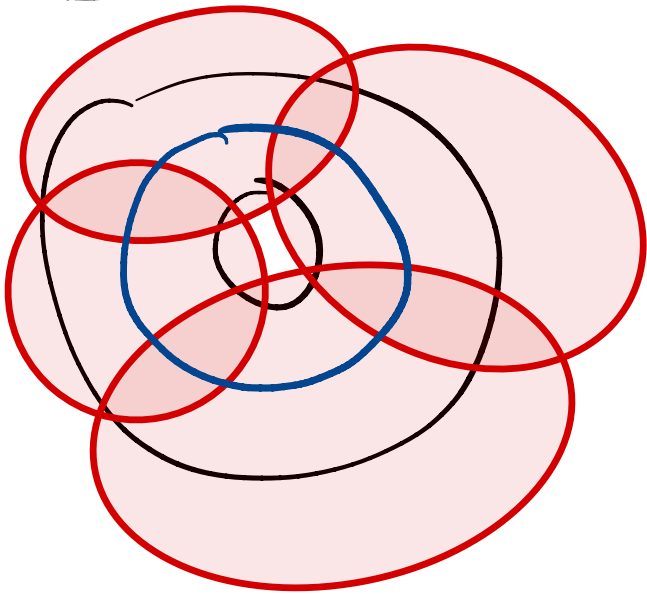


$$H^1_{dR}(M) \simeq H^1_D \simeq H^1_V(\mathcal{U}, \mathbb{R})$$

\Rightarrow each columnology is ind of
 good covers!

2.7 Local reconstructability (of quantum states)

Key feature of T.O.: reduced density matrices
 on patches \equiv marginals



don't determine the global state

$\{ \text{density matrices} \} = \text{convex set.}$

can't see this as the coefficients

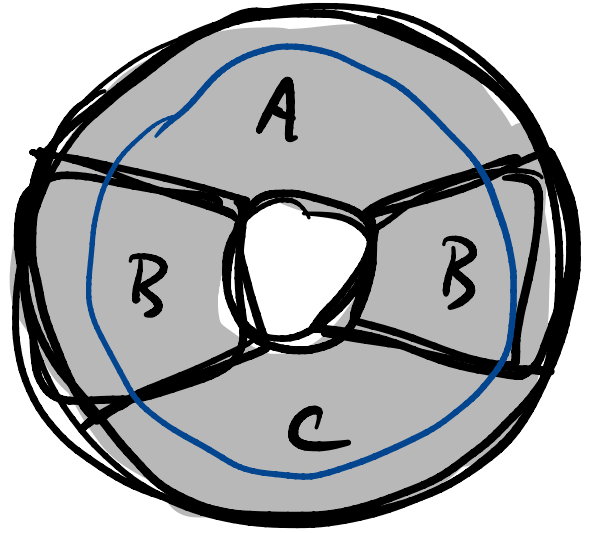
$\rho_1 + \rho_2$ is not
 a density matrix

$(\text{tr } \rho = 1)$

Obstruction to global reconstructibility:

$$I(A:C|B) \equiv \underbrace{S_{AB} + S_{BC} - S_B - S_{ABC}}$$

Q: from P_{AB}, P_{BC}, P_B
can we reconstruct P_{ABC} ?



If $I(A:C|B) = 0$

$P_{ABC} \equiv$ "Markov chain"

can be uniquely reconstructed
acting locally on B.

[Pet 2]

$$\text{If } S_A = |A| - \gamma b_0(\partial A)$$

for some γ ind. of A .

$$\Rightarrow I(A:C|B) = 2\gamma \geq 0 \quad (\text{strong subadditivity})$$

TEE (top entanglement entropy)

(1) Ind. of n deformations of the regions
continuous

(2) Ind of def. of the state keeping
the correlator length \ll size of regions.

