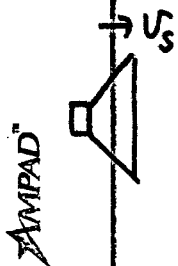


Lect 18 Doppler effect, 4-vectors

§ non-relativistic sound wave Doppler effect. (sound velocity ω)

• source moving velocity v_s .



The two successive sound pulse generated at source at time interval

$$\Delta T = \frac{1}{\nu_0}$$

ⓐ

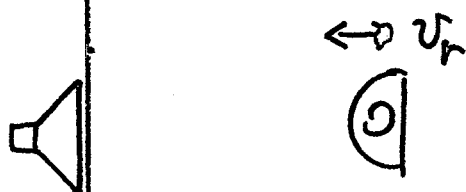
receiver for the receiver to receive them is

$$\Delta T - \frac{v_s \Delta T}{\omega} = \Delta T \left(1 - \frac{v_s}{\omega}\right)$$

$$\Rightarrow \nu_D = \frac{1}{\Delta T \left(1 - \frac{v_s}{\omega}\right)} = \frac{\nu_0}{1 - \frac{v_s}{\omega}}$$

↑
Source

• receiver moving:



The time interval for the receiver to receive two successive pulse @ $\Delta T'$

$$\Delta T = \frac{\Delta T' v_r}{\omega} = \Delta T'$$

$$\Rightarrow \Delta T' = \Delta T / \left(1 + \frac{v_r}{\omega}\right) \Rightarrow \nu_D = \nu_0 \left(1 + \frac{v_r}{\omega}\right)$$

↑
receiver

⇒ Combine together

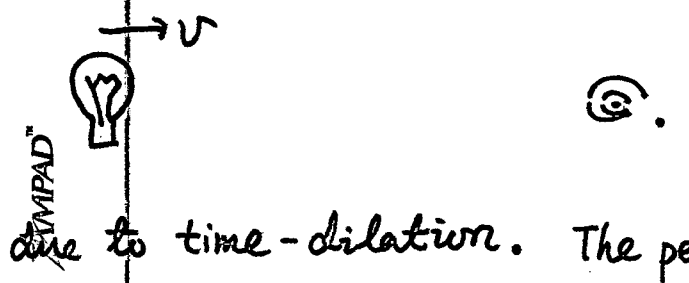
$$\nu_D = \frac{1 + \frac{v_r}{\omega}}{1 - \frac{v_s}{\omega}} \nu$$

receiver/source motion are non-equivalent

source motion can has singularity.

§ Relativistic Doppler effect for light

A light source flashes at $\tau_0 = \frac{1}{\nu_0}$ in its own ^{rest} frame. Now the light source is moving at v toward the receiver, which is at rest.



In the receiver's frame, the period of light pulses is $\tau = \gamma \tau_0$

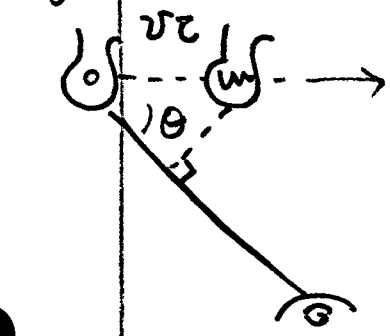
due to time-dilatation. The period at the receiver

$$\tau - \frac{v}{c} \tau = \gamma \tau_0 (1 - \frac{v}{c})$$

$$\Rightarrow \text{receiver's frequency } f = \frac{1}{\tau} = \frac{1}{\tau_0} \frac{\sqrt{1 - \frac{v^2}{c^2}}}{(1 - \frac{v}{c})} = f_0 \sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}}$$

Only the relative motion between the light source and the receiver's matters, because there's no 3rd frame of media like in the case of sound wave.

* if the observer is not at the direction of motion of the light source



Similarly, the propagation distance between two successive pulses is shortened by $\frac{v \tau \cos \theta}{c}$

$$\Rightarrow \tau - \frac{v \tau \cos \theta}{c} = \gamma \tau_0 (1 - \frac{v}{c} \cos \theta), \text{ transverse Doppler}$$

\Rightarrow receiver's frequency

$$f = \frac{1}{\tau} = f_0 \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 - \frac{v}{c} \cos \theta}$$

even if $\theta = 90^\circ$, we have $f = f_0 \sqrt{1 - \frac{v^2}{c^2}}$

Example: light source red $\lambda_0 = 650\text{nm}$, how large the speed is for a driver to see it is green $\lambda = 500\text{nm}$?

