

University of California at San Diego – Department of Physics – Prof. John McGreevy  
**Physics 230 Quantum Phases of Matter, Spr 2024**  
**Assignment 2 – Solutions**

Due 11pm Thursday, April 18, 2024

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1. Anyons in the toric code.

- (a) Show that when acting on a toric code groundstate the operator

$$W_C = \prod_{\ell \in C} X_\ell$$

creates a state which violates only the star operators at the sites in the boundary of  $C$ ,  $\partial C$ , a pair of  $e$ -particles.

Clearly  $W_C$  commutes with  $B_p$  since they are both made of just  $X$ s. If a site  $j$  touches  $C$  in the middle somewhere,  $A_j$  shares two edges with  $W_C$  and therefore  $[A_j, W_C] = 0$ . At the end of the curve is a site which shares only one edge with  $A_j$  and therefore they anticommute. Therefore acting with  $W_C$  changes the  $A_j$ -eigenvalue of the state from 1 to  $-1$ .

- (b) Show that when acting on a toric code groundstate the operator

$$V_{\check{C}} = \prod_{\ell \perp \check{C}} Z_\ell$$

creates a state which violates only the plaquette operators in the boundary of  $\check{C}$ ,  $\partial\check{C}$ .

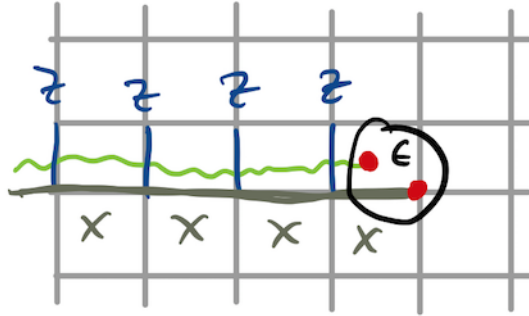
This question is related to the previous by the duality map which takes the lattice to the dual lattice and interchanges  $X \leftrightarrow Z$ .

- (c) Show that a boundstate of an  $e$  particle and an  $m$  particle in the 2d toric code must be a fermion.

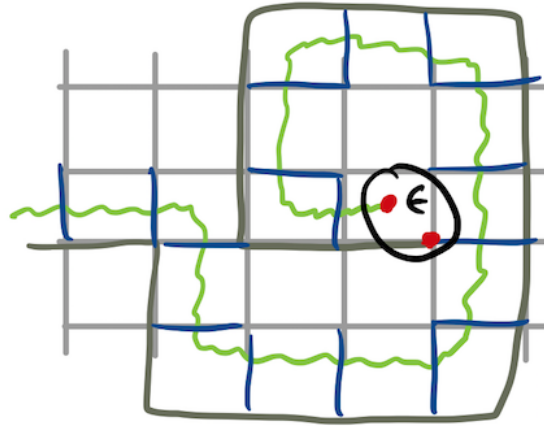
Recall that a particle is a fermion if after rotating it by  $2\pi$  its state picks up a minus sign:

$$| \text{---} \circlearrowleft \rangle = - | \text{---} \bullet \rangle$$

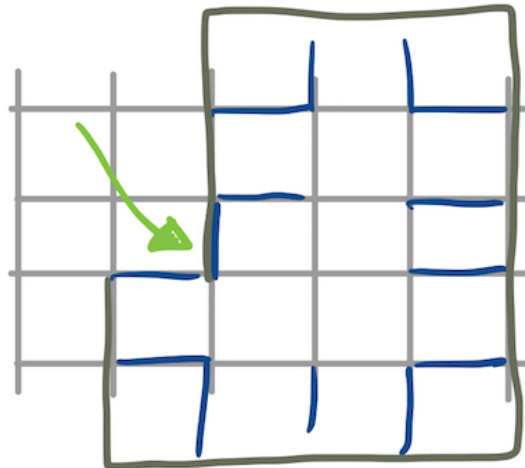
The operator which creates the epsilon particle looks like this, call it  $\mathcal{O}_1$ :



The operator which creates an epsilon particle and rotates it by  $2\pi$  looks like this, call it  $\mathcal{O}_2$ :



The product of the two is this:



(1)

