

Solution to HW 3

①

1.a) $\psi_{\sigma}(r) = \sum_{i\sigma} \phi_{i\sigma}(r) a_{i\sigma}$

$$H_1 = \int dr \psi_{\sigma}^{\dagger}(r) \left(-\frac{\hbar^2}{2m} \nabla^2 + U(r) \right) \psi_{\sigma}(r) = \sum_{i\sigma j\sigma'} \int \phi_{i\sigma}^{\dagger}(r) \overbrace{\phi_{j\sigma'}(r)}^{(-\frac{\hbar^2}{2m} \nabla^2 + U(r))} dr a_{i\sigma}^{\dagger} a_{j\sigma}$$

$$= \sum_{i,j,\sigma} \int dr \phi_{i\sigma}^{\dagger}(r) \left[-\frac{\hbar^2}{2m} \nabla^2 + U(r) \right] \phi_{j\sigma}(r) a_{i\sigma}^{\dagger} a_{j\sigma}$$

$$H_2 = \frac{1}{2} \int dr_1 dr_2 \psi_{\sigma}^{\dagger}(r_1) \psi_{\sigma'}^{\dagger}(r_2) V(r_1 - r_2) \psi_{\sigma'}(r_2) \psi_{\sigma}(r_1)$$

$$= \frac{1}{2} \int dr_1 dr_2 \sum_{\substack{ij\ell k \\ \sigma\sigma'}} \underbrace{\phi_{i\sigma}^{\dagger}(r_1) a_{i\sigma}^{\dagger} \phi_{j\sigma'}^{\dagger}(r_2) a_{j\sigma'}^{\dagger}}_{V(r_1 - r_2)} \phi_{\ell\sigma'}(r_2) a_{\ell\sigma'} \phi_{k\sigma}(r_1) a_{k\sigma}$$

for Coulomb interaction:

$$H_2 = \frac{e^2}{2} \sum_{\substack{ij\ell k \\ \sigma\sigma'}} \int dr_1 dr_2 \frac{\phi_{i\sigma}^{\dagger}(r_1) \phi_{j\sigma'}^{\dagger}(r_2) \phi_{\ell\sigma'}(r_2) \phi_{k\sigma}(r_1)}{|r - r'|} a_{i\sigma}^{\dagger} a_{j\sigma'}^{\dagger} a_{\ell\sigma'} a_{k\sigma}$$

1.b) in the plane-wave basis.

$$\phi_{k\sigma}(r) = \frac{1}{\sqrt{V}} e^{i\vec{k}\cdot\vec{r}}, \text{ and if } U(r) \text{ is a constant}$$

$$H_1 = \sum_{\mathbf{k}} \frac{\hbar^2 k^2}{2m} a_{k\sigma}^{\dagger} a_{k\sigma}$$

$$H_2 = \frac{e^2}{2} \sum_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4} \int dr_1 dr_2 \frac{e^{-i\mathbf{k}_1 \cdot r_1} e^{-i\mathbf{k}_2 \cdot r_2} e^{i\mathbf{k}_3 \cdot r_2} e^{i\mathbf{k}_4 \cdot r_1}}{|r_1 - r_2|} a_{\mathbf{k}_1, \sigma}^{\dagger} a_{\mathbf{k}_2, \sigma'}^{\dagger} a_{\mathbf{k}_3, \sigma'} a_{\mathbf{k}_4, \sigma}$$

(2)

introduce center of mass coordinate

$$R = \frac{r_1 + r_2}{2}, \quad r = r_1 - r_2$$

$$\Rightarrow r_1 = R + \frac{r}{2}$$

$$e^{-i(k_1 - k_2 + k_3 - k_4) \cdot \frac{r_1 + r_2}{2}} = e^{-i(k_1 - k_2 + k_3 - k_4) \cdot R}$$

$$\times V(r)$$

$$\int dr_1 dr_2 \dots = \int dR dr e^{-i(k_1 + k_2 - k_3 - k_4) \cdot R}$$

$$= \delta_{k_1 + k_2 = k_3 + k_4} \cdot V\left[\frac{(k_1 - k_2 + k_3 - k_4)}{2}\right]$$

$$\text{set } k_3 = k_2 - q, \quad k_4 = k_1 + q \Rightarrow$$

$$H_2 = \frac{1}{2} \sum_{k_1 k_2 q} \widehat{V(q)} a_{k_1 \sigma}^+ a_{k_2 \sigma'}^+ a_{k_2 - q \sigma'} a_{k_1 + q \sigma}, \quad \text{where } V(q) = \frac{4\pi e^2}{q^2}$$

