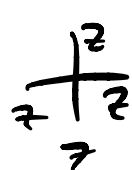
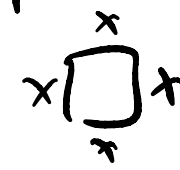


# 2d Toric code, continued

$$H_{TC} = - \sum_j A_j - \sum_p B_p$$

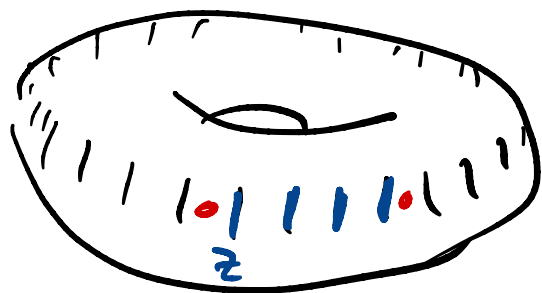




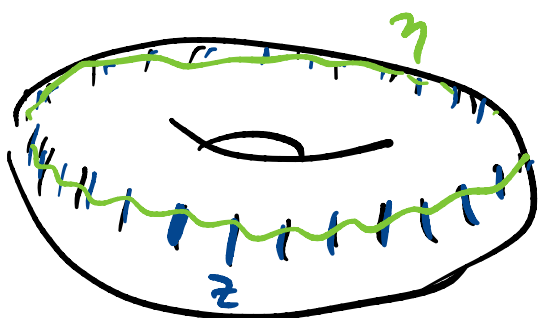


$$W(\sigma) = \prod_{l \in \sigma} X_l$$

~~~~~



$$W(\sigma) V(\eta) = -V(\eta) W(\sigma)$$

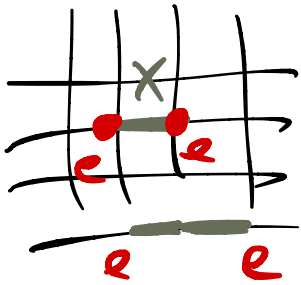


$$V(\eta) = \prod_{l \perp \eta} Z_l$$

~~~~~

GSD  $\Leftrightarrow$  braiding of top. excitations

# Phase Diagram



$$\Delta H = - \sum_e (g X_e + h Z_e)$$

creation/annihilation  
of fermions  
+ kinetic term for  
fermions

... for  
n particles

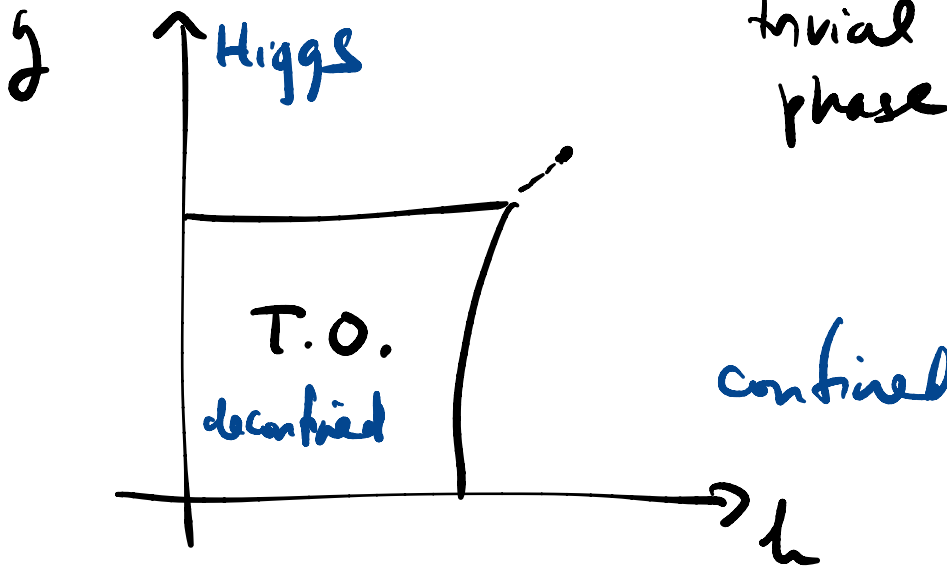
AND

a string tension  
Energy/length  
of electric  
flux strings.