$$\frac{f}{f_0} = \sqrt{\frac{1+\beta}{1-\beta}} = \frac{\lambda_0}{\lambda} = \frac{650}{500} \Rightarrow \frac{v}{c} = \frac{1 - (\frac{\lambda}{\lambda_0})^2}{1 + (\frac{\lambda}{\lambda_0})^2} = 0.26c.$$

SSC AMPAD
4-vector

we have

$$\begin{pmatrix} cdt' \\ dx' \\ dy' \\ dz' \end{pmatrix} = \begin{pmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} cdt \\ dx \\ dy \\ dz \end{pmatrix} \leftarrow \text{4-vector}$$

$(cdt)^2 - [(dx)^2 + (dy)^2 + (dz)^2]$
 $= (cdt')^2 - [(dx')^2 + (dy')^2 + (dz')^2]$

For a particle with the rest mass m_0 , ~~and~~ and velocity \vec{u} , we

define its proper time $d\tau = \sqrt{1 - \frac{u^2}{c^2}} dt$
time interval measured in the particle rest frame. time interval measured in the lab frame

$d\tau$ and m_0 is independent of frame transformation.

Define $\left(\frac{m_0 c dt}{d\tau}, \frac{m_0 dx}{d\tau}, \frac{m_0 dy}{d\tau}, \frac{m_0 dz}{d\tau} \right)$ multiply $m_0/d\tau$ to the space-time interval

$$= [\gamma m_0 c, \gamma m_0 u_x, \gamma m_0 u_y, \gamma m_0 u_z] = \left[\frac{E}{c}, p_x, p_y, p_z \right]$$

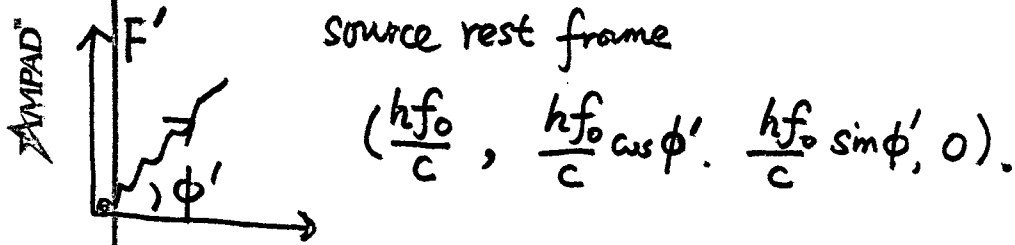
Energy and momentum follow the same transformation of (cdt, dx, dy, dz)
relativistic

$\Rightarrow \left(\frac{E}{c}\right)^2 - (\vec{P})^2$ is invariant, which equals to $(m_0 c)^2$.

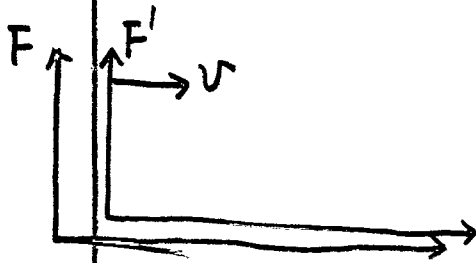
we can also apply it to light. For light $E = hf$, and $p = \frac{hf}{c}$.

Rederive Doppler effect.

and $\left(\frac{E}{c}\right)^2 - p^2 \equiv 0$



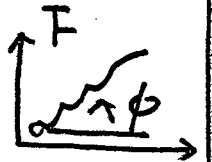
in the lab frame (receiver's)



in F frame, the energy-momentum vector

$\left(\frac{hf}{c}, \frac{hf}{c} \cos \phi, \frac{hf_0}{c} \sin \phi, 0\right)$

$$\Rightarrow \begin{matrix} \text{lab} \\ F' \end{matrix} \rightarrow \begin{pmatrix} \frac{hf_0}{c} \\ hf_0 \cos \phi' \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} \frac{hf}{c} \\ \frac{hf}{c} \cos \phi \end{pmatrix} \leftarrow F$$



$\Rightarrow \frac{hf_0}{c} = \gamma \frac{hf}{c} - \beta\gamma \frac{hc}{c} \cos \phi$

$\Rightarrow f_0 = \gamma(1 - \beta \cos \phi) f$

$$\Rightarrow f = \frac{\sqrt{1 - \frac{v^2}{c^2}}}{1 - \frac{v}{c} \cos \phi} f_0$$