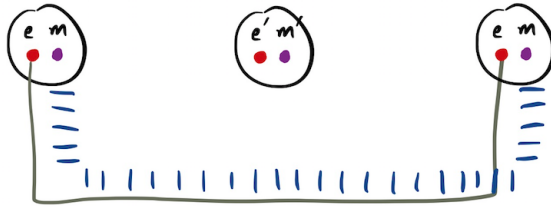
except I need to tell you the order in which the  $X$  and  $Z$  act on the indicated link. Therefore

$$\langle \text{---} \circlearrowleft | \text{---} \cdot \rangle = \langle \text{gs} | \mathcal{O}_2 \mathcal{O}_1 | \text{gs} \rangle .$$

If we can separate the two loops of  $X$ s and  $Z$ s in (??) they will give  $W_C V_{\bar{C}}$  for contractible curves, which gives 1 when acting on the groundstate. But: the  $X$  and  $Z$  on the indicated link need to be moved past each other first. This costs a minus sign and therefore

$$\langle \text{---} \circlearrowleft | \text{---} \cdot \rangle = \langle \text{gs} | \mathcal{O}_2 \mathcal{O}_1 | \text{gs} \rangle = -1 .$$

Alternatively, we can think about *exchange*. Exchanging two particles can be accomplished by first rotating one around the other by a  $\pi$  rotation, and then translating both of them by their separation. As you can see in this figure:



the first step requires the string creating the  $e$  particle to cross that creating the  $m$  particle on an odd number of links. (The second step is innocuous.)

2. **Toric code as  $\mathbb{Z}_2$  gauge theory with matter.** Consider a model with qubits on the links of a lattice (with Pauli operators  $X_\ell, Z_\ell, X_\ell Z_\ell = -Z_\ell X_\ell$ ) and qubits on the sites of the lattice (with Pauli operators  $\sigma_j^x, \sigma_j^z, \sigma_j^x \sigma_j^z = -\sigma_j^z \sigma_j^x$ ).

(a) Show that the operator

$$G_j \equiv A_j \sigma_j^z$$

(where  $A_j$  is the star operator) generates the gauge transformation<sup>1</sup>

$$\sigma_j^x \rightarrow (-1)^{s_j} \sigma_j^x, \quad Z_{ij} \rightarrow (-1)^{s_i} Z_{ij} (-1)^{s_j}, \quad \sigma_j^z \rightarrow \sigma_j^z, \quad X_{ij} \rightarrow X_{ij} \quad (2)$$

(where  $i, j$  are the sites at the ends of the link labelled  $ij$ ). By *generates* here I mean that an operator  $\mathcal{O}$  transforms as

$$\mathcal{O} \rightarrow \mathcal{G}_s^\dagger \mathcal{O} \mathcal{G}_s, \quad \mathcal{G}_s \equiv \prod_j G_j^{s_j}$$

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<sup>1</sup>In the first version of this problem set I had reversed my convention about whether  $A$  is made of  $X$ s or  $Z$ s.

with  $s_j = 0, 1$ .

The relations involving  $Z$  and  $\sigma^z$  follow because they commute with  $G_j$ .  $X_\ell$  transforms under  $G_j$  if the site  $j \in \partial\ell$ .

(b) Show that the Hamiltonian

$$\mathbf{H} = - \sum_j G_j - \sum_p B_p - h \sum_{ij} \sigma_i^x Z_{ij} \sigma_j^x - g \sum_{\langle ij \rangle} X_{ij}$$

is gauge invariant.

$G_j$  commutes with  $G_{j'}$ .  $B_p$  commutes since it is a closed loop of  $Z$ s. The third term is a kinetic term for the  $e$  particles, which is invariant because the transformation of the  $\sigma^x$ s cancels that of the  $Z$ . The last term is invariant because it is made of  $X$ s.

Here we can identify  $\sigma_j^x$  as the operator which creates an  $e$  particle at site  $j$ . And we can identify  $\sigma_j^z = (-1)^{n_j}$  as the parity of the number operator.

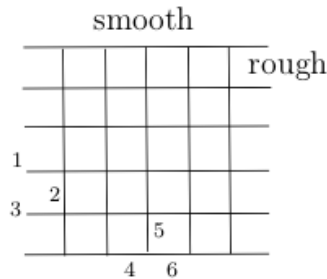
(c) Show that if we set  $\sigma_j^x = 1$  and  $\sigma_j^z = 1$  for all  $j$  we get back the (perturbed) toric code.

Bonus problem: interpret this operation as a choice of gauge in the model where  $G_j = 1$  is imposed as a constraint on physical states.

Clearly if we just erase all the  $\sigma_j^x$ s and  $\sigma_j^z$ s we get back the toric code Hamiltonian.