(3) It can be related to the anyon data.

$$\gamma = \log \sum_a d_a^2 \quad d_a \equiv \text{quantum dim of anyon } a.$$

$$\left[\dim \mathcal{H}_k \text{ anyons of type } a \stackrel{k \rightarrow \infty}{\sim} d_a^k \right]$$

$$-2\gamma = -I(A:C|B) = \underbrace{(S_{ABC} - S_{BC})}_{\Delta S \text{ from closing } A \text{ on } BC} - \underbrace{(S_{AB} - S_B)}_{\Delta S \text{ from closing } B \text{ on } AB}$$

ΔS from closing A on BC

ΔS from closing B on AB

Pf of 1: let $D = \overline{ABC}$

$\Rightarrow ABCD$ is a pure state.

$\Rightarrow \underline{S_A = S_{BCD}}$.

Under SC:

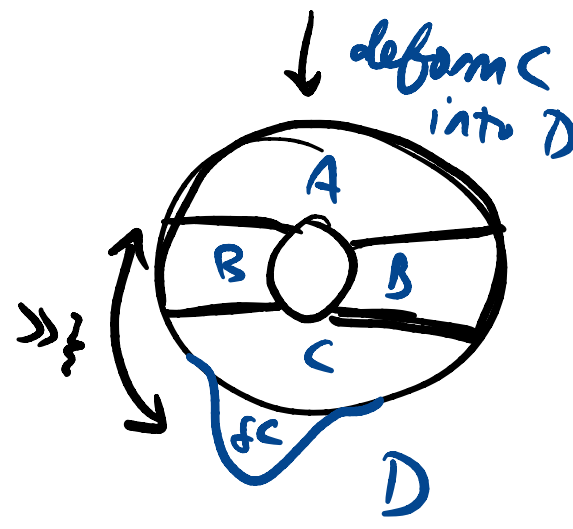
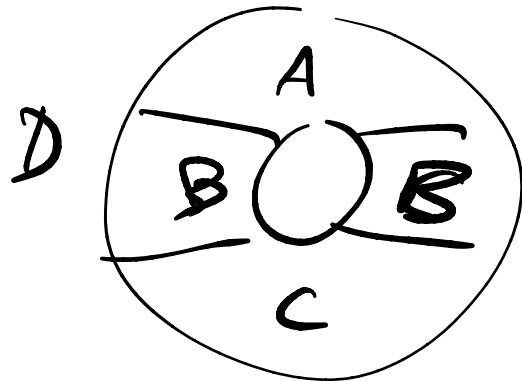
$S_{AB} - S_B$ can't change.

$S_{ABC} - S_{BC} \stackrel{\text{pure}}{=} S_D - S_{AD}$

expect $\Delta_{SC} (S_D - S_{AD}) = 0$

Suppose state is gs. of a local H.

Pf of 2: ΔH localized inside some region
doesn't $S_A, S_B \dots$



if ΔH localized on overlap of A & B:
move A & B so that $\Delta H \in A$.



Fusion rules: $a \times b = \sum_c N_{ab}^c c$

$\Rightarrow a \times a \dots \times a$
k times

$$d_a d_b = \sum_c N_{ab}^c d_c$$
$$= \sum_c (N_a)^c_b d_c$$

d_a are positive eigenvectors of $(N_a)^c_b$.

3. (Quantum Double Model and) Homotopy (Groups)

Def. A homotopy is a family of maps

$$f_t : X \rightarrow Y \quad t \in I = [0, 1]$$

s.t. $F : X \times I \rightarrow Y$ is continuous.

$$(x, t) \mapsto f_t(x)$$

$f_0, f_1 : X \rightarrow Y$ are homotopic. $f_0 \simeq f_1$.