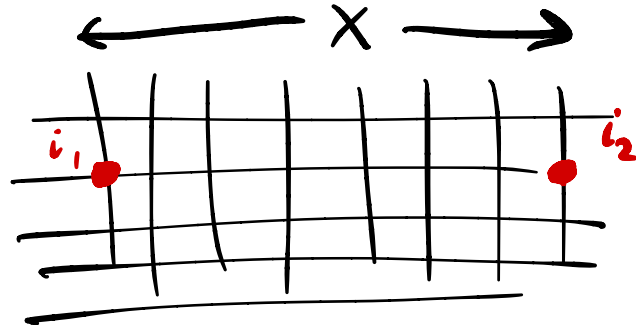
$|S_0\rangle \rightarrow$

Large g:  $\otimes_e |+\rangle_e$  trivial.

Large h:  $\otimes_e |0\rangle_e$  trivial.



Large h: no strings in g.s.



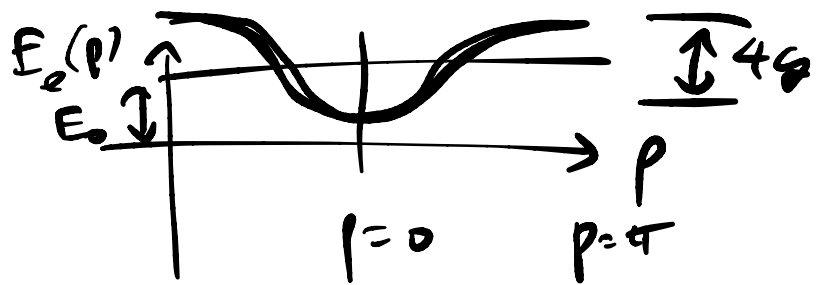
Insert external charge:

demand  $-1 = A_{i_1} = A_{i_2}$ .

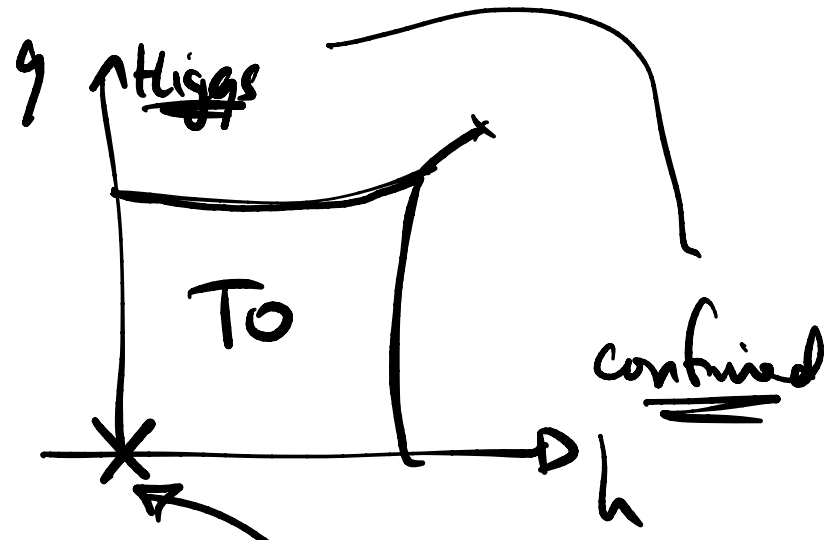
Q: what is  $V(x)$ ?