2a): $|\psi\rangle = a_{i_1\sigma_1}^+ a_{i_2\sigma_2}^+ \dots a_{i_n\sigma_n}^+ |0\rangle$

$$\langle\psi|H_0|\psi\rangle = \langle 0| a_{i_n\sigma_n} \dots a_{i_1\sigma_1} \left[\sum_{ij\sigma} \langle i|\hat{h}|j\rangle a_{i\sigma}^+ a_{j\sigma} \right] a_{i_1\sigma_1}^+ \dots a_{i_n\sigma_n}^+ |0\rangle$$

we have to set $i=j$, which can be any one of i_1, \dots, i_n

$$\Rightarrow \langle\psi|H_0|\psi\rangle = \sum_{i\sigma} \langle i|\hat{h}|i\rangle n_{i\sigma} = \sum_{i\sigma} n_{i\sigma} \int dr \left[\phi_i^*(r) \left(-\frac{\hbar^2}{2m} \nabla^2 + U(r) \right) \phi_i(r) \right]$$

$$\langle\psi|H_2|\psi\rangle = \langle 0| a_{i_n\sigma_n} \dots a_{i_1\sigma_1} \left(\frac{1}{2} \sum_{\substack{ij\ell k \\ \sigma\sigma'}} \langle ij|V|\ell k\rangle a_{i\sigma}^+ a_{j\sigma'}^+ a_{\ell\sigma} a_{k\sigma'} \right) a_{i_1\sigma_1}^+ \dots a_{i_n\sigma_n}^+ |0\rangle$$

we have to pair the indices $ij\ell k$

① direct channel $a_{i\sigma}^+ a_{j\sigma'}^+ a_{\ell\sigma} a_{k\sigma'}$

$i=k$, and $j=l$, each of them can take i_1, \dots, i_n

$$\Rightarrow \frac{e^2}{2} \sum_{ij\sigma\sigma'} n_{i\sigma} n_{j\sigma'} \int dr dr' \frac{|\phi_i(r)|^2 |\phi_j(r')|^2}{|r-r'|} \leftarrow \langle ij|V|ij\rangle$$

② exchange channel $a_{i\sigma}^+ a_{j\sigma'}^+ a_{\ell\sigma} a_{k\sigma'}$

$i=l$, $j=k$ and $\sigma=\sigma'$, due to the cross, an extra minus sign

$$\Rightarrow \frac{e^2}{2} (-) \sum_{ij\sigma\sigma'} n_{i\sigma} n_{j\sigma'} \delta_{\sigma\sigma'} \langle ij|V|ji\rangle = -\frac{e^2}{2} \sum_{ij\sigma\sigma'} n_{i\sigma} n_{j\sigma'} \underbrace{\delta_{\sigma\sigma'}}_{\text{cross}} \frac{\int dr dr' (\phi_i^*(r) \phi_j^*(r') \times \phi_j(r) \phi_i(r'))}{|r-r'|}$$

add together \Rightarrow

$$\langle \Psi | H | \Psi \rangle = \sum_{i\sigma} n_{i\sigma} \int d\mathbf{r} \left\{ \phi_i^*(\mathbf{r}) \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{r}) \right) \phi_i(\mathbf{r}) \right\}$$

$$+ \frac{e^2}{2} \sum_{i,j\sigma,\sigma'} n_{i\sigma} n_{j\sigma'} \int d\mathbf{r} d\mathbf{r}' \left\{ \frac{|\phi_i(\mathbf{r})|^2 |\phi_j(\mathbf{r}')|^2}{|\mathbf{r}-\mathbf{r}'|} - \delta_{\sigma\sigma'} \frac{\phi_i^*(\mathbf{r}) \phi_j^*(\mathbf{r}') \phi_j(\mathbf{r}) \phi_i(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} \right\}$$

b) introduce Lagrange multiplier λ_i for each $\int d\mathbf{r} |\phi_i(\mathbf{r})|^2 = 1$.