every map
is continuous
in the following

A deformation retraction of X onto $A < X$

is a homotopy from

$$f_0 = \text{id}: X \rightarrow X$$

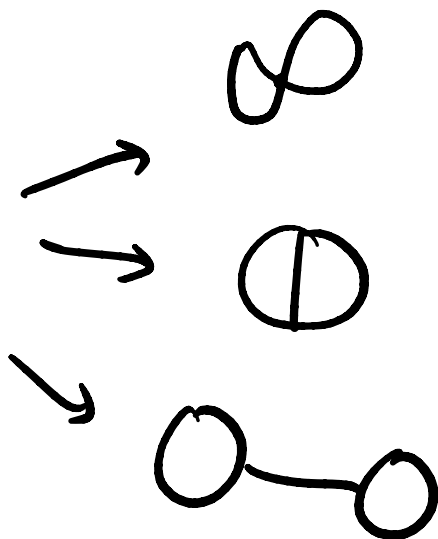
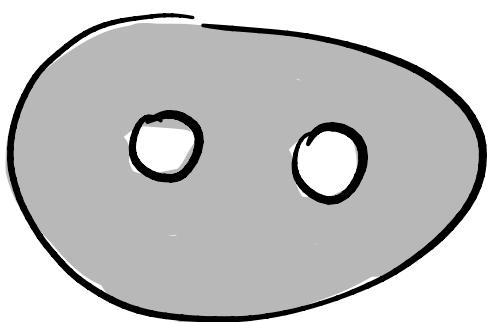
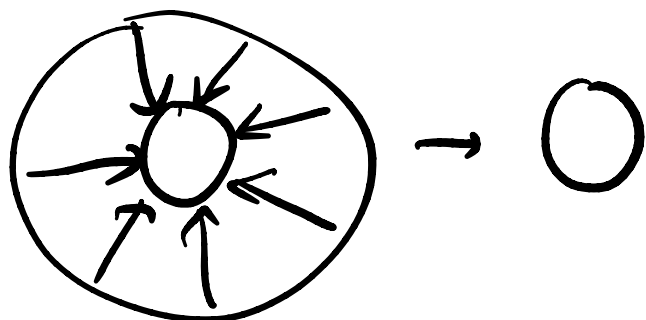
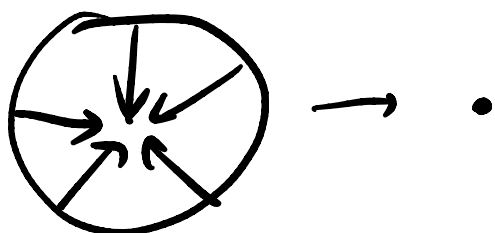
to

$$f_1 = r: X \rightarrow X$$

a retraction

$$r(X) = A \subset X \quad \text{like a projector}$$

$$r^2 = r.$$



\simeq is an eq. relation.

X deformation retracts to B is not.

Def: $X \simeq Y$ X is homotopy equivalent to Y

$\exists f: X \rightarrow Y$ and $g: Y \rightarrow X$ s.t. $f \circ g \simeq \mathbb{1}$
 $g \circ f \simeq \mathbb{1}$.

eg: if X def. retracts to $A \subset X$

via $f: X \times I \rightarrow X$

$r: X \rightarrow A$

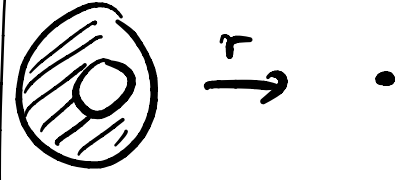
$i: A \hookrightarrow X$

then $r \circ i = \text{id}$ on A

and $i \circ r \simeq \text{id}$ by f .

$\Rightarrow X \simeq A$.

$\bullet \bullet \xrightarrow{r} \bullet$
is not a
def. retract



$X \simeq Y \Leftrightarrow \exists Z$ which deformation retracts to $X \simeq Y$.

If $X \simeq pt$, " X is contractible".

"same"

eg: category
spaces w/ metrics
 $\mathcal{T}_{\text{eucl}(X)}$

• smooth manifolds

• top. mflds

not in of same

isom by

invertible
diffeomorphism

homeomorphism

Homotopy equivalence \neq homeomorphic.

eg: Ball & point (# of dims is a homeomorphism invariant)
eg: $L_{1/7}$ and $L_{2/7}$ $(\cong S^3/\Gamma)$

are homotopy equivalent
but not homeomorphic.