And clearly we can choose the gauge parameter  $s_j$  in (1) to set  $\sigma^x = 1$ . The tricky part of this is that we can also erase all the appearances of  $\sigma^z$ . This wouldn't make sense on the full Hilbert space  $\otimes \mathcal{H}_2$ , since  $\sigma^z$  and  $\sigma^x$  do not commute. The claim is that on the space of physical states of the gauge theory, which satisfy  $G_j |\text{phys}\rangle = |\text{phys}\rangle$  for all  $j$ , we can do this. It's because on such states, the action of  $\sigma_j^z$  can be replaced by  $A_j$ . If everything is gauge-invariant, any appearance of  $G_j$  can be moved onto the states, and replaced with 1.

3. **A topological qubit.** Consider the toric code on this cell complex:

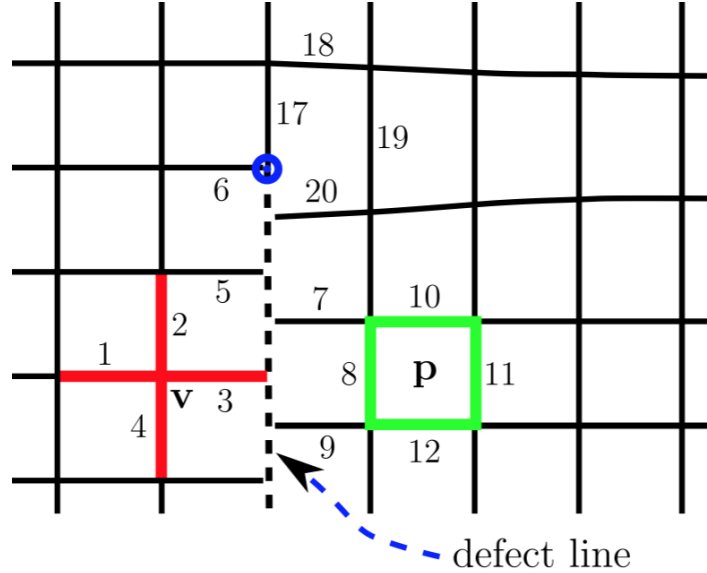


Recall that rough boundary conditions means that plaquette terms get truncated, such as the term  $-X_1X_2X_3$ , while smooth boundary conditions mean that star terms get truncated, such as the term  $-Z_4Z_5Z_6$ .

Show that there is a two-dimensional space of groundstates. A good way to do this is using the algebra of string operators which terminate at the various components of the boundary without creating excitations.

An electric string  $W_C = \prod_{\ell \in C} X_\ell$  can end without creating excitations on the rough boundaries. A magnetic string  $V_{\hat{C}} = \prod_{\ell \perp \hat{C}} Z_\ell$  can end without creating excitations on the smooth boundaries.  $VW = -WV$ , and they commute with the toric code Hamiltonian, so there is a pair of degenerate groundstates.

4. **Duality wall.** [Bonus problem] Show that the following hamiltonian realizes a *duality wall* in the toric code.



What this means is that when crossing the wall, an  $e$  particle turns into an  $m$  particle and vice versa. (More precisely, there is a string operator which transports an  $e$  particle to the wall, and can be completed by a string operator transporting an  $m$  particle away from the wall, without creating any excitations.)

To be more precise about the figure: the dotted line carries no degrees of freedom. The terms in the hamiltonian along the wall are of the form

$$H = \dots - X_2X_5X_3Z_7 - Z_3X_7X_8X_9$$

and there is a term at the end of the wall (the little blue circle) of the form  $-Z_6Y_{17}X_{18}X_{19}X_{20}$ . Show that these terms commute with each other and all the usual star and plaquette terms, such as  $A = Z_1Z_2Z_3Z_4$  and  $B = X_8X_{10}X_{11}X_{12}$ .

What can you say about the end of the duality wall (the little blue circle)?

The figure (and the result) comes from [this paper](#) by Kitaev and Kong.

A string operator which is the usual  $W$  operator made of  $X$ s along the curve on the left (and hence transports an  $e$  particle) can be connected to a  $V$  operator made of a string of  $Z$ s crossing the curve on the right (and hence transports an  $m$  particle) without violating any of the terms in the hamiltonian. And vice-versa.

The endpoint of the duality wall can absorb a fermion (the boundstate of  $e$  and  $m$  particles).