

Lect 5: Momentum and angular momentum

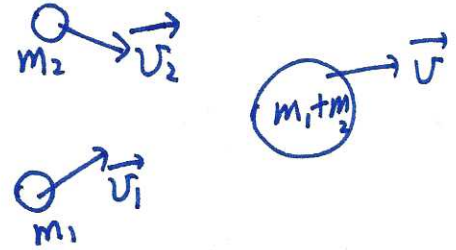
§: Conservation of momentum — space translation symmetry

- Consider a N-particle system, the total momentum $\vec{P}_t = \sum_{\alpha=1}^N \vec{P}_{\alpha}$ is conserved if the net external force \vec{F}^{ext} is zero. More generally,

$$\dot{\vec{P}}_t = \vec{F}^{ext}$$

Internal force only changes the distribution of \vec{P}_t among different particles, but do not change \vec{P}_t .

Example: inelastic collision of 2-bodies



$$m_1 \vec{v}_1 + m_2 \vec{v}_2 = (m_1 + m_2) \vec{v}$$

$$\Rightarrow \vec{v} = \frac{m_1 \vec{v}_1 + m_2 \vec{v}_2}{m_1 + m_2}$$

Continuum

$$\vec{R} = \frac{1}{M} \int \vec{r} dm = \frac{1}{M} \int \rho \vec{r} dV$$

- Center of mass:

$$\vec{R} = \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \vec{r}_{\alpha} = \frac{m_1 \vec{r}_1 + \dots + m_N \vec{r}_N}{M} \Rightarrow \vec{P}_t = M \dot{\vec{R}}$$

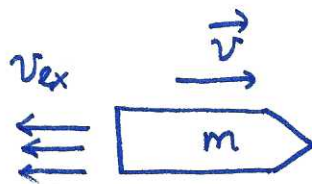
$$\dot{\vec{P}}_t = \vec{F}^{ext} \Rightarrow \boxed{M \ddot{\vec{R}} = \vec{F}^{ext}}$$

\Rightarrow The center of mass motion of a N-particle system is the same as a point particle. This is the reason we can idealize the center of mass motion (translation) of an extended object as mass point motion.

For example, the earth motion around the sun.

Example: Recoil - rockets

the spent fuel is ejected at speed v_{ex} relative to the rocket.



At time t : $P(t) = mv$

At time $t + \Delta t$: the rocket mass $m + dm$, ($dm < 0$)

the spent fuel: $-dm$

$$\Rightarrow P(t + \Delta t) = (m + dm)(v + dv) - dm(v - v_{ex})$$

$$\approx mv + mdv + v_{ex}dm$$

$$P(t + \Delta t) = P(t) \Rightarrow mdv + v_{ex}dm = 0$$

$$dv = -v_{ex} \frac{dm}{m} \Rightarrow v - v_0 = -v_{ex} \int_{m_0}^m \frac{dm'}{m'} = v_{ex} \ln \frac{m_0}{m}$$

Consider a rocket filled 90% fuel, then $\frac{m_0}{m} = 10 \Rightarrow v - v_0 = v_{ex} \ln 10 = 2.3 v_{ex}$.

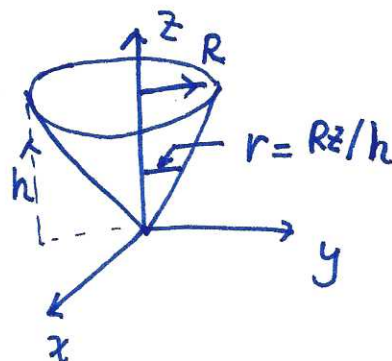
2: The CM of a solid cone

$$\vec{R} = \frac{1}{M} \int \rho \vec{r} dv = \frac{\rho}{M} \int \vec{r} dv$$

$$R_x = \frac{\rho}{M} \iiint dx dy dz \cdot x = 0$$

$$R_y = \frac{\rho}{M} \iiint dx dy dz \cdot y = 0$$

$$R_z = \frac{\rho}{M} \iiint dx dy dz \cdot z = \frac{\rho \int_0^h z dz \int_0^{Rz/h} r dr \int_0^{2\pi} d\theta}{\rho \int_0^h dz \int_0^{Rz/h} r dr \int_0^{2\pi} d\theta} = \frac{\int_0^h dz \cdot \frac{z}{2} \left(\frac{Rz}{h}\right)^2}{\int_0^h dz \cdot \frac{1}{2} \left(\frac{Rz}{h}\right)^2} = \frac{\int_0^h z^3 dz}{\int_0^h z^2 dz} = \frac{\frac{1}{4} h^4}{\frac{1}{3} h^3} = \frac{3}{4} h = R_z$$