$$F = - \frac{\partial V}{\partial x}$$

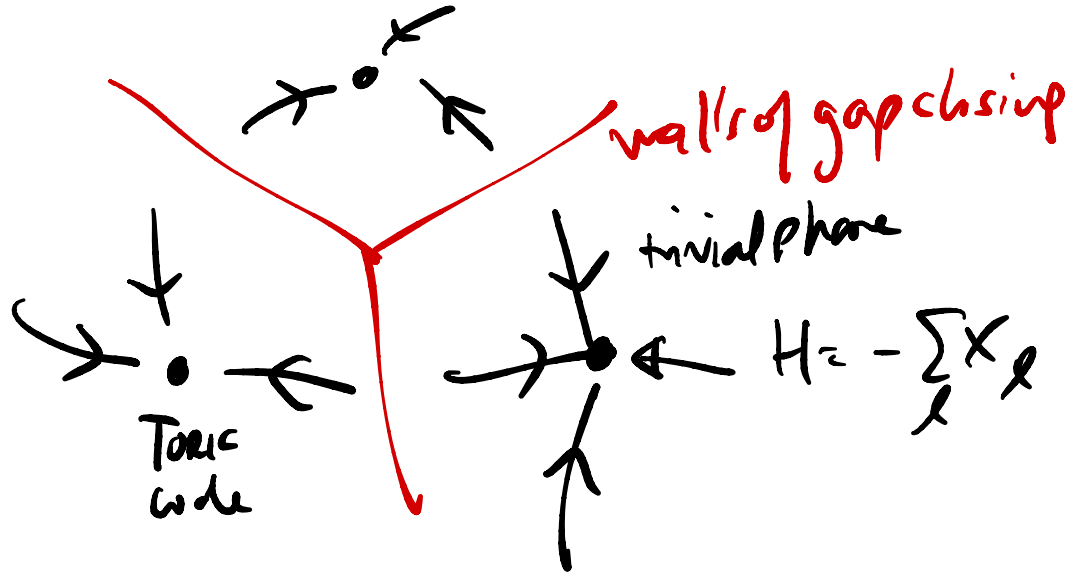




when  $4g \geq 1$   
 $\Rightarrow$  Condense e particle.



RG comment: Topological Code is the RG fixed pt. in the deconfined phase



# Toric code & gauge theory

TC =  $\mathbb{Z}_2$  gauge theory w/ charged bosonic matter.

gauge redundancy:

$$* \left. \begin{array}{l} \phi_i \rightarrow s_i \phi_i \\ X_{ij} \rightarrow s_i X_{ij} s_j \end{array} \right\} \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} \\ \\ \\ \end{array}$$

$s_i = \pm 1$

a field  $\phi_i$   
 $\mathbb{Z}_2$ -valued

$$H = - \sum_p B_p - \sum_l Z_l - g \sum_{ij} \phi_i X_{ij} \phi_j$$

$$- \sum_j A_j (-1)^{n_j}$$

$G_j \equiv \underline{\underline{A_j (-1)^{n_j}}}$  generates the twist \*

$$\left( \begin{array}{l} (-1)^n \phi = -\phi (-1)^n \\ n = 0, 1 = \# \text{ of particles} \end{array} \right)$$

$$\left( \begin{array}{l} G_j = 1 \text{ is} \\ \text{like} \\ \nabla \cdot E = 4\pi\rho \end{array} \right)$$

Set  $\phi = 1 \rightarrow$  T.C.

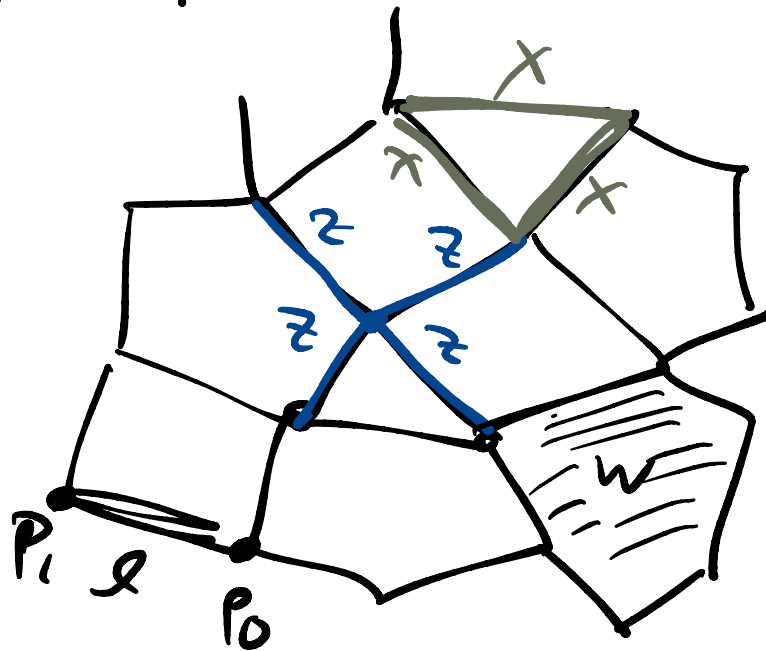
"unitary gauge".