"Combinatorial torsion" distinguishes them.



same homology $\not\Rightarrow$ homotopy equivalence

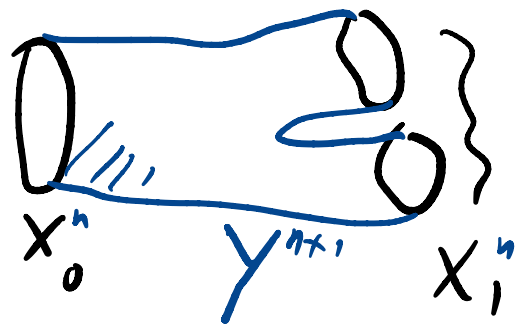
eg $L_{(2/3)}$ and $L_{(1/2)}$

have the same $H^*(L, \mathbb{Z})$

but are not homotopy equivalent.

another equivalence relation:

$X_0 \stackrel{B}{\simeq} X_1$ if $\exists Y$ s.t.
 $X_0 - X_1 = \partial Y$.
a manifold



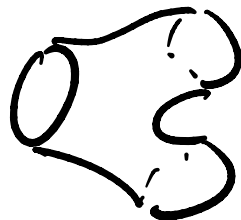
"Bordism"

eg: if $\partial Y = \emptyset$ bordism between \emptyset and \emptyset .

n pts $\stackrel{B}{\simeq} (n+2k)$ pts.

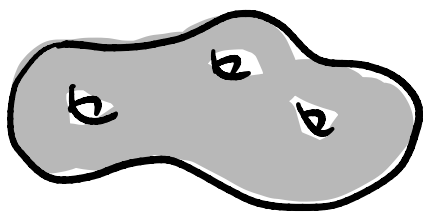
not a mfld.

1d | d:

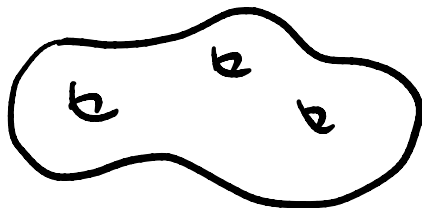


...

1d | 2d:



=



- \emptyset

for any closed R.S.

Boundary doesn't preserve $\chi = 2 - 2g$.

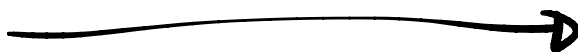
Some surfaces are not the body of anything.

(eg Klein bottle.)

Atiyah-Segal approach to TQFT:

path integral

n-mfld
w/o bdy

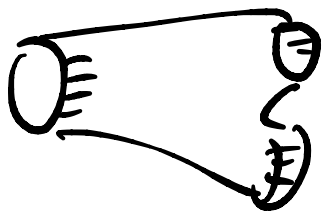


\mathbb{Z}

n-mfld
w/ bdy

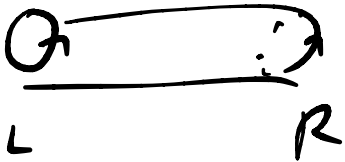


$\Psi[\text{bdy data}] \in \mathcal{Z}(\text{bdy})$



$$\longrightarrow \mathcal{H}_L \otimes \mathcal{H}_R^* \otimes \mathcal{H}_R^*$$

$$\simeq \text{operator: } \mathcal{H}_L \rightarrow \mathcal{H}_R \otimes \mathcal{H}_R$$



$$\longrightarrow \mathcal{H}_L \otimes \mathcal{H}_R^*$$

$$\text{operator: } \mathcal{H}_L \rightarrow \mathcal{H}_R$$

cf: Chen knows this:
[Swingle-Walter ...]

entanglement $\eta \Psi$ (diagram)

g. [irrational CFT data \longrightarrow
"Gaffnian"

non-unitary?
or gapped?

[S. Simon
et al]