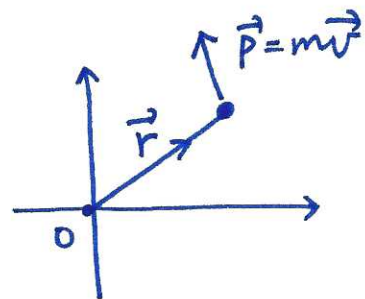
§ Angular momentum for a single particle

$$\vec{L} = \vec{r} \times \vec{p}$$

$$\dot{\vec{L}} = \dot{\vec{r}} \times \vec{p} + \vec{r} \times \dot{\vec{p}} = \vec{v} \times m\vec{v} + \vec{r} \times \vec{F} = \vec{r} \times \vec{F}$$

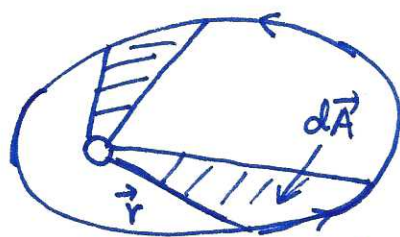
For the central force field, $\vec{F} \parallel \vec{r}$, $\Rightarrow \dot{\vec{L}} = 0$,

\Rightarrow angular momentum is conserved.



Example: Kepler's problem

1st law: The planet orbit is planar and closed forming an ellipse.



$$d\vec{r} = \vec{v} dt$$

2nd law: a line drawn from the planet to the sun

sweeps out equal areas in equal times.

① \vec{L} is conserved: \rightarrow the orbit is planar. The orbital plane is perpendicular to \vec{L} .

② The magnitude of \vec{L} : $\vec{L} = \vec{r} \times m\vec{v}$
 $d\vec{A} = \vec{r} \times d\vec{r} \Rightarrow \frac{d\vec{A}}{dt} = \vec{r} \times \vec{v} = \frac{\vec{L}}{m}$

Kepler's 2nd law is equivalent to angular momentum conservation

The closed orbit and ellipse requires stronger condition $\rightarrow 1/r$ potential

or $-1/r^2 \hat{e}_r$ force.

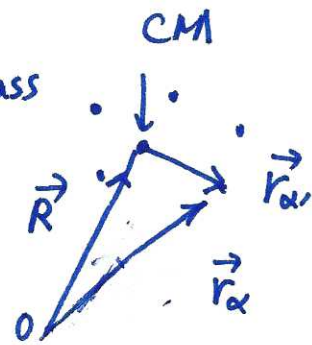
Angular momentum for several particles

$$\alpha = 1, 2, \dots, N$$

$$\vec{l}_\alpha = \vec{r}_\alpha \times \vec{p}_\alpha$$

$$\vec{L}_{tot} = \sum_{\alpha=1}^N \vec{l}_\alpha = \sum_{\alpha=1}^N \vec{r}_\alpha \times \vec{p}_\alpha$$

\vec{R} : center of mass coordinate
 \vec{r}'_α : center of mass vector



$$\vec{L}_{tot} = \sum_{\alpha=1}^N (\vec{R} + \vec{r}'_\alpha) \times \vec{p}_\alpha = \vec{R} \times \sum_{\alpha=1}^N \vec{p}_\alpha + \sum_{\alpha=1}^N \vec{r}'_\alpha \times m_\alpha \dot{\vec{r}}_\alpha$$

$$= \vec{R} \times \vec{p}_{tot} + \sum_{\alpha=1}^N \vec{r}'_\alpha \times m_\alpha (\dot{\vec{r}}_\alpha + \dot{\vec{R}})$$

$$= \vec{R} \times \vec{p}_{tot} + \sum_{\alpha=1}^N \vec{r}'_\alpha \times m_\alpha \dot{\vec{r}}_\alpha + \left(\sum_{\alpha=1}^N m_\alpha \vec{r}'_\alpha \right) \times \dot{\vec{R}}$$

"0" please prove it

orbit \vec{L}_{orbit}

internal angular momentum with respect to CM: \vec{L}_{spin}

$$\vec{L}_{tot} = \vec{L}_{orbit} + \vec{L}_{spin}$$

$$\dot{\vec{L}}_{orbit} = \dot{\vec{R}} \times \vec{p}_{tot} + \vec{R} \times \dot{\vec{p}}_{tot} = \vec{R} \times \vec{F}^{ext}$$

"0"

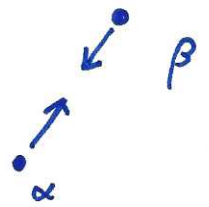
$$\dot{\vec{L}}_{tot} = \sum_{\alpha=1}^N \vec{r}_\alpha \times (\vec{F}_\alpha^{ex} + \sum_{\beta \neq \alpha} \vec{F}_{\alpha\beta}) = \sum_{\alpha=1}^N \vec{r}_\alpha \times \vec{F}_\alpha^{ex} + \sum_{\alpha \neq \beta} \vec{r}_\alpha \times \vec{F}_{\alpha\beta}$$

$$= \sum_{\alpha=1}^N \vec{r}_\alpha \times \vec{F}_\alpha^{ex} + \frac{1}{2} \sum_{\alpha \neq \beta} (\vec{r}_\alpha \times \vec{F}_{\alpha\beta} + \vec{r}_\beta \times \vec{F}_{\beta\alpha})$$