# General graphs (cell complex)

$$B_p = \sum_{\ell \in \partial p} \ell$$

$$A_j = \sum_{\ell \in V(j)} \ell$$

$$[A, B] = 0.$$



$$H = - \sum_i A_i - \sum_p B_p.$$

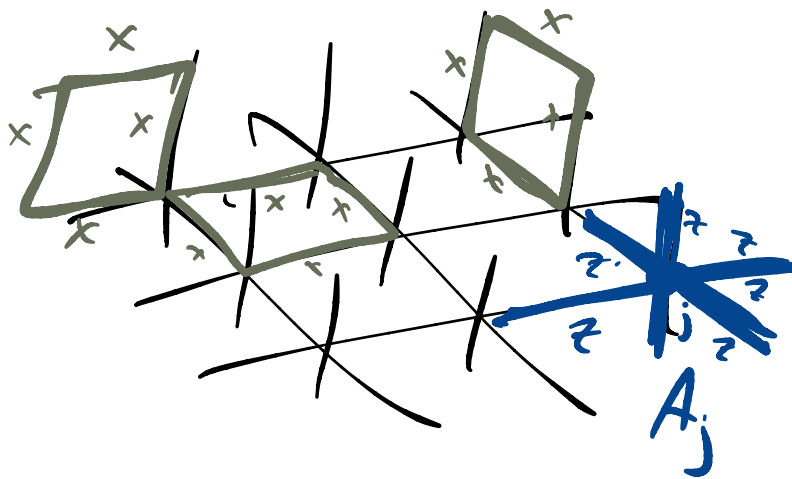
extra data:  $\{ \text{faces} \}_p$  in notion of boundary  $\rightarrow B_p$

$\{ \text{links} \}_\ell$  in  $\partial \ell = \{ p_0, p_1 \}$

$\rightarrow A_j$

$\partial p$   
= collection of links.

eg: <sup>toric code on</sup> 3d<sup>v</sup> cubic lattice

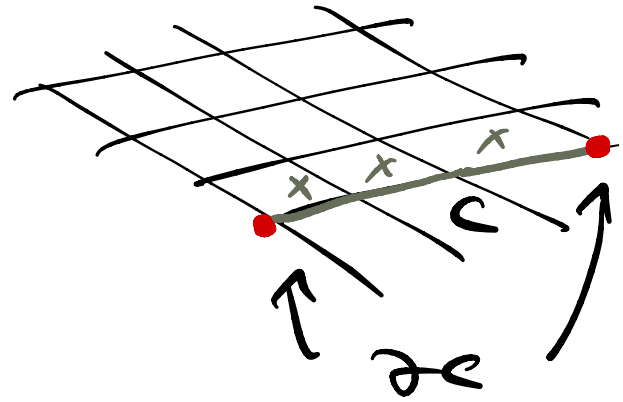


# excitations of 3d toric code:

- violations of  $A_i$  are  $e$ -particles created by

$$W_c = \prod_{l \in c} X_l$$

at  $\partial c$ .

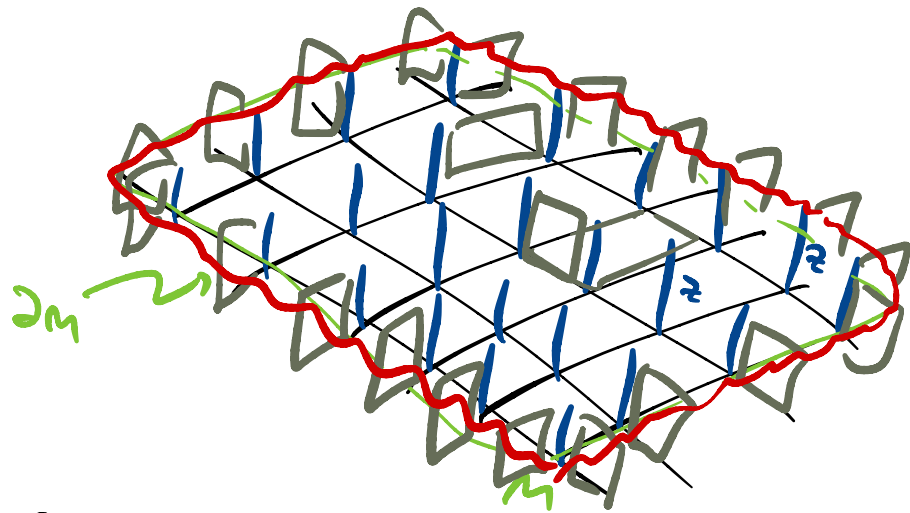


- violations of  $B_p$ .

$$V_M = \prod_{l \perp M} Z_l$$

if  $M$  is closed  $\partial M = 0$

then  $[V_M, H] = 0$ .



