

## Physics 215B QFT Winter 2019 Assignment 5

Due 12:30pm Wednesday, February 13, 2019

### 1. Electron-photon scattering at low energy.

Consider the process  $e\gamma \rightarrow e\gamma$  in QED at leading order.

- (a) Draw and evaluate the two diagrams.
- (b) Find  $\frac{1}{4} \sum_{\text{spins/polarizations}} |\mathcal{M}|^2$ .
- (c) Construct the two-body final-state phase space measure in the limit where the photon frequency is  $\omega \ll m$  (the electron mass), in the rest frame of the electron. I suggest the following kinematical variables:

$$p_1 = (\omega, 0, 0, \omega), p_2 = (m, 0, 0, 0), p_4 = (\omega', \omega' \sin \theta, 0, \omega' \cos \theta), p_3 = p_1 + p_2 - p_4 = (E', p')$$

for the incoming photon, incoming electron, outgoing photon and outgoing electron respectively.

- (d) Find the differential cross section  $\frac{d\sigma}{d\cos\theta}$  as a function of  $\omega, \theta, m$ . (The expression can be prettified by using the on-shell condition  $p_3^2 = m^2$  to relate  $\omega'$  to  $\omega, \theta$ . It is named after Klein and Nishina.)
- (e) Show that the limit  $E \ll m$  gives the (Thomson) scattering cross section for classical electromagnetic radiation from a free electron.

### 2. Non-relativistic limits.

#### (a) Coulomb potential.

Derive from QED that the force between non-relativistic electrons is a repulsive  $1/r^2$  force law!

#### (b) Pseudoscalar Yukawa theory.

Consider the theory of a massive Dirac fermion  $\Psi$  and a massive pseudoscalar  $\varphi$  interacting via the term

$$V_5 \equiv g_5 \bar{\Psi} \gamma^5 \Psi \varphi.$$

Convince yourself that this theory is parity invariant.

List the Feynman rules.

Draw and evaluate the diagrams contributing to  $\Psi\Psi \rightarrow \Psi\Psi$  scattering at leading order in  $g_5$ .

Consider the non-relativistic limit,  $m \gg |\vec{p}|$  and find the effective interaction hamiltonian. If you happen to find zero for the leading term, then it's not the leading term.

### 3. Scale invariant quantum mechanics.

Consider the action for one quantum variable  $r$  with  $r > 0$  and

$$S[r] = \int dt \left( \frac{1}{2} m \dot{r}^2 - V(r) \right), \quad V(r) = \frac{\lambda}{r^2}.$$

- (a) Show that the (non-relativistic) mass parameter  $m$  can be eliminated by a multiplicative redefinition of the field  $r$  or of the time  $t$ . As a result, convince yourself that the physics of interest here should only depend on the combination  $m\lambda$ . Show that the coupling  $m\lambda$  is dimensionless:  $[m\lambda] = 0$ .
- (b) Show that this action is *scale invariant*, *i.e.* show that the transformation

$$r(t) \rightarrow s^\alpha \cdot r(st) \tag{1}$$

(for some  $\alpha$  which you must determine), (with  $s \in \mathbb{R}^+$ ) is a symmetry. Find the associated Noether charge  $\mathcal{D}$ . For this last step, it will be useful to note that the infinitesimal version of (1) is ( $s = e^a$ ,  $a \ll 1$ )

$$\delta r(t) = a \left( \alpha + t \frac{d}{dt} \right) r(t).$$

- (c) Find the position-space Hamiltonian  $\mathbf{H}$  governing the dynamics of  $r$ . Show that the Schrödinger equation is Bessel's equation

$$\left( -\frac{\partial_r^2}{2m} + \frac{\lambda}{r^2} \right) \psi_E(r) = E \psi_E(r).$$

Show that the Noether charge associated  $\mathcal{D}$  with scale transformations ( $\equiv$  dilatations) satisfies:  $[\mathcal{D}, \mathbf{H}] = \mathbf{iH}$ . This equation says that the Hamiltonian has a definite scaling dimension, *i.e.* that its scale transformation is  $\delta \mathbf{H} = \mathbf{ia}[\mathcal{D}, \mathbf{H}] = -\mathbf{aH}$ . Note that you should not need to use arcane facts about Bessel functions, only the asymptotic analysis of the equation, in subsequent parts of the problem.

- (d) Describe the behavior of the solutions to this equation as  $r \rightarrow 0$ . [Hint: in this limit you can ignore the RHS. Make a power-law ansatz:  $\psi(r) \sim r^\Delta$  and find  $\Delta$ .]

- (e) What happens if  $2m\lambda < -\frac{1}{4}$ ? It looks like there is a continuum of negative-energy solutions (boundstates). This is another example of a *too-attractive* potential.
- (f) A hermitian operator has orthogonal eigenvectors. We will show next that to make  $\mathbf{H}$  hermitian when  $2m\lambda < -\frac{1}{4}$ , we must impose a constraint on the wavefunctions:

$$(\psi_E^* \partial_r \psi_E - \psi_E \partial_r \psi_E^*)|_{r=0} = 0 \quad (2)$$

There are two useful perspectives on this condition: one is that the LHS is the probability current passing through the point  $r = 0$ .

The other perspective is the following. Consider two eigenfunctions:

$$\mathbf{H}\psi_E = E\psi_E, \quad \mathbf{H}\psi_{E'} = E'\psi_{E'}.$$

Multiply the first equation by  $\psi_{E'}^*$  and integrate; multiply the second by  $\psi_E^*$  and integrate; take the difference. Show that the result is a boundary term which must vanish when  $E = E'$ .

- (g) Show that the condition (2) is empty for  $2m\lambda > -\frac{1}{4}$ . Impose the condition (2) on the eigenfunctions for  $2m\lambda < -\frac{1}{4}$ . Show that the resulting spectrum of boundstates has a *discrete* scale invariance.

[Cultural remark: For some reason I don't know, restricting the Hilbert space in this way is called a *self-adjoint extension*.]<sup>1</sup>

- (h) [Extra credit] Consider instead a particle moving in  $\mathbb{R}^d$  with a central  $1/r^2$  potential,  $r^2 \equiv \vec{x} \cdot \vec{x}$ ,

$$S[\vec{x}] = \int dt \left( \frac{1}{2} m \dot{\vec{x}} \cdot \dot{\vec{x}} - \frac{\lambda}{r^2} \right).$$

Show that the same analysis applies (*e.g.* to the s-wave states) with minor modifications.

[A useful intermediate result is the following representation of (minus) the laplacian in  $\mathbb{R}^d$ :

$$\vec{p}^2 = -\frac{1}{r^{d-1}} \partial_r (r^{d-1} \partial_r) + \frac{\hat{L}^2}{r^2}, \quad \hat{L}^2 \equiv \frac{1}{2} \hat{L}_{ij} \hat{L}_{ij}, \quad L_{ij} = -\mathbf{i} (x_i \partial_j - x_j \partial_i),$$

where  $r^2 \equiv x^i x^i$ . By 's-wave states' I mean those annihilated by  $\hat{L}^2$ .]

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<sup>1</sup>This model has been studied extensively, beginning, I think, with K.M. Case, *Phys Rev* **80** (1950) 797. More recent literature includes Hammer and Swingle, [arXiv:quant-ph/0503074](https://arxiv.org/abs/quant-ph/0503074), *Annals Phys.* 321 (2006) 306-317. The associated Schrödinger equation also arises as the scalar wave equation for a field in anti de Sitter space. A recent paper which discusses connections with the renormalization group in more detail is [this one](#), by S. Paik.