

1) Consider a system of spin- $\frac{1}{2}$  neutral fermions at zero temperature interacting through the Yukawa potential

$$V(r) = \left( +\frac{g}{4\pi r} \right) e^{-r/a}$$

a) Calculate the Landau interaction function  $f_{pp',\infty'}^{(0)}$  at the Hartree-Fock level, and get Landau parameters  $F_0^s F_0^a, F_1^s, F_1^a$  up to  $a^4$ .

b) determine the factors of the specific heat, compressibility and magnetic susceptibility in the Fermi liquid theory compared to the non-interacting Fermi gas with the same bare mass and density.

c) Apply the above result to the screened Coulomb potential in the interacting electron gas by using Thomas-Fermi model.

Consider the low density limit ( $k_f \ll k_{FT}$ ). Compare the magnetic susceptibility and compressibility to the "free-gas" value defined in 2)? Should we take this result seriously?

(from Leggett's HW 3, Phy 490 Fall 2001).

a) Derive the effective mass renormalization.

a) Consider a fermi-liquid system. By summing the Boltzmann equation over momentum  $p$ , and using the fact that the collision integral must conserve the local particle density, derive the continuity equation

$$\frac{\partial P}{\partial t} + \nabla_i j_i = 0, \text{ where } \vec{j} \text{ is particle current}$$

defined as  $j_i = \frac{1}{V_{P^0}} \sum_{p_i} \frac{\partial \epsilon_{p^0}(r,t)}{\partial p_i} n_{p^0}(r,t)$

b) Linearize the expression of  $j_i$ , and keep in mind that  $\epsilon_{p^0}(r,t)$  also depends on the distribution charge  $\delta n_{p^0}(r,t) = n_{p^0}(r,t) - n_{p^0}^0$ .

Show  $j_i(r,t) = \frac{1 + F_i/3}{m^*} \frac{1}{V_{P^0}} \sum_{p_i} p_i \delta n_{p^0}(r,t)$

c) In Galilean invariant system, by relativistic continuity principle. Argue

that  $j_i(r,t) = \frac{1}{V_{P^0}} \sum_{p_i} \frac{p_i}{m} \delta n_{p^0}(r,t)$ , show that  $\frac{m^*}{m} = 1 + F_i/3$ .

d) Can you apply the above result to spin current? Please show

$$\vec{j}_i = \frac{1 + F_i/3}{m^*} \frac{1}{V} \sum_{p_i} p_i \vec{\sigma} \delta n_{p^0}(r,t), \text{ and we can define}$$

the spin effective mass  $\frac{1}{m_s^*} = \frac{1 + F_i/3}{m^*}$ , i.e.  $m_s^* = \frac{1 + F_i/3}{1 + F_i/3} m$ .

Why  $m_s^*$  is not necessarily equal to  $m$ ?

### 3. Sound excitations

- ① First sound or hydrodynamic sound. Prove sound velocity equals

$v^2 = \frac{1}{\lambda\rho}$ , where  $\lambda$  is the compressibility,  $\rho$  is the mass density. Show  $v^2 = \frac{n}{m} \left( \frac{\partial K}{\partial n} \right)$ . In the ideal fermi gas  $\frac{v_{\text{ideal}}}{v_{\text{show}}} = \sqrt{\frac{4}{3}}$ .

$$\text{and in Fermi liquid } \left( \frac{v_{\text{FL}}}{v_{\text{ideal}}} \right)^2 = \frac{1 + F_0^3}{1 + \frac{1}{3} F_1^3}$$

- ② in the class, we have derived the equation for the zero sound

we write  $\delta n(p, r, t) = \sum_{\vec{q}} \delta n(\vec{p}) e^{i(\vec{qr}-\omega t)}$  and

$$\delta n(\vec{p}) = -\frac{\partial n_{\vec{p}}}{\partial \epsilon_{\vec{p}}} v_{\vec{p}}.$$

$$\text{and } (\omega - \omega_0) v_{\vec{p}} - \omega_0 \sum_{\vec{p}'} f_{\vec{p}\vec{p}'} \left( -\frac{\partial n_{\vec{p}}}{\partial \epsilon_{\vec{p}}} \right) v_{\vec{p}'} = 0.$$

By expanding  $v_{\vec{p}} = \sum_{\ell=0}^{\infty} Y_{\ell 0}(\hat{\vec{p}}) u_{\ell 0}$  ( $Y_{\ell 0}(\hat{\vec{p}})$  is the spherical harmonic function)  
(set  $\hat{\vec{q}}$  along  $\hat{\vec{z}}$ -axis).

Show that

$$\frac{u_{\ell 0}}{2\ell+1} + \sum_{\ell'} \sqrt{2} u_{\ell'}(s) F_{\ell'}^s \frac{u_{\ell' 0}}{2\ell'+1} = 0, \text{ where}$$

$$\sqrt{2} u_{\ell'}(s) = \int_0^{\pi} d\cos\theta \int_0^{2\pi} d\phi \quad Y_{\ell 0}(\cos\theta) \frac{\cos\theta}{\cos\theta - s} \quad Y_{\ell' 0}(\cos\theta)$$

if we truncate at  $\ell=0$ , we get  $F_0^s \int \frac{d\Omega}{4\pi} \frac{\cos\theta}{\cos\theta - s} = 1$ .

and show  $S \rightarrow 1 + 2e^{-\frac{2}{3}F_0}$  ( $F_0 \rightarrow 0$ ), and

$$S \rightarrow \sqrt{\frac{F_0}{3}} \quad (F_0 \rightarrow +\infty)$$

③ Prove that

$$\sqrt{2_{00}} = 1 - \frac{s}{2} \ln \left( \left| \frac{s+1}{s-1} \right| \right)$$

$$\sqrt{2_{01}} = \sqrt{2_{10}} = s\sqrt{2_{00}}, \quad \sqrt{2_{11}} = s^2\sqrt{2_{00}} + 1/3$$

If only  $F_0^s \neq 0$ , the eigenmode of the zero sound has the configuration of  $U_{l0} = (2l+1) \frac{\sqrt{2_{l0}}}{\sqrt{2_{00}}} \sin(l\pi x) \quad (l \geq 1)$ .

From the symmetry point of view, explain why  $\{U_{l0}\} \quad (l=0, 1, \dots)$  can couple together.

If we keep both  $F_0^s$  and  $F_1^s$  nonzero show that the zero sound frequency satisfies

$$\frac{1}{2}s \ln \left( \left| \frac{s+1}{s-1} \right| \right) - 1 = \frac{1 + \frac{1}{3}F_1}{F_0 + s^2 F_1 + \frac{1}{3}F_0 F_1}$$