$V_M$  creates a string excitation on  $\partial M$ .

The boundary of a boundary is empty

$$\underline{\partial M = 0}$$

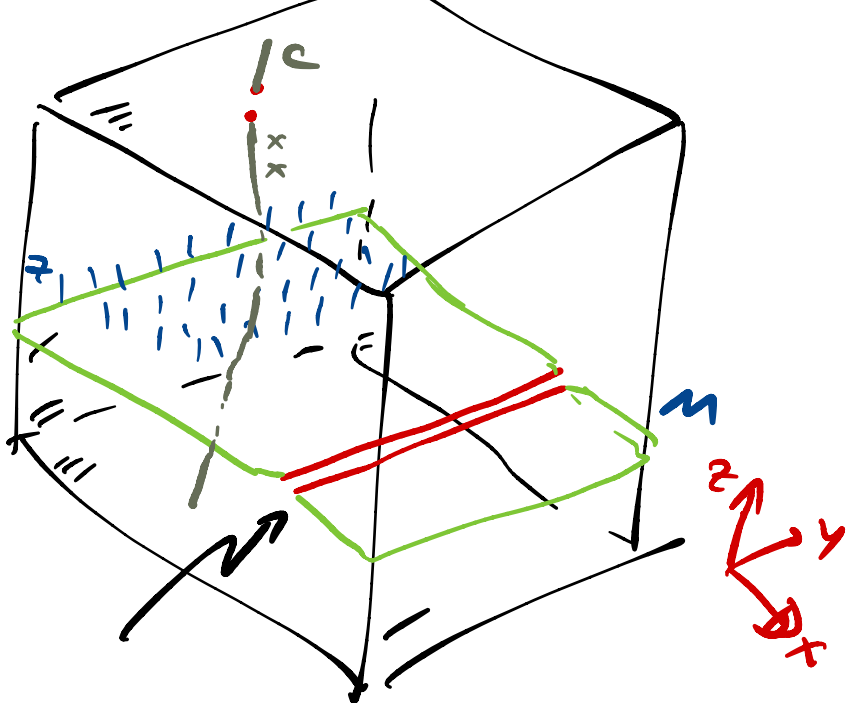
$\Rightarrow$  string is closed.

$$T^3 = S^1 \times S^1 \times S^1$$

$$V_m W_c = -W_c V_m$$

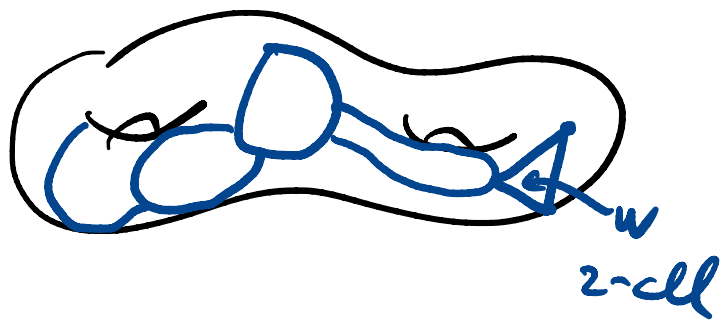
$e$  &  $m$  string  
particle  
braid nontrivially.

$\Leftrightarrow$  nontrivial GSD on  $T^3$ .



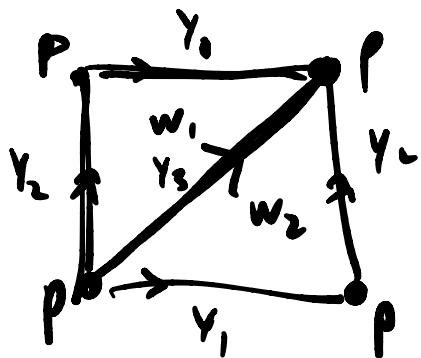
## 1.1 Cell complexes & homology

Take a  $d$ -manifold  $X$  & decomp it into cells.



each is  $\cong$  a ball.

choose an abelian group  $A$ .  
eg  $\mathbb{Z}_2 = \langle 0, 1 \rangle$ .  $1+1=0$ .