$$E = \langle \Psi | H | \Psi \rangle - \sum_{i\sigma} \lambda_{i\sigma} \int |\phi_{i\sigma}(\mathbf{r})|^2 d\mathbf{r}, \quad \text{where } \phi_{i\sigma}(\mathbf{r}) = \phi_i(\mathbf{r}) \chi_\sigma$$

χ_σ is the spin WF

$$\frac{\delta E}{\delta \phi_{i\sigma}^*} = \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{r}) \right) \phi_i(\mathbf{r}) \chi_\sigma$$

$$- \frac{e^2}{2} \sum_{j\sigma'} \delta_{\sigma\sigma'} n_{j\sigma'} \int d\mathbf{r}' \frac{\phi_j^*(\mathbf{r}') \phi_j(\mathbf{r})}{|\mathbf{r}-\mathbf{r}'|} \phi_i(\mathbf{r}') \chi_\sigma$$

$$+ \frac{e^2}{2} \sum_{j\sigma'} n_{j\sigma'} \int d\mathbf{r}' \frac{|\phi_j(\mathbf{r}')|^2}{|\mathbf{r}-\mathbf{r}'|} \phi_i(\mathbf{r}) \chi_\sigma - \lambda_{i\sigma} \phi_i(\mathbf{r}) \chi_\sigma = 0$$

(set $n_{i\sigma} = 1$ (occupied)).

i.e

$$\left\{ -\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{r}) + \sum_{j\sigma'} n_{j\sigma'} \int d\mathbf{r}' \frac{|\phi_j(\mathbf{r}')|^2}{|\mathbf{r}-\mathbf{r}'|} \right\} \phi_i(\mathbf{r}) \chi_\sigma$$

$$- \sum_j \delta_{\sigma\sigma'} n_{j\sigma'} \int d\mathbf{r}' \frac{\phi_j^*(\mathbf{r}') \phi_j(\mathbf{r})}{|\mathbf{r}-\mathbf{r}'|} \phi_i(\mathbf{r}') \chi_\sigma = \lambda_{i\sigma} \phi_i(\mathbf{r})$$

(5)

3c) $\phi_{i\sigma}(r) = e^{ik_i \cdot r} \chi_\sigma$, plug in to the HF equation

and set $U(r) = 0$.

$$\begin{aligned} & \left\{ -\frac{\hbar^2}{2m} \nabla^2 + \sum_{j\sigma'} \int dr' \frac{n_{j\sigma'}}{|r-r'|} \right\} e^{ik_i \cdot r} \chi_\sigma \\ & - \sum_j \delta_{\sigma\sigma'} n_{j\sigma'} \int dr' \frac{e^{ik_j \cdot (r-r')}}{|r-r'|} \cdot e^{ik_i \cdot r'} \chi_\sigma \\ & = \left\{ -\frac{\hbar^2}{2m} \nabla^2 + \sum_{j\sigma'} \int dr' \frac{n_{j\sigma'}}{|r-r'|} \right\} e^{ik_i \cdot r} \chi_\sigma \\ & - \left\{ \sum_j \delta_{\sigma\sigma'} n_{j\sigma'} \int dr' \frac{e^{i(k_j - k_i) \cdot (r-r')}}{|r-r'|} \right\} e^{ik_i \cdot r} \chi_\sigma = \lambda_{i\sigma} e^{ik_i \cdot r} \chi_\sigma \end{aligned}$$

$$\text{Thus } \lambda_{i\sigma} = \frac{\hbar^2 k_i^2}{2m} + \sum_{j\sigma'} \int dr' \frac{n_{j\sigma'}}{|r-r'|}$$

$$- \sum_j \delta_{\sigma\sigma'} n_{j\sigma'} \int dr' \frac{e^{i(k_j - k_i) \cdot (r-r')}}{|r-r'|}$$

thus $\lambda_{i\sigma}$ has the meaning of particle energy

$$\begin{aligned} & = \text{Kinetic} + \text{Hartree} + \text{Fock} \\ & \quad \quad \quad (\text{direct}) \quad \quad (\text{exchange}) \end{aligned}$$

$$\textcircled{3} \quad \rho_{\sigma}(r) = \psi_{\sigma}^{\dagger}(r) \psi_{\sigma}(r) = \sum_{ij} \phi_i^*(r) \phi_j(r) a_{i\sigma}^{\dagger} a_{j\sigma}$$

