

Lecture 3 Coriolis force

$$\vec{F}_{\text{cor}} = 2m \vec{V} \times \vec{\Omega}$$

① velocity-dependent force, which cannot be written as gradient of a scalar potential, but as a curl of vector potential

② dissipationless $\vec{F} \perp \vec{V}$ c.f. friction force $\parallel \vec{V}$.

Analogy to Lorentz force $\vec{F}_L = q \frac{\vec{V}}{c} \times \vec{B}$. $\vec{B} \leftrightarrow \vec{\Omega}$

Example: vortices in type II superconductor $\leftrightarrow \vec{B} = \nabla \times \vec{A}$
 vortices in superfluid ${}^4\text{He}$ $\leftrightarrow \vec{\Omega} = \nabla \times \vec{A}_{\Omega}$

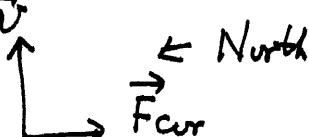
§ Direction of \vec{F}_{cor} on the earth.

For object moving on the surface of the earth

\vec{V} is in the tangent plane, only the normal component of $\vec{\Omega}_{\perp}$ can result in an \vec{F}_{cor} in-plane.

In the north-hemisphere $\vec{\Omega}_{\perp}$ point out of the ground

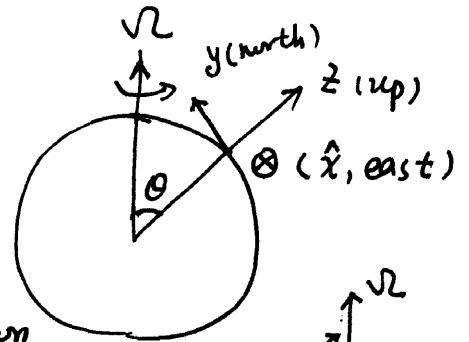
$$\vec{F}_{\text{cor, inplane}} = \cancel{\vec{V} \times \vec{\Omega}_{\perp}} \rightarrow \begin{matrix} \text{right hand side of } \vec{V} \\ \vec{V} \end{matrix}$$



In the south-hemisphere, $\vec{F}_{\text{cor}} \rightarrow$ left hand side of \vec{V} .

§ 2. Easterly deflection of a falling object

$$m \ddot{\vec{r}} = \underbrace{m \vec{g}_0}_{m\vec{g}} + \vec{F}_{cf} + \vec{F}_{cor}$$



define up. direction (\hat{z}) as the opposite direction of \vec{g} , north (\hat{y}) perpendicular to \hat{z} and \vec{v} lies in the yz plane. \hat{x} is the east.

$$\vec{v} = (0, \sqrt{2} \sin \theta, \sqrt{2} \cos \theta), \quad \theta \text{ the angle between } \hat{v} \text{ and } \hat{z}\text{-axis.}$$

$$\dot{\vec{r}} \times \vec{v} = (\dot{y} \sqrt{2} \cos \theta - \dot{z} \sqrt{2} \sin \theta, -\dot{x} \sqrt{2} \cos \theta, \dot{x} \sqrt{2} \sin \theta)$$

$$\Rightarrow \ddot{\vec{r}} = -g \hat{z} + 2 \dot{\vec{r}} \times \vec{v}$$

$$\ddot{x} = 2\sqrt{2}(\dot{y} \cos \theta - \dot{z} \sin \theta), \quad \ddot{y} = -2\sqrt{2} \dot{x} \cos \theta, \quad \ddot{z} = -g + 2\sqrt{2} \dot{x} \sin \theta$$

- zero order $\Rightarrow x = y = 0 \text{ & } z = h - \frac{1}{2} g t^2$

- first order $\ddot{x} = +2\sqrt{2} g t \sin \theta \Rightarrow x = \frac{1}{3} \sqrt{2} g t^3 \sin \theta$

earth spins from west to east; the linear velocity at higher place is larger than that at lower altitude.

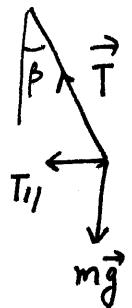
Estimat: for $\theta = 90^\circ$, $h = 100m$

$$\Delta x = \frac{1}{3} \sqrt{2} g \left(\frac{2h}{g} \right)^{\frac{3}{2}} \approx 2 \text{ cm.}$$

§4: Foucault pendulum: (See the spin of the earth)

The rotation of the oscillation plane.

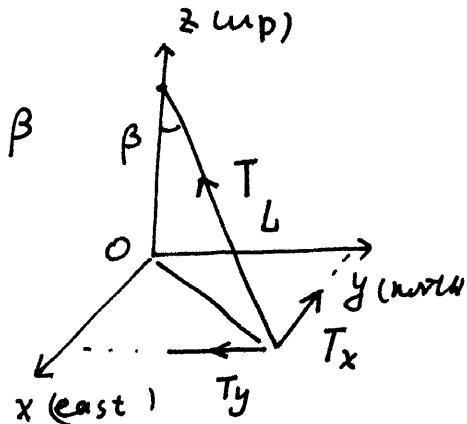
$$m \ddot{\vec{r}} = \vec{T} + \underbrace{m \vec{g}_0 + m(\vec{\omega} \times \vec{r}) \times \vec{\omega}}_{m \vec{g}} + 2m \dot{\vec{r}} \times \vec{\omega}$$



$$T_{||} \approx T \cdot \beta \approx mg \beta$$

$$T_{||x} \approx -mg \frac{x}{L}$$

$$T_y \approx -mg \frac{y}{L}$$



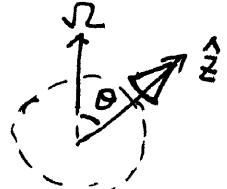
$$\Rightarrow \ddot{x} = -\frac{g x}{L} + 2\dot{y}\sqrt{\omega^2 \cos^2 \theta} \\ \ddot{y} = -\frac{g y}{L} - 2\dot{x}\sqrt{\omega^2 \cos^2 \theta}$$

$\sqrt{\omega_z^2} = \sqrt{2}\omega \sin \theta$ is the component of $\vec{\omega}$ perpendicular the XY plane

define $\omega_0^2 = \frac{g}{L}$ \Rightarrow

θ the angle between $\hat{\omega}$ & \hat{z}

$$\ddot{x} - 2\sqrt{\omega_z} \dot{y} + \omega_0^2 x = 0 \\ \ddot{y} - 2\sqrt{\omega_z} \dot{x} + \omega_0^2 y = 0$$



$$\text{define } \eta = x + iy \Rightarrow \ddot{\eta} + 2i\sqrt{\omega_z} \dot{\eta} + \omega_0^2 \eta = 0$$

eigen-frequency

$$e^{-i\alpha t} \Rightarrow -\alpha^2 + 2\sqrt{\omega_z} \alpha - \omega_0^2 = 0$$

$$\alpha = \sqrt{\omega_z} \pm \sqrt{\sqrt{\omega_z^2} + \omega_0^2} \approx \sqrt{\omega_z} \pm \omega_0$$

$$\Rightarrow \eta(t) = e^{-i\sqrt{\omega_z}t} [C_1 e^{i\omega_0 t} + C_2 e^{-i\omega_0 t}]$$

if we choose linear polarization at $t=0$, i.e. $C_1 = C_2$

$$\Rightarrow x(t) + iy(t) = \boxed{e^{-i\sqrt{\omega_z}t}} \cos \omega_0 t$$

rotation of oscillation plane at $\sqrt{\omega_z} = \sqrt{2}\omega \sin \theta$.