2 cells  $w_1, w_2$

1 cells  $y_1, y_2, y_3$

0 cells  $p$

$$\begin{cases} \partial w_1 = y_3 - y_1 - y_2 \\ \partial w_2 = y_1 + y_2 - y_3 \end{cases}$$

$$\begin{cases} \partial y_1 = p - p = 0 = \partial y_2 \\ \quad \quad \quad = \partial y_3 \end{cases}$$

$$\partial p = 0$$



$$\Omega_k \equiv \Omega_k(\Delta, A) \equiv \text{span}_A \{ \sigma, \sigma \in \Delta_k \}$$

$\uparrow$   
 triangulation  
 of  $X$

$\nwarrow$   
 $\Delta_k \equiv \{ k\text{-cells} \}$   
 in  $\Delta$

$C \in \Omega_k \equiv$  formal linear combination  
 of  $k$ -cells  
 $\equiv k$ -chain.

$$C + C' \in \Omega_k$$

orientation:  $\sigma_{\bar{l}} = -\sigma_l$  .  $\bar{l} \equiv$  reverse  
 orientation

Bdy map:  $\partial_k : \Omega_k \rightarrow \Omega_{k-1}$   
linear  $\sigma_k \mapsto (\partial \sigma_k)$

This data is called a cell complex.

The bdy of a bdy is empty  $\partial_{k-1} \circ \partial_k \equiv \partial^2 = 0$ .  
 $\Rightarrow$  A cell complex is a chain complex

A  $\wedge$  chain (  $\partial C = 0$  ) is called a cycle  
 closed ie  $C \in \text{Ker } \partial$ .

$$\partial^2 = 0 \implies \text{Ker } \partial_k \supset \text{Im } \partial_{k+1}$$

$$\begin{array}{ccccccc} \dots & \rightarrow & \Omega_{k+1} & \xrightarrow{\partial_{k+1}} & \Omega_k & \xrightarrow{\partial_k} & \Omega_{k-1} \rightarrow \dots \\ & & C_{k+1} & \xrightarrow{\partial_{k+1}} & \underbrace{\partial_{k+1}(C_{k+1})} & \xrightarrow{\partial_{k+1}} & 0 \end{array}$$

Homology of this chain complex

$$\underline{\underline{H_k(\Delta, A)}} \equiv \frac{\text{Ker}(\partial_k: \Omega_k \rightarrow \Omega_{k-1})}{\text{Im}(\partial_{k+1}: \Omega_{k+1} \rightarrow \Omega_k)}$$

Then:  $\dim H_k$  are top. inits of  $X$ .

Compare to toric code:  $\Omega_1(\Delta, \mathbb{Z}_2)$  label a basis  
 of  $H_{TC}$  on  $\Delta$

$\mathbb{1}$ -chain  $\rightarrow$   $\left. \begin{array}{l} \text{draws blue} \\ \text{line on 1-} \\ \text{cells in } \sigma \end{array} \right\} \in H_{TC}$

$$\text{Ker } \partial_1 \subset \Omega_1 \iff$$

closed string states  
satisfying  $A_j = 1$

$$\text{Im } \partial_2 : \Omega_2 \rightarrow \Omega_1 \iff$$

states madeable  
by the action of  $(\mathcal{B}_p)$ .

$$H_1 \iff$$

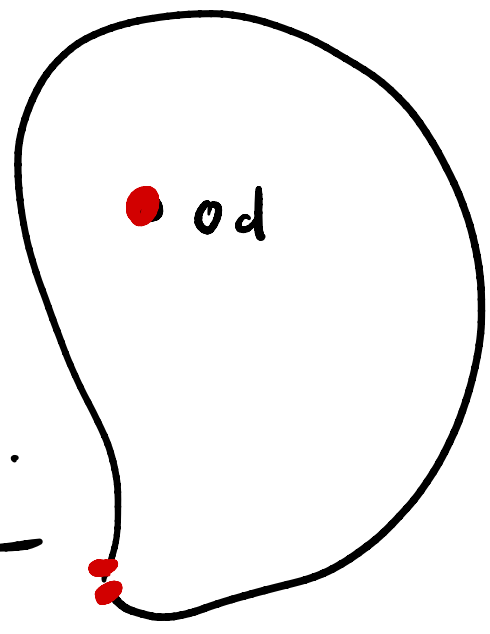
groundstates.

so for  $A = \mathbb{Z}_2$   
and  $p=1$

next: general,  $A$ , general  $p$ .  
abelian

In 2d:

1d worldline  
links w a  
 $d-2$  dim'l object.



$p$ -dim'l worldvolume  
links w a  
 $d-1-p$  dim'l object

In 3d:



eg: in  $d=4$   
braids. string  
w/ string.