$$\langle \psi | \rho_{\sigma}(r) \rho_{\sigma'}(r') | \psi \rangle = \langle \psi | \sum_{ij\sigma} \phi_i^*(r) \phi_j(r) \phi_{j'}^*(r') \phi_{i'}(r') a_{i\sigma}^{\dagger} a_{j\sigma} a_{j'\sigma'}^{\dagger} a_{i'\sigma'} | \psi \rangle$$

direct channel $i=j$ and $j'=i'$. \Rightarrow $\underbrace{a_{i\sigma}^{\dagger} a_{j\sigma}}_{+} \underbrace{a_{j'\sigma'}^{\dagger} a_{i'\sigma'}}_{+}$

$$\sum_{ij\sigma\sigma'} |\phi_i(r)|^2 |\phi_j(r')|^2 n_{i\sigma} n_{j\sigma'} \quad (\text{for } i \neq j \text{ or } \sigma \neq \sigma')$$

which is just $\langle \rho_{\sigma}(r) \rangle \langle \rho_{\sigma'}(r') \rangle$

exchange channel $i=i'$ & $j=j'$ & $\sigma=\sigma'$ $\underbrace{a_{i\sigma}^{\dagger} a_{j\sigma} a_{j'\sigma'}^{\dagger} a_{i'\sigma'}}_{+}$

$$\rightarrow - \sum_{ij\sigma'} \delta_{\sigma\sigma'} \phi_i^*(r) \phi_j^*(r') \phi_j(r) \phi_i(r') n_{i\sigma} (n_{j\sigma} - 1)$$

$$= - \sum_{ij\sigma'} \delta_{\sigma\sigma'} \phi_i^*(r) \phi_j^*(r') \phi_j(r) \phi_i(r') n_{i\sigma} n_{j\sigma} + \sum_{ij\sigma'} \delta_{\sigma\sigma'} \phi_i^*(r) \phi_i(r') n_{i\sigma} \cdot \sum_j \phi_j(r) \phi_j^*(r')$$

$$= - \sum_{ij\sigma'} \delta_{\sigma\sigma'} \phi_i^*(r) \phi_j^*(r') \phi_j(r) \phi_i(r') n_{i\sigma} n_{j\sigma} + \sum_{i\sigma'} \delta_{\sigma\sigma'} |\phi_i(r)|^2 n_{i\sigma} \delta(r-r')$$

\hookrightarrow this an artifact due to normalization

$$\Rightarrow G_{\sigma\sigma'}(r, r') = \langle \psi | \rho_{\sigma}(r) \rho_{\sigma'}(r') | \psi \rangle$$

$$= \langle \psi | \rho_{\sigma}(r) | \psi \rangle \langle \psi | \rho_{\sigma'}(r') | \psi \rangle$$

at $r \neq r' \Rightarrow$

$$G_{\sigma\sigma'}(r, r') = 0 \quad \text{if } \sigma \neq \sigma'$$

$$= - \sum_{ij} \varphi_i^*(r) \varphi_j^*(r') \varphi_j(r) \varphi_i(r') n_{i\sigma} n_{j\sigma} \quad \text{if } \sigma = \sigma'$$

For plane wave \Rightarrow

$$G_{\sigma\sigma'}(r, r') = - \frac{1}{V} \sum_{\vec{k}_i \vec{k}_j} e^{i\vec{k}_i \cdot (r-r')} e^{i\vec{k}_j \cdot (r'-r)} n_{\vec{k}_i \sigma} n_{\vec{k}_j \sigma}$$

$$= - \left[\sum_{\vec{k}_i} e^{i\vec{k}_i \cdot (r-r')} n_{\vec{k}_i \sigma} \right]^2$$

$$= - \left| \frac{1}{(2\pi)^3} \int d^3\vec{k} e^{i\vec{k} \cdot (\vec{r}-\vec{r}')} \Theta(k < k_F) \right|^2 \quad (\text{for } \sigma = \sigma')$$

The detail of how to do the integral is in the lecture notes

and I won't repeat.

$$G_{\sigma\sigma'}(r, r') / \langle \rho_\sigma \rangle \langle \rho_{\sigma'} \rangle$$

$$= -9 \left(\frac{\chi \cos \chi - \sin \chi}{\chi^3} \right)^2$$

