University of California at San Diego - Department of Physics - Prof. John McGreevy

# Quantum Field Theory C (215C) Spring 2015 Assignment 1 

Posted March 30, 2015
Due 11am, Thursday, April 9, 2015

Please remember to put your name at the top of your homework.

## Problem Set 1

1. Scale invariant quantum mechanics

Consider the action for one quantum variable $r$ with $r>0$ and

$$
S[r]=\int d t\left(\frac{1}{2} m \dot{r}^{2}-V(r)\right), \quad V(r)=\frac{\lambda}{r^{2}} .
$$

(a) Show that the (non-relativistic) mass parameter $m$ can be eliminated by a multiplicative redefinition of the field $r$ or of the time $t$. As a result, convince yourself that the physics of interest here should only depend on the combination $m \lambda$. Show that the coupling $m \lambda$ is dimensionless: $[m \lambda]=0$.
(b) Show that this action is scale invariant, i.e. show that the transformation

$$
\begin{equation*}
r(t) \rightarrow s^{\alpha} \cdot r(s t) \tag{1}
\end{equation*}
$$

(for some $\alpha$ which you must determine), (with $s \in \mathbb{R}^{+}$) is a symmetry. Find the associated Noether charge $\mathcal{D}$. For this last step, it will be useful to note that the infinitesimal version of (1) is ( $s=e^{a}, a \ll 1$ )

$$
\delta r(t)=a\left(\alpha+t \frac{d}{d t}\right) r(t)
$$

(c) Describe the behavior of the solutions to this equation as $r \rightarrow 0$. [Hint: in this limit you can ignore the RHS. Make a power-law ansatz: $\psi(r) \sim r^{\Delta}$ and find $\Delta$.]
(d) What happens if $2 m \lambda<-\frac{1}{4}$ ? It looks like there is a continuum of negative-energy solutions (boundstates). This is another example of a too-attractive potential.
(e) A hermitian operator has orthogonal eigenvectors. We will show next that to make $\mathbf{H}$ hermitian when $2 m \lambda<-\frac{1}{4}$, we must impose a constraint on the wavefunctions:

$$
\begin{equation*}
\left.\left(\psi_{E}^{\star} \partial_{r} \psi_{E}-\psi_{E} \partial_{r} \psi_{E}^{\star}\right)\right|_{r=0}=0 \tag{2}
\end{equation*}
$$

There are two useful perspectives on this condition: one is that the LHS is the probability current passing through the point $r=0$.

The other perspective is the following. Consider two eigenfunctions:

$$
\mathbf{H} \psi_{E}=E \psi_{E}, \quad \mathbf{H} \psi_{E^{\prime}}=E^{\prime} \psi_{E^{\prime}}
$$

Multiply the first equation by $\psi_{E^{\prime}}^{\star}$ and integrate; multiply the second by $\psi_{E}^{\star}$ and integrate; take the difference. Show that the result is a boundary term which must vanish when $E=E^{\prime}$.
(f) Show that the condition (2) is empty for $2 m \lambda>-\frac{1}{4}$. Impose the condition (2) on the eigenfunctions for $2 m \lambda<-\frac{1}{4}$. Show that the resulting spectrum of boundstates has a discrete scale invariance.
[Cultural remark: For some silly reason, restricting the Hilbert space in this way is called a self-adjoint extension.] ${ }^{1}$
(g) [Extra credit] Consider instead a particle moving in $\mathbb{R}^{d}$ with a central $1 / r^{2}$ potential, $r^{2} \equiv \vec{x} \cdot \vec{x}$,

$$
S[\vec{x}]=\int d t\left(\frac{1}{2} m \dot{\vec{x}} \cdot \dot{\vec{x}}-\frac{\lambda}{r^{2}}\right) .
$$

Show that the same analysis applies (e.g. to the s-wave states) with minor modifications.
[A useful intermediate result is the following representation of (minus) the laplacian in $\mathbb{R}^{d}$ :

$$
\vec{p}^{2}=-\frac{1}{r^{d-1}} \partial_{r}\left(r^{d-1} \partial_{r}\right)+\frac{\hat{L}^{2}}{r^{2}}, \quad \hat{L}^{2} \equiv \frac{1}{2} \hat{L}_{i j} \hat{L}_{i j}, \quad L_{i j}=-\mathbf{i}\left(x_{i} \partial_{j}-x_{j} \partial_{i}\right),
$$

where $r^{2} \equiv x^{i} x^{i}$. By 's-wave states' I mean those annihilated by $\hat{L}^{2}$.]

