

Physics 212C QM Spring 2020 Assignment 7

Due 12:30pm Monday, May 18, 2020

Please refer to the first homework for submission format and procedures (and replace hw01 by hw07 in the relevant places, of course).

1. **Brain-warmer.** Suppose we have a wavefunction Ψ of N bosons on a thin ring of radius R , governed by a Hamiltonian of the form

$$\mathbf{H} = \sum_i \frac{p_i^2}{2m} + \sum_{i < j} V(|r_{ij}|).$$

Let θ_i be the angular coordinate of the i th particle. Now suppose we make a new state $\Psi' = e^{-i\sum_i \theta_i} \Psi$. Show that

$$\langle \Psi' | \mathbf{H} | \Psi' \rangle = \langle \Psi | \mathbf{H} | \Psi \rangle - \omega_c L + \frac{1}{2} I_{cl} \omega_c^2$$

where L is the expected angular momentum of the state $|\Psi\rangle$, $\omega_c \equiv \frac{\hbar}{mR^2}$, $I_{cl} = NmR^2$.

2. **Density matrix and correlation functions.** Recall the single-particle density matrix in the state $\rho = \sum_s p_s |\Psi_s\rangle\langle\Psi_s|$, as defined in lecture,

$$\rho_1(r, r') \equiv \sum_s p_s \Psi^*(r, r_2, \dots, r_N) \Psi(r', r_2, \dots, r_N).$$

This can be defined for either bosons or fermions.

- (a) Check that ρ_1 is proportional to the two-point correlation function

$$\rho_1(r, r') \propto \text{tr} \rho \psi^\dagger(r) \psi(r') \equiv \langle \psi^\dagger(r) \psi(r') \rangle$$

and find the proportionality constant. Check that it works for both bosons and fermions.

- (b) [Bonus question] Prove that for a fermionic state the eigenvalues of $\rho_1(r, r')$ are between 0 and 1.

[Hint: for fermions, the expectation value of the number operator $\psi^\dagger(r) \psi(r)$ in any state is ≤ 1 .]

3. Topological terms, particle on a ring.

The purpose of this problem is to demonstrate that total derivative terms in the action do affect the physics quantum mechanically.

The euclidean path integral for a particle on a ring with magnetic flux $\theta = \int \vec{B} \cdot d\vec{a}$ through the ring is given by

$$Z = \int [D\phi] e^{-\int_0^\beta d\tau (\frac{m}{2} \dot{\phi}^2 - i \frac{\theta}{2\pi} \dot{\phi})}.$$

Here

$$\phi \equiv \phi + 2\pi \tag{1}$$

is a coordinate on the ring. Because of the identification (1), ϕ need not be a single-valued function of τ – it can wind around the ring. On the other hand, $\dot{\phi}$ is single-valued and periodic and hence has an ordinary Fourier decomposition. This means that we can expand the field as

$$\phi(\tau) = \frac{2\pi}{\beta} Q\tau + \sum_{\ell \in \mathbb{Z} \setminus 0} \phi_\ell e^{i \frac{2\pi}{\beta} \ell \tau}. \tag{2}$$

- (a) Show that the $\dot{\phi}$ term in the action does not affect the classical equations of motion. In this sense, it is a topological term.
- (b) Using the decomposition (2), write the partition function as a sum over topological sectors labelled by the *winding number* $Q \in \mathbb{Z}$ and calculate it explicitly.

[Hint: use the Poisson resummation formula

$$\sum_n f(n) = \sum_l \hat{f}(2\pi l)$$

where $\hat{f}(p) = \int dx e^{-ipx} f(x)$ is the fourier transform of f .]

- (c) Use the result from the previous part to determine the energy spectrum as a function of θ .
- (d) Derive the canonical momentum and Hamiltonian from the action above and verify the spectrum.

4. Excitations and energy of a weakly-interacting bose fluid.

We can see the appearance of the sound mode in the continuum by the following treatment, where we regard the interaction strength as weak.

(a) Write the Hamiltonian

$$\mathbf{H} = \sum_p \mathbf{b}_p^\dagger \mathbf{b}_p \frac{p^2}{2m} + \int d^d r \int d^d r' V(|r - r'|) \mathbf{b}_r^\dagger \mathbf{b}_{r'}^\dagger \mathbf{b}_{r'} \mathbf{b}_r$$

entirely in terms of the momentum-space operators $\mathbf{b}_p, \mathbf{b}_p^\dagger$. After doing this, specialize to the case of contact interactions, *i.e.* $V(r_{ij}) = U\delta^d(r_{ij})$.

(b) Here is the trick. In the BEC state, $\mathbf{b}_0^\dagger \mathbf{b}_0 = N_0 \sim N$. We will approximate such a state as a coherent state for \mathbf{b}_0 with eigenvalue $b_0 \sim \sqrt{N} \gg 1$, a complex number. Since the commutator $[\mathbf{b}_0^\dagger, \mathbf{b}] = 1 \ll N$ we make a small error by this approximation. We will treat all the other creation and annihilation operators $\mathbf{b}_{p \neq 0}$ as operators, but as *small*. That is, $\mathbf{b}_0 \sim b_0 = \mathcal{O}(N)$, $\mathbf{b}_p = \mathcal{O}(N^0)$.

Expand the Hamiltonian keeping only the terms of order N and larger.

To eliminate b_0 , use the fact that the total number of particles is

$$N = |b_0|^2 + \sum_{p \neq 0} \mathbf{b}_p^\dagger \mathbf{b}_p$$

to write it in terms of $\mathbf{b}_{p \neq 0}, \mathbf{b}_{p \neq 0}^\dagger$.

(c) In the previous part you should find a hamiltonian which is quadratic in the $\mathbf{b}_{p \neq 0}, \mathbf{b}_{p \neq 0}^\dagger$. This is an Easy Problem (according to our classification). One way to solve it is to substitute

$$\mathbf{b}_p^\dagger = u_p \mathbf{a}_p + v_p \mathbf{a}_{-p}^\dagger, \quad \mathbf{b}_p = u_p \mathbf{a}_p^\dagger + v_p \mathbf{a}_{-p}$$

where the coefficients u_p, v_p are to be determined and can be assumed to be real, and even functions of p .

Show that demanding that $\mathbf{a}_p, \mathbf{a}_p^\dagger$ satisfy canonical boson commutation relations implies $u_p^2 - v_p^2 = 1$. Find \mathbf{a}_p in terms of $\mathbf{b}_p, \mathbf{b}_p^\dagger$.

(d) To determine u_p, v_p (or α_p in $u_p = \cosh \alpha_p$) demand that the off-diagonal terms drop out of the Hamiltonian

$$\mathbf{H} = \sum_{p \neq 0} \epsilon_p \mathbf{a}_p^\dagger \mathbf{a}_p + C.$$

[Hint: to solve this condition, I recommend making the substitution $u_p = \frac{1}{\sqrt{A_p^2 - 1}}$, and solving for A_p .]

Find the resulting energy spectrum ϵ_p and [bonus problem] the c-number term C .

- (e) Expand ϵ_p about $p = 0$ and find a phonon mode. Determine the speed of sound. What does ϵ_p do at large p ?
- (f) [bonus problem] How does the groundstate energy density E/V depend on the density $n = N/V$ and the scattering length $a \equiv \frac{mU}{4\pi}$?