Due to Newton's 3rd law $\vec{F}_{\alpha\beta} = -\vec{F}_{\beta\alpha}$

$$\sum_{\alpha \neq \beta} (\vec{r}_\alpha \times \vec{F}_{\alpha\beta} + \vec{r}_\beta \times \vec{F}_{\beta\alpha}) = \sum_{\alpha \neq \beta} (\vec{r}_\alpha - \vec{r}_\beta) \times \vec{F}_{\alpha\beta}$$

for central force $\vec{F}_{\alpha\beta} \parallel \vec{r}_\alpha - \vec{r}_\beta \Rightarrow \parallel 0$



$$\Rightarrow \dot{\vec{L}}_{tot} = \sum_{\alpha=1}^N \vec{r}_\alpha \times \vec{F}_\alpha^{ex} = \dot{\vec{P}}^{ex}$$

$$= \sum_{\alpha=1}^N (\vec{r}'_\alpha + \vec{R}) \times \vec{F}_\alpha^{ex} = \dot{\vec{L}}_{orbit} + \sum_{\alpha=1}^N \vec{r}'_\alpha \times \vec{F}_\alpha^{ex}$$

$$\Rightarrow \boxed{\dot{\vec{L}}_{spin} = \sum_{\alpha=1}^N \vec{r}'_\alpha \times \vec{F}_\alpha^{ex}} \leftarrow \text{torque respect to CM.}$$

In QM, electron has both orbit angular momentum + spin.

spin is quantized at $\hbar/2$.

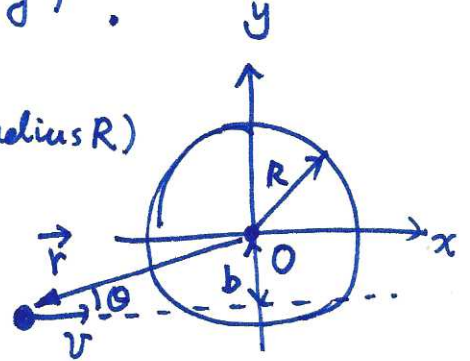
§ Moment of Inertia (fixed axis rotation)

$L_z = I_{zz} \omega$, where ω is along z-axis.

$$I_{zz} = \sum_{\alpha} m_{\alpha} (x_{\alpha}^2 + y_{\alpha}^2) = \int \rho(x^2 + y^2) dx dy dz$$

Examples: ① A uniform circular wheel (mass M , radius R)

A lump of putty (mass m), speed v with a line of approach of distance b .



$$L_z^{in} = mrv \sin \theta = mrb$$

$$L_z^{fin} = I_{zz} \omega = (I_{wheel} + mR^2) \omega$$

$$I_{wheel} = \int_0^R \rho r^3 dr d\theta = \rho \int_0^{2\pi} d\theta \cdot \frac{R^3}{4} = \frac{\pi}{2} \rho R^3 = \frac{MR^2}{2}$$

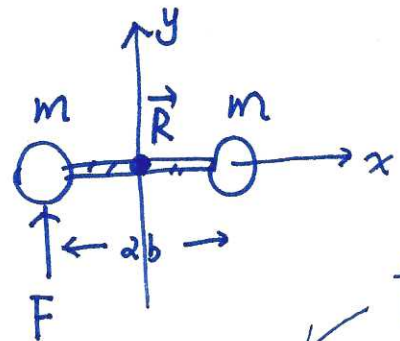
$$M = \int_0^R \rho r^2 dr d\theta = \rho \pi R^2$$

$$\Rightarrow L_z^{fin} = (m + M/2) R^2 \omega \Rightarrow \omega = \frac{m}{m + M/2} \frac{vb}{R^2}$$

② A sliding and spinning dumbbell

At time $t=0$, an moment of impulse F is acted on the left mass for time t .

What will happen?



$$\dot{\vec{p}} = F \Rightarrow \Delta p = Ft$$

The center of mass will move at $\vec{V}_{cm} = \dot{\vec{R}} = \frac{Ft}{2m}$

\vec{F} also has torque: with respect to \vec{R}

$$\vec{F} = F \cdot b, \quad I = 2Mb^2 \quad \Rightarrow \quad \omega = \frac{\tau_{\text{net}}}{I} = \frac{F \cdot b \cdot t}{2Mb}$$

so

$$\begin{cases} v_{\text{left}} = v_{\text{cm}} + \omega b = \frac{F \cdot b \cdot t}{m} \\ v_{\text{right}} = v_{\text{cm}} - \omega b = 0 \end{cases}$$