## 2. Gaussian integrals are your friend

(a) Show that

$$
\int_{-\infty}^{\infty} d x e^{-\frac{1}{2} a x^{2}+j x}=\sqrt{\frac{2 \pi}{a}} e^{\frac{j^{2}}{2 a}}
$$

(b) Consider a collection of variables $x_{i}, i=1 . . N$ and a hermitian matrix $a_{i j}$. Show that

$$
\int \prod_{i=1}^{N} d x_{i} e^{-\frac{1}{2} x_{i} a_{i j} x_{j}+J^{i} x_{i}}=\frac{(2 \pi)^{N / 2}}{\sqrt{\operatorname{det} a}} e^{\frac{1}{2} J^{i} a_{i j}^{-1} J^{j}}
$$

(Summation convention in effect, as always.)
[Hint: change into variables to diagonalize $a$. $\operatorname{det} a=\prod a_{i}$, where $a_{i}$ are the eigenvalues of $a$.]

[^0](c) Consider a Gaussian field $X$, governed by the (quadratic) euclidean action
$$
S[x]=\int d t \frac{1}{2}\left(\dot{X}^{2}+\Omega^{2} X^{2}\right) .
$$

Show that

$$
\left\langle e^{-\int d s J(s) X(s)}\right\rangle_{X}=\mathcal{N} e^{+\frac{1}{2} \int d s d t J(s) G(s, t) J(t)}
$$

where $G$ is the (Feynman) Green's function for $X$, satisfying:

$$
\left(-\partial_{s}^{2}+\Omega^{2}\right) G(s, t)=\delta(s-t)
$$

Here $\mathcal{N}$ is a normalization factor which is independent of $J$. Note the similarity with the previous problem, under the replacement

$$
a=-\partial_{s}^{2}+\Omega^{2}, \quad a^{-1}=G
$$

## 3. An application of effective field theory in quantum mechanics.

[I learned this example from Z. Komargodski.]
Consider a model of two canonical quantum variables $\left(\left[\mathbf{x}, \mathbf{p}_{x}\right]=\mathbf{i}=\left[\mathbf{y}, \mathbf{p}_{y}\right], 0=\right.$ $\left[\mathbf{x}, \mathbf{p}_{y}\right]=[\mathbf{x}, \mathbf{y}]$, etc) with Hamiltonian

$$
\mathbf{H}=\mathbf{p}_{x}^{2}+\mathbf{p}_{y}^{2}+\lambda \mathbf{x}^{2} \mathbf{y}^{2}
$$

(This is similar to the degenerate limit of the model studied in lecture with two QM variables where both natural frequencies are taken to zero.)
(a) Based on a semiclassical analysis, would you think that the spectrum is discrete?
(b) Study large, fixed $x$ near $y=0$. We will treat $x$ as the slow (= low-energy) variable, while $y$ gets a large restoring force from the background $x$ value. Solve the $y$ dynamics, and find the groundstate energy as a function of $x$ :

$$
V_{\text {eff }}(x)=E_{\text {g.s. of } y}(x)
$$

(c) The result is not analytic in $x$ at $x=0$. Why?
(d) Is the spectrum of the resulting 1 d model with

$$
\mathbf{H}_{\mathrm{eff}}=\mathbf{p}_{x}^{2}+V_{\mathrm{eff}}(\mathbf{x})
$$

discrete? Is this description valid in the regime which matters for the semiclassical analysis?
[Bonus: determine the spectrum $\mathbf{H}_{\text {eff }}$.]


[^0]:    ${ }^{1}$ This model has been studied extensively, beginning, I think, with K.M. Case, Phys Rev 80 (1950) 797. More recent literature includes Hammer and Swingle, arXiv:quant-ph/0503074, Annals Phys. 321 (2006) 306-317. The associated Schrödinger equation also arises as the scalar wave equation for a field in anti de Sitter space.

