

Prob 1)

$$\textcircled{1} \quad F = eE - \frac{m\ddot{v}}{2} = m\ddot{v}$$

$$\Rightarrow \ddot{v} = \frac{eE}{m} - \frac{v}{2} \quad \Rightarrow \frac{v_\infty}{2} = \frac{eE}{m} \text{ with } v_\infty = \frac{eE\tau}{m}$$

$$\Rightarrow v(t) = A e^{-t/\tau} + v_\infty$$

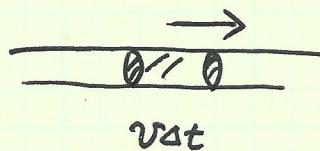
$$\Rightarrow v(t) = (v_0 - v_\infty) e^{-t/\tau} + v_\infty,$$

$v_\infty$  does not depend on  $v_0$ .

\textcircled{3} Consider the long time limit  $t \rightarrow \infty, v(t) \rightarrow v_\infty$ .

then  $j = nev_\infty = \frac{e^2 n E \tau}{m}$

$$\Rightarrow \boxed{\sigma = \frac{ne^2 \tau}{m}}$$



$$\textcircled{2} \quad I = \frac{\Delta Q}{\Delta t} = \frac{env\Delta t A}{\Delta t} = envA$$

$$j = I/A = env$$

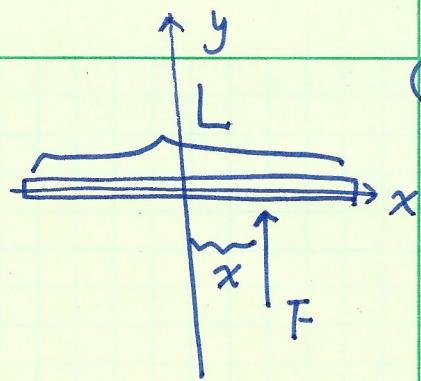
at  $t \gg \tau$ ,  $j$  becomes steady, hence  $v = v_\infty$

$$j = env_\infty = \frac{e^2 n E \tau}{m}$$

Prob 3)

$$\textcircled{1} \quad \vec{P} = F \sin t \hat{y}$$

$$\Rightarrow \textcircled{2} \quad M \vec{v}_{cm} = \vec{P} \Rightarrow \vec{v}_{cm} = \frac{F \sin t}{M} \hat{y}$$



\textcircled{2} in the center of mass frame, the rod becomes to rotate around the z-axis.

$$\frac{d\vec{L}_{cm}}{dt} = \vec{P}_{cm}, \quad \vec{P}_{cm} = \vec{r} \times \vec{F} = x F \hat{z}$$

The moment of inertia in the center of mass frame

$$I = 2 \int_{-L/2}^{L/2} p x^2 dx = \frac{1}{3} p x^3 \Big|_{-L/2}^{L/2} = \frac{2}{3} p (L/2)^3 = \frac{1}{12} p L^3$$

$$M = \int p \star dx = pL$$

$$\Rightarrow I = \frac{1}{12} M L^2$$

$$\Rightarrow I \omega = \Gamma_{cm} \sin t \Rightarrow \omega = \frac{12 x F \sin t}{M L^2} \hat{z}$$

\textcircled{3} Left end  $\vec{v}_L = \vec{v}_{cm} - \frac{L}{2} \omega \hat{z} + \vec{\omega} \times \vec{r}$

$\Rightarrow \vec{v}_L$  and  $\vec{v}_R$  are along  $\hat{y}$ -direction in the Lab frame.

in the CM frame  $\vec{v}_{L,cm} = -\hat{y} \omega \frac{L}{2} = -\hat{y} \frac{6 x F \sin t}{M L}$

\textcircled{4}  $\vec{v}_{R,cm} = \hat{y} \omega \frac{L}{2} = \hat{y} \frac{6 x F \sin t}{M L}$

$$\vec{v}_L = \vec{v}_{Cm} + \vec{v}_{L,Cm} = \frac{F_{st}}{M} \left(1 - \frac{6x}{L}\right) \hat{y}$$

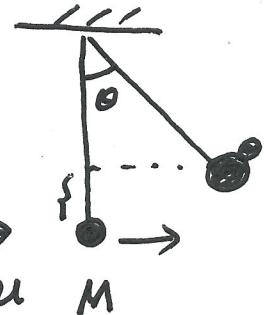
$$\vec{v}_R = \vec{v}_{Cm} + \vec{v}_{R,Cm} = \frac{F_{st}}{M} \left(1 + \frac{6x}{L}\right) \hat{y}$$

Prob 3): during the collision, the pendulum

does not have time to move. There's no

horizontal external force. Hence the momentum  $\overset{\rightarrow}{m u}$

along the horizontal direction is conserved.



$$m u = (m+M) v_0 \Rightarrow v_0 = \frac{m u}{m+M}$$

After the collision, the mechanical energy is conserved.

$$E_K = \frac{1}{2} (m+M) v_0^2, \quad E_G = (M+m) g l (1 - \cos \theta)$$

$\Rightarrow$  At the  $\theta = \theta_m$ ,  $E_K = 0$

$$(M+m) g l (1 - \cos \theta_m) = \frac{1}{2} (m+M) v_0^2$$

$$1 - \cos \theta_m = \frac{v_0^2}{2 g l} \Rightarrow \cos \theta_m = 1 - \frac{u^2}{2 g l} \left(\frac{m}{m+M}\right)^2$$

$$\theta \theta_m = \cos^{-1} \left[ 1 - \frac{u^2}{2 g l} \left(\frac{m}{m+M}\right)^2 \right]$$

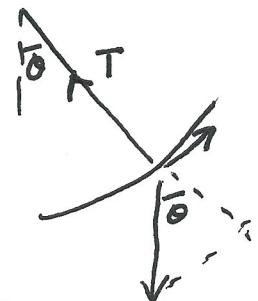
(4)

b) at a general angle  $\theta$

$$(m+M)gl(1-\cos\theta) + \frac{1}{2}(m+M)v^2 = \frac{1}{2}(m+M)v_0^2$$

$$v^2 = v_0^2 - 2gl(1-\cos\theta) = \frac{m^2 u^2}{(m+M)^2} - 2gl(1-\cos\theta)$$

$$T - G \cos\theta = (M+m)v^2/l$$



$$T = (M+m)g \cos\theta + \frac{(M+m)}{l} [v_0^2 - 2gl(1-\cos\theta)] \quad G$$

$$= (M+m)g [-2 + 3\cos\theta] + \frac{m}{M+m} \frac{m u^